

AN H^1 -GALERKIN METHOD FOR A STEFAN PROBLEM WITH A QUASILINEAR PARABOLIC EQUATION IN NON-DIVERGENCE FORM

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ABSTRACT. Optimal error estimates in L^2 , H^1 and H^2 -norms are established for a single phase Stefan problem with quasilinear parabolic equation in non-divergence form by an H^1 -Galerkin procedure.

KEY WORDS AND PHRASES. H^1 -Galerkin procedure, finite element approximation, nonlinear Stefan problem, non-divergence form, error analysis.

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1. INTRODUCTION.

With the help of Galerkin finite element methods, Nitsche in his pioneering works [1]-[3] established error estimates for linear problems, proposed earlier by Magenes [4]. We extended his analysis to nonlinear problems in divergence form [5]-[6]. In the present work, a single phase Stefan problem with quasilinear parabolic equation in non-divergence form is considered and under appropriate conditions optimal error estimates for Galerkin approximation in L^2 , H^1 as well as H^2 norms are established. We require more regularity assumptions for the present one than for the cases discussed in [5]-[6], and consequently we improve upon the estimates in L^2 -norm.

The organization of the paper is as follows: In section 2, the description of the problem and the transformed system with some preliminaries are presented. The weak formulation and H^1 -Galerkin procedure are discussed in section 3. Section 4 deals with an auxiliary projection and some approximation Lemmas. In section 5 optimal error estimates in L^2 , H^1 and H^2 -norms for continuous time Galerkin approximations are established, assuming existence of the approximate solution. Finally, in section 6 the question of global existence and uniqueness of the Galerkin approximation is discussed.

2. PROBLEM DESCRIPTION AND DOMAIN FIXING.

The nonlinear heat conduction with change of phase can be modelled as a single phase nonlinear Stefan problem in a variable domain $\Omega(\tau) \times (0, T_0)$, where $\Omega(\tau) = \{y \in (0, S(\tau))\}$ and $S(\tau)$ known to be the free boundary. We state this problem as follows:

Find a pair $\{U, S\}$, $U = U(y, \tau)$, $S = S(\tau)$ such that U satisfies

$$U_\tau - a(U) U_{yy} = 0, \text{ for } (y, \tau) \in \Omega(\tau) \times (0, T_0] \tag{2.1}$$

with initial and boundary conditions

$$U(y, 0) = g(y), \text{ for } y \in I = (0, 1) \tag{2.2}$$

$$U_y(0, \tau) = 0$$

for $\tau > 0$

for $\tau > 0$

$$U(S(\tau), \tau) = 0, \tag{2.3}$$

and S , the free boundary satisfies

$$S_\tau = -U_y(S(\tau), \tau), \text{ for } \tau > 0 \tag{2.4}$$

with $S(0) = 1$. The above problems is a special case of the general situation discussed in Fasano et. al. [7], where 'a' depends only on U , $q \equiv 0$ and $\phi = -U_y(S(\tau), \tau)$ in their notations.

We use the following notiations. Let $\Omega(\tau) \subset \mathbb{R}$ be a bounded domain for $\tau \geq 0$. Let $(u, v) = \int_{\Omega(\tau)} uv \, dx$ and $\|u\|^2 = (u, u)$. For each nonnegative integer m , let $H^m(\Omega(\tau))$ be the usual Sobolev space $W^{m, p}(\Omega(\tau))$, for $p = 2$ with the norm

$$\|u\|_{H^m(\Omega(\tau))}^2 = \sum_{j=0}^m \left\| \frac{\partial^j u}{\partial x^j} \right\|^2 \, dx.$$

Further, $W^{m, \infty}(\Omega(\tau))$ is defined as usual with the norm

$$\|u\|_{W^{m, \infty}(\Omega(\tau))} = \sum_{j=0}^m \left\| \frac{\partial^j u}{\partial x^j} \right\|_{L^\infty(\Omega(\tau))}.$$

In case $I = \Omega(\tau)$, we shall omit I from $H^m(I)$, $L^\infty(I)$ and $W^{m, \infty}(I)$ and norm $H^m(I)$ is denoted by $\|\cdot\|_m$.

If X be a normed linear space with norm $\|\cdot\|_X$ and $\phi: (a, b) \rightarrow X$, then we denote by

$$\|\phi\|_{W^{k, q}(a, b; X)}^q = \sum_{\beta=0}^k \left\| \frac{\partial^\beta \phi}{\partial t^\beta} \right\|_{L^q(a, b; X)}^q, \quad 1 \leq q < \infty$$

and

$$\|\phi\|_{W^{k, q}(a, b; X)} \text{ is accordingly defined.}$$

In case $(a,b) = (0,T)$ and $X = H^m$ or $W^{m,\infty}$, we write simply $\|\phi\|_{W^{k,q}(H^m)}$ for $\|\phi\|_{W^{k,q}(0,T;H^m(I))}$ or $\|\phi\|_{W^{k,q}(W^{m,\infty})}$ for $\|\phi\|_{W^{k,q}(0,T;W^{m,\infty}(I))}$. For convenience, we use $\phi_x = \frac{\partial\phi}{\partial x}$, $\phi_{xx} = \frac{\partial^2\phi}{\partial x^2}$; $\phi_t = \frac{\partial\phi}{\partial t}$ and $\phi(1) = \phi(1,t)$, if $\phi = \phi(x,t)$.

Throughout this work, K will always denote a generic constant. On occasion, we will show that a constant depends on certain parameters, while independent of others.

We shall now state our main assumption on $a(\cdot)$, g and the solution U,S , and call them collectively 'condition B'.

CONDITION B.

- (i) For $p \in R$, $a(p) \geq \alpha$, where α is a positive constant.
- (ii) For $p \in R$, $a(p) \in C^3(R)$ and there is a common bound $K_1 > 0$ such that $|a|$, $|a_p|$, $|a_{pp}|$ and $|a_{ppp}| \leq K_1$.
- (iii) The initial function g is sufficiently smooth and satisfies the compatibility condition that is $g_y(0) = g(1) = 0$.
- (iv) The problem (2.1)-(2.4) has a unique solution.

For the existence and uniqueness of the solution of (2.1)-(2.4), see Fasano et.al. [7].

Further it is assumed that the solution U,S of (2.1)-(2.4) satisfies the following regularity condition. For an integer $r \geq 1$,

$$U \in L^\infty(0,T_0; H^{r+1}(\Omega(\tau))) \cap W^{1,2}(0,T_0; H^{r+1}(\Omega(\tau))) \cap W^{1,\infty}(0,T_0; W^{2,\infty}(\Omega(\tau))),$$

$$S \in W^{1,\infty}(0,T_0).$$

Let \tilde{K}_2 be the bound for the functions in above mentioned norms.

We fix the free boundary, using Landau type transformation [8]

$$x = S^{-1}(\tau)y, \quad \tau \geq 0. \tag{2.5}$$

Further, we introduce an additional transformation in time scale given by

$$t = t(\tau) = \int_0^\tau S^{-2}(\tau')d\tau', \tag{2.6}$$

in order to decouple the resulting transformed system. A routine calculation shows that the function $u(x,t) = U(y,\tau)$ satisfies

$$u_t - a(u)u_{xx} = -u_x(1)x u_x, \quad x \in I, \quad t \in (0,T] \tag{2.7}$$

$$u(x,0) = g(x), \quad x \in I; \tag{2.8}$$

$$u_x(0,t) = u(1,t) = 0, \quad t > 0 \tag{2.9}$$

and the function $s(t) = S(\tau)$ satisfies

$$\frac{ds}{dt} = -u_x(1)s, \quad t > 0 \tag{2.10}$$

with $s(0) = 1$.

Here, $t = T$ corresponds to $\tau = T_0$. Note that all the regularity assumptions for U, S are carried over to u, s with the bound say K_2 and the new regularity assumptions are collectively called R_1 . Further, the integral (2.6) can be rewritten as

$$\frac{d\tau}{dt} = s^2(t), \quad \text{with} \quad \tau(0) = 0. \tag{2.11}$$

3. WEAK FORMULATION AND H^1 -GALERKIN PROCEDURE.

Consider the space:

$$H^2(I) = \{v \in H^2(I) : v_x(0) = v(1) = 0\},$$

The weak formulation of (2.7)-(2.9) is given by

$$(u_{tx}, v_x) + (a(u)u_{xx}, v_{xx}) = u_x(1)(xu_x, v_{xx}), \quad v \in H^2(I) \quad \text{and} \quad t > 0 \tag{3.1}$$

with $u(x,0) = g(x)$.

H^1 -Galerkin Procedure. Let S_h^0 ($0 < h \leq 1$) be a finite dimensional subspace of $H^2(I)$ belonging to regular $S_h^{0r,2}$ family, for a definition see Oden et. al. [9] and

satisfying the following approximation and inverse properties:

(i) For $v \in H^m(I)$ $H^2(I)$, there is a constant K_0 independent of h such that

$$\inf_{\chi \in S_h^0} \|v - \chi\|_j \leq K_0 h^{m-j} \|v\|_m, \quad \text{for } j = 0, 1, 2 \text{ and } 2 \leq m \leq r+1;$$

(ii) For $\chi \in S_h^0$, $\|\chi\|_2 \leq K_0 h^{-1} \|\chi\|_1$.

Now we call $u^h : (0, T] \rightarrow S_h^0$ an H^1 -Galerkin approximation of u , if it satisfies

$$(u_{tx}^h, \chi_x) + (a(u^h)u_{xx}^h, \chi_{xx}) = u_x^h(1)(xu_x^h, \chi_{xx}), \quad \chi \in S_h^0 \tag{3.2}$$

and the initial condition

$$u^h(x, 0) = Q_h g(x), \tag{3.3}$$

where Q_h is an appropriate projection of u onto S_h^0 at $t = 0$, to be defined later. Further, the Galerkin approximations s_h and τ_h of s and τ respectively are given by

$$\frac{ds_h}{dt} = -u_x^h(1)s_h, \quad \text{with } s_h(0) = 1 \tag{3.4}$$

and

$$\frac{d\tau_h}{dt} = s_h^2, \quad \text{with } \tau_h(0) = 0. \tag{3.5}$$

4. SOME APPROXIMATION LEMMAS.

Set

$$A(u; v, w) = (a(u)v_{xx}, w_{xx}) - u_x(1)(xv_x, w_{xx}); \text{ for } u \in W^{1, \infty}, v \text{ and } w \in H^2. \tag{4.1}$$

The boundedness and Garding type inequality for A can be established by standard arguments.

LEMMA 4.1. For $u \in W^{1, \infty}, v$ and $w \in H^2(I)$

$$|A(u; v, w)| \leq M \|v_{xx}\| \|w_{xx}\| \tag{4.2}$$

and

$$A(u; v, v) \geq \tilde{\alpha} \|v_{xx}\|^2 - \rho \|v_x\|^2, \tag{4.3}$$

where $M, \tilde{\alpha}$ and ρ are constants, but M and ρ may depend on $\|u_x\|_{L^\infty}$.

Define

$$A_p(u; v, w) = A(u; v, w) + \rho(v_x, w_x).$$

Note that $A_p(u, \dots)$ is coercive in H^2 , that is

$$A_p(u; v, v) \geq \tilde{\alpha} \|v_{xx}\|^2. \tag{4.4}$$

Let $\tilde{u} \in S_h^0$ be an approximation of u with respect to the form A_p :

$$A_p(u; u - \tilde{u}, \chi) = 0, \quad \chi \in S_h^0, \tag{4.5}$$

Now, an application of Lax-Milgram theorem shows the existence of a unique solution \tilde{u} of equation (4.5).

Consider

$$L^*(u)\phi = (a(u)\phi_{xx}, \phi_{xx}) + u_x(1)(x\phi_{xx}, \phi_x) - \rho\phi_{xx}, \quad u \in H^2. \tag{4.6}$$

For $\Psi \in L^2(I)$, define $\phi \in H^4 \cap H^2$ by

$$L^*(u)\phi = \Psi; \quad x \in I \tag{4.7}$$

$$\phi_{xx}|_{x=1} = \phi_{xxx}|_{x=0} = 0.$$

Then, for $v \in H^2_0(I)$ we get

$$(v, L^*(u)\phi) = A_\rho(u;v,\phi). \tag{4.8}$$

Thus, defining $D(L^*)$ as

$$D(L^*) = \{\phi \in H^4 \cap H^2: \phi_{xx}(1) = \phi_{xxx}(0) = 0\},$$

we have from the positivity and boundedness of A_ρ that at least a weak solution $\phi \in D(L^*)$ of (4.7) for each $\Psi \in L^2$ exists and the regularity

$$\|\phi\|_4 \leq C_0 \|\Psi\|, \tag{4.9}$$

where C_0 depends on u and its derivatives, holds.

Let $\eta = u - \tilde{u}$. We now need to obtain some estimates of η and its temporal derivatives η_t , for our future use. The following Lemma proves very convenient for our purpose.

LEMMA 4.2. Let $\psi \in H^2_0(I)$ and satisfy

$$A_\rho(u;\psi,\chi) = F(\chi), \chi \in S_h, \tag{4.10}$$

where $F: H^2_0(I) \rightarrow \mathbb{R}$ and linear. Let there exist constants M_1 and M_2 with $M_1 \geq M_2$ such that

$$|F(\phi)| \leq M_1 \|\phi_{xx}\|, \quad \phi \in H^2_0 \tag{4.11}$$

and

$$|F(\phi)| \leq M_2 \|\phi\|_4, \quad \phi \in D(L^*). \tag{4.12}$$

Then, for sufficiently small h

$$\|\phi_{xx}\| \leq K_3 [M_1 + \inf_0 \|\phi - \chi\|_2] \tag{4.13}$$

$$\chi \in S_h$$

and

$$\|\psi\| \leq K_3 [(M_1 + \inf_0 \|\phi - \chi\|_2)h^2 + M_2], \tag{4.14}$$

$$\chi \in S_h$$

where $K_3 = K_3(\alpha, \rho, M, C_0; K_0)$ is used as generic constant.

PROOF. Note that

$$A_\rho(u;\psi,\psi) = A_\rho(u;\psi,\psi-\chi) - F(\psi-\chi) + F(\psi), \chi \in S_h.$$

By coercive property (4.4) for A_ρ , we get

$$\|\psi_{xx}\|^2 \leq (\tilde{\alpha})^{-1} [M \|\psi_{xx}\| + M_1 \inf_0 \|\psi-\chi\|_2 + M_1 \|\psi_{xx}\|].$$

$$\chi \in S_h$$

For $\phi \in H^2_0$, $\|\phi\|_2 \leq K \|\phi_{xx}\|$. Therefore,

$$\|\phi_{xx}\| \leq (\tilde{\alpha})^{-1} [M \inf_{\chi \in S_h} \|\phi - \chi\|_2 + 2KM_1]$$

and the estimate (4.13) follows. In order to get an L^2 -estimate, we follow here Aubin-Nitsche duality arguments. For $\Psi \in L^2(I)$, define $\phi \in D(L^*)$ by (4.7). Multiply both the sides of (4.7) by ϕ to obtain for $u \in H^2_0$,

$$\begin{aligned} (\phi, \Psi) &= (\phi, L^*(u)\phi) = A(u; \phi, \phi) \\ &= A(u; \phi, \phi - \chi) + F(\chi - \phi) + F(\phi) \\ &\leq [M \|\phi_{xx}\| \inf_{\chi \in S_h} \|\phi - \chi\|_2 + M_1 \inf_{\chi \in S_h} \|\phi - \chi\|_2] + M_2 \|\phi\|_4 \\ &\leq [M \|\phi_{xx}\| K_0 h^2 + M_1 K_0 h^2 + M_2] \|\phi\|_4. \end{aligned} \tag{4.15}$$

From (4.9), (4.13) and (4.15), we obtain the required estimate (4.14).

The next Lemma contains the error estimates related to η and η_t .

Lemma 4.3. For $t \in [0, T]$, the following estimates

$$\|\eta\|_j \leq K_4 h^{m-j} \|u\|_m \tag{4.16}$$

and

$$\begin{aligned} \|\eta_t\|_j &\leq K_5 h^{m-j} (\|u\|_m + \|u_t\|_m), \\ j &= 0, 1, 2 \text{ and } 2 \leq m \leq r+1, \end{aligned} \tag{4.17}$$

hold. Here K_4 and K_5 are positive constants depending on parameters expressed through the following expressions that is

$$K_4 = K_4(K_0, K_3) \text{ and } K_5 = K_5(K_0, K_1, K_3, K_4, \|u_t\|_{w, 2, \infty} \text{ and } \|u\|_{w, 2, \infty}).$$

PROOF. Put $\psi = \eta$ and $F = 0$ in the previous Lemma 4.2 to get

$$\begin{aligned} \|\eta_{xx}\| &\leq K_3 \inf_{\chi \in S_h} \|\eta - \chi\|_2 \\ &\leq K_3 \inf_{\chi \in S_h} \|u - v\|_2, \quad v = \chi + u \in S_h \\ &\leq K_0 K_3 h^{m-2} \|u\|_m, \quad 2 \leq m \leq r+1. \end{aligned}$$

For $\eta \in H^2_0$, $||\eta|| \leq ||\eta_x||$ and $||\eta_x|| \leq ||\eta_{xx}||$. Hence the result (4.16) for $j = 2$. Similarly, we get the estimate (4.16), for $j = 0$, consequently, the estimate for $||\eta||_1$ follows from the interpolation inequality,

$$||\eta||_1 \leq ||\eta||^{1/2} ||\eta||_2^{1/2}.$$

In order to estimate η_t , we differentiate (4.5) with respect to 't' and obtain

$$A_p(u; \eta_t, \chi) = - \left(\frac{da(u)}{dt} \right) \eta_{xx}, \chi_{xx} + u_{tx}(1) (\chi \eta_x, \chi_{xx}). \tag{4.18}$$

Identifying the right hand side of (4.18) with $F(\chi)$, we see that for $\phi \in H^2_0(I)$

$$|F(\phi)| \leq K_6 ||\eta_{xx}|| ||\phi_{xx}||,$$

where K_6 depends on K_1 and $||u_t||_{W^{1,\infty}}$.

Further, for $\phi \in D(L^*)$ and $u \in H^2_0$, we get on integration by parts

$$\begin{aligned} F(\phi) &= (\eta_x, (a_t(u)\phi_{xx})_x) - u_{tx}(1)(\eta, (\chi\phi_{xx})_x) \\ &= - (\eta, (a_t(u)\phi_{xx})_{xx}) + u_{tx}(1)(\eta, (\chi\phi_{xx})_x) \end{aligned}$$

and

$$|F(\phi)| \leq K_7 ||\eta|| ||\phi||_4,$$

where $K_7 = K_7(K_1, ||u_{txx}||_{L^\infty}, ||u_{xx}||_{L^\infty}$ and $||u_t||_{W^{1,\infty}})$.

Thus, Lemma 4.2 is applicable with $M_1 = K_6 ||\eta_{xx}||$ and $M_2 = K_7 ||\eta||$ and we get the desired estimate (4.17) for $j = 0, 2$. For $j = 1$, as usual we make use of the interpolation inequality. We shall also need later the following estimate for $\eta_x(1)$.

LEMMA 4.4. There is a constant $K_8 = K_8(\alpha, K_0, M; K_4)$ such that for $2 \leq m \leq r+1$.

$$|\eta_x(1)| \leq K_8 n^{2(m-2)} ||u||_m.$$

PROOF. Define an auxiliary function $\phi \in H^4 \cap H^2_0$ as a solution of

$$L^*(u) \phi = 0, \chi \in I$$

$$\phi_{xxx}|_{x=0} = 0;$$

$$\phi_{xx}|_{x=1} = 1.$$

Multiplying by η the first equation and integrating by parts, we obtain

$$\begin{aligned} \alpha |\eta_x(1)| &\leq |a(0) \eta_x(1)| \leq |A_\rho(u; \eta, \phi)| \\ &\leq A_\rho(u; \eta, \phi - \chi), \chi \in S_h^0 \\ &\leq M \|\eta\|_2 \inf_{\chi \in S_h} \|\phi - \chi\|_2 \\ &\leq MK_4 K_0 h^{2(m-2)} \|u\|_m \|\phi\|_m \end{aligned}$$

Hence, the result follows.

5. A PRIORI ERROR ESTIMATES FOR CONTINUOUS TIME GALERKIN APPROXIMATION.

Throughout this section, we assume that there are positive constants K^* and h_0 such that a Galerkin approximation $u^h \in S_h^0$ in (3.2) exists and satisfies,

$$\|u^h\|_{K(H^2)} \leq K^*, \text{ for } 0 < h \leq h_0, \tag{5.1}$$

where $u^h(x, 0)$ is defined as $Q_h g$, satisfying

$$A_\rho(g; g - Q_h g, \chi) = 0, \chi \in S_h^0. \tag{5.2}$$

Clearly, $u^h(x, 0) \equiv \tilde{u}(x, 0)$.

Let $\zeta = u^h - \tilde{u}$ and $e = u - u^h = \eta - \zeta$.

THEOREM 5.1. Suppose $\eta = u - \tilde{u}$ satisfies (4.5) and u^h , the Galerkin approximation of u is defined by (3.2) with Q_h given as in (5.2). Further, assume that (5.1) holds. Then, there is a constant $K_9 = K_9(\alpha, \rho, K^*, K_1, K_4, K_5 \text{ and } K_8)$ such that for $m \geq 4$

$$\|\zeta_x\|_{L^\infty(L^2)} + \beta \|\zeta_{xx}\|_{L^2(L^2)} \leq K_9 h^{\frac{m}{2}} (\|u_t\|_{L^2(H^m)} + \|u\|_{L^2(H^m)}). \tag{5.3}$$

PROOF. From (4.5) and (3.1) with $v = \chi$, we get

$$(\tilde{u}_{tx}, \chi_x) + A_\rho(u; \tilde{u}, \chi) = -(\eta_{tx}, \chi_x) + \rho(u_x, \chi_x), \chi \in S_h^0.$$

Subtracting this from (3.2), we obtain

$$\begin{aligned} (\zeta_{tx}, \chi_x) + A_\rho(u^h; u^h, \chi) - A_\rho(u; \tilde{u}, \chi) &= (\eta_{tx}, \chi_x) - \rho(\eta_x, \chi_x) \\ &\quad + \rho(\zeta_x, \chi_x) \end{aligned} \tag{5.4}$$

But

$$\begin{aligned}
 A_\rho(u^h; u^h, \chi) - A_\rho(u; \tilde{u}, \chi) &= (a(u^h) \zeta_{xx}, \chi_{xx}) \\
 &+ ([a(u^h) - a(u)] \tilde{u}_{xx}, \chi_{xx}) - u_x(1)(x\zeta_x, \chi_{xx}) + \eta_x(1)(xu_x^h, \chi_{xx}) \\
 &- \zeta_x(1)(xu_x^h, \chi_{xx}) + \rho(\zeta_x, \chi_x). \tag{5.5}
 \end{aligned}$$

From (5.4)-(5.5) with $\chi = \zeta$, it follows on integrating by parts with respect to x the two terms on the right hand side of (5.4),

$$\begin{aligned}
 \frac{1}{2} \frac{d}{dt} \|\zeta_x\|^2 + (a(u^h) \zeta_{xx}, \zeta_{xx}) &= -(\eta_t, \zeta_{xx}) + \rho(\eta, \zeta_{xx}) + u_x(1)(x\zeta_x, \zeta_{xx}) \\
 &+ ([a(u) - a(u^h)] \tilde{u}_{xx}, \zeta_{xx}) - \eta_x(1)(xu_x^h, \zeta_{xx}) + \zeta_x(1)(xu_x^h, \zeta_{xx}).
 \end{aligned}$$

Using $a(\cdot) \geq \alpha$, (5.1) and replacing \tilde{u} by $u - \eta$, we obtain

$$\begin{aligned}
 \frac{1}{2} \frac{d}{dt} \|\zeta_x\|^2 + \alpha \|\zeta_{xx}\|^2 &\leq \{ \|\eta_t\| + \rho \|\eta\| + K_2 \|\zeta_x\| + K_2(\|\eta\|_{L^\infty} + \|\zeta\|_{L^\infty}) \\
 \|\eta\|_2 + K_1 K_2(\|\eta\| + \|\zeta\|) + K^* |\eta_x(1)| \} &\|\zeta_{xx}\| + K^* |\zeta_x(1)| \|\zeta_{xx}\|. \tag{5.6}
 \end{aligned}$$

Since $|\zeta_x(1)| \leq \|\zeta_x\|^{1/2} \|\zeta_{xx}\|^{1/2}$ for $\zeta \in H^2_0$, applying Young's inequality for the last term and the inequality $ab \leq \frac{a^2}{2\epsilon} + \frac{\epsilon}{2} b^2$, $a, b \geq 0$; $\epsilon > 0$ for the remaining terms in (5.6), we get using $\|\phi\|_{L^\infty} \leq \|\phi_x\|$ for $\phi \in H^2_0$

$$\begin{aligned}
 \frac{1}{2} \frac{d}{dt} \|\zeta_x\|^2 + \alpha \|\zeta_{xx}\|^2 &\leq K_{10}(\epsilon) \|\zeta_{xx}\|^2 + K(K_1, K_2, K^*, \rho; \epsilon) (\|\eta_t\|^2 \\
 &+ \|\eta\|^2 + |\eta_x(1)|^2 + \|\eta\|_1^2 \|\eta\|_2^2) \\
 &+ K(K_2; \epsilon) \|\eta\|_2^2 \|\zeta_x\|^2 + K(K_1, K_2, K^*, \epsilon) \|\zeta_x\|^2.
 \end{aligned}$$

Now with appropriate choice of ϵ , $K_{10}(\epsilon)$ can be made less than or equal to $\alpha/2$. With this choice of ϵ , we get by integrating with respect to 't' and using Gronwall's inequality the following

$$\begin{aligned}
 \|\zeta_x\|^2(t) + \alpha \int_0^t \|\zeta_{xx}\|^2 dt' &\leq K(K_1, K_2, K^*; \rho) \int_0^t (\|\eta_t\|^2 + \|\eta\|^2 \\
 &+ |\eta_x(1)|^2 + \|\eta\|_1^2 \|\eta\|_2^2) dt'.
 \end{aligned}$$

From (4.16)-(4.17) and (4.19) with $2(m-2) \geq m$ and $2m - 3 \geq m$ that is $m \geq 4$, we get the desired estimate (5.3).

COROLLARY 5.2. Let all the assumptions of the previous theorem hold and the finite dimensional subspace S_h^0 satisfy the inverse property. Then there is a constant K_{11} depending on K_9 and K_0 such that for $r + 1 \geq m \geq 4$,

$$\begin{aligned} & \|r\|_{L^\infty(L^2)} + \|r\|_{L^\infty(H^1)} + h \|r\|_{L^\infty(H^2)} \\ & \leq K_{11} h^m (\|u\|_{L^2(H^m)} + \|u_t\|_{L^2(H^m)}). \end{aligned} \tag{5.7}$$

PROOF. From the estimate (5.3) and $\|\zeta\| \leq \|\zeta_x\|$ for $\zeta \in S_h^0$, we get

$$\|\zeta\|_{L^\infty(L^2)} + \|\zeta\|_{L^\infty(H^1)} \leq K_{11} h^m (\|u\|_{L^2(H^m)} + \|u_t\|_{L^2(H^m)}).$$

By inverse property for S_h^0 , we have

$$\|\zeta\|_{L^\infty(H^2)} \leq K_0 h^{-1} \|\zeta\|_{L^\infty(H^1)}, \quad \zeta \in S_h^0.$$

Hence the result (5.7). From Theorem 5.1, Corollary 5.2, Lemma 4.3 and triangle inequality we get the following theorem.

THEOREM 5.3. Let the solution u of (2.7)-(2.9) satisfy the regularity condition R_1 . Further, suppose that there are positive constants h_0 and K^* ($K^* \geq 2K_2$) such that an approximate solution $u^h \in S_h^0$ of (3.2) satisfying (5.1) exists in $I \times (0, T]$ for $0 < h \leq h_0$. Then, the following estimates hold for $r \geq 3$,

$$\|e\|_{L^\infty(H^j)} \leq K_{12} h^{r+1-j}, \quad j = 0, 1, 2, \tag{5.8}$$

where $K_{12} = K_{12}(K_4, K_{11}$ and $K_2)$. Besides, for sufficiently small h and $r \geq 3$,

$$\|u^h\|_{L^\infty(H^2)} \leq 2K_2 \leq K^* \tag{5.9}$$

and consequently, K_{12} can be chosen independent of K^* .

PROOF. The estimates (5.8) for $j = 0, 1, 2$ are immediate from the Theorem 5.1, Corollary 5.2 and Lemma 4.3 by triangle inequality. To prove (5.9), we note

$$\begin{aligned} \|u^h\|_{L^\infty(H^2)} & \leq \|u\|_{L^\infty(H^2)} + \|e\|_{L^\infty(H^2)} \\ & \leq K_2 + K_{12} h^{r-1} \end{aligned}$$

$\leq 2K_2$, for sufficiently small h and $r \geq 3$.

We are now looking for approximations of U and S , where the pair $\{U, S\}$ is the solution of (2.1)-(2.4). The Galerkin approximations U^h and S_h are given by

$$U^h(y, \tau) = u^h(x, t) \tag{5.10}$$

$$S^h(\tau) = s_h(t) \tag{5.11}$$

where,

$$y = s_h(t)x, \tag{5.12}$$

$$\tau = \tau_h.$$

and s_h, τ_h are defined by (3.4), (3.5) respectively.

THEOREM 5.4. Suppose that the condition B and the regularity condition \tilde{R}_1 are satisfied. Then the following estimates hold for $r \geq 3$,

$$\| |s - s_h| \|_{L^\infty(0, T_0)} = O(h^{r+1}) \tag{5.13}$$

$$\| |\tau - \tau_h| \|_{L^\infty(0, T_0)} = O(h^{r+1}) \tag{5.14}$$

and

$$\| |U - U^h| \|_{L^\infty(0, T_0; H^j(\tilde{\Omega}(\tau)))} = O(h^{r+1-j}), \quad j = 0, 1, 2, \tag{5.15}$$

where $\| |\cdot| \|$ is interpreted as

$$\| |\phi| \|_{L^\infty(0, T_0; H^j(\tilde{\Omega}(\tau)))} = \int_0^{T_0} \| |\phi| \|_{H^j(\tilde{\Omega}(\tau))} dt$$

with $\tilde{\Omega}(\tau) = (0, \min(S(\tau), S_h(\tau)))$.

PROOF. From (2.10) and (3.4), we have

$$|s - s_h| \leq \int_0^t (|\eta_x(1)| + |\zeta_x(1)|) |s| dt' + \int_0^t |u_x^h(1)| |s - s_h| dt'.$$

An application of Gronwall's inequality and the estimates (4.19), (5.3), for $m = r + 1$ gives

$$\begin{aligned} \| |s - s_h| \|_{L^\infty(0, T)} &\leq K(K_2) \{ \| |\eta_x(1)| \|_{L^2(0, T)} + \| |\zeta_x(1)| \|_{L^2(0, T)} \} \\ &\leq K(K_2, K_8) \{ h^{2(r-1)} \| |u| \|_{L^2(H^{r+1})} + \| |\zeta_{xx}| \|_{L^2(L^2)} \} \\ &\leq K_{13} h^{r+1}, \quad \text{for } r \geq 3. \end{aligned} \tag{5.16}$$

The estimate (5.13) is immediate from (5.16), if we note that

$$\|s - s_h\|_{L^\infty(0, T_0)} = \|s - s_h\|_{L^\infty(0, T)}.$$

Further, the estimate (5.14) follows from (2.11), (3.5) and (5.16). Finally since

$$\| |u - u^h| \|_{L^\infty(0, T_0; H^j(\tilde{\Omega}(\tau)))} \leq \| |u - u^h| \|_{L^\infty(0, T; H^j(I))}$$

we obtain the required estimate (5.15).

6. GLOBAL EXISTENCE AND UNIQUENESS OF THE GALERKIN APPROXIMATION.

Now we consider the problem of existence of the Galerkin approximation u^h in the domain of existence of u . Towards this, let us recall (5.4) and note

$$\begin{aligned} A_p(u^h; u^h, \chi) - A_p(u; \tilde{u}, \chi) &= A_p(u; \zeta, \chi) + ([a(u^h) - a(u)] u_{xx}^h, \chi_{xx}) \\ &\quad + \eta_x(1)(xu_x^h, \chi_{xx}) - \zeta_x(1)(xu_x^h, \chi_{xx}). \end{aligned}$$

From the above, we get

$$\begin{aligned} (\zeta_{tx}, \chi_x) + A_p(u; \zeta, \chi) &= (\eta_{tx}, \chi_x) - \rho(\eta_x, \zeta_x) + \rho(\zeta_x, \chi_x) + ([a(u) - a(u^h)] \\ &\quad u_{xx}^h, \chi_{xx}) - \eta_x(1)(xu_x^h, \chi_{xx}) + \zeta_x(1)(xu_x^h, \chi_{xx}). \end{aligned} \tag{6.1}$$

But,

$$a(u) - a(u^h) = \tilde{a}_u e = - \int_0^1 \frac{\partial a}{\partial u}(u - \xi e) e \, d\xi. \tag{6.2}$$

Replacing u^h by $u - e$ in (6.1) with (6.2), we have

$$\begin{aligned} (\zeta_{tx}, \chi_x) + A_p(u; \zeta, \chi) &= -(\eta_t, \chi_{xx}) + \rho(\eta, \chi_{xx}) + \rho(\zeta_x, \chi_x) \\ &\quad - \left(\int_0^1 \frac{\partial a}{\partial u}(u - \xi e) (\eta - \zeta) d\xi (u_{xx} - e_{xx}), \chi_{xx} \right) - \eta_x(1)(x(u_x - e_x), \chi_{xx}) \\ &\quad + \zeta_x(1)(x(u_x - e_x), \chi_{xx}). \end{aligned}$$

Substitute e by $E(x, t)$, where $E \in H^2_0$. Then we get

$$\begin{aligned} (\zeta_{tx}, \chi_x) + A_p(u; \zeta, \chi) &= -(\eta_t, \chi_{xx}) + \rho(\eta, \chi_{xx}) + \rho(\zeta_x, \chi_x) \\ &\quad - \left(\int_0^1 \frac{\partial a}{\partial u}(u - \xi E) (\eta - \zeta) d\xi (u_{xx} - E_{xx}), \chi_{xx} \right) - \eta_x(1)(x(u_x - E_x), \chi_{xx}) \\ &\quad + \zeta_x(1)(x(u_x - E_x), \chi_{xx}), \end{aligned} \tag{6.3}$$

which is a linear ordinary differential equation in ζ . Therefore, for any $E = E(x,t)$ there exists a unique solution ζ of (6.3) with

$$\zeta(x, 0) = 0 \tag{6.4}$$

in the interval $(0,T]$.

The equation (6.3) defines an operator \mathcal{J} such that $\zeta = \mathcal{J}(E)$, for each $E \in H^2$. Since $e = \eta - \zeta$, therefore

$$e = \eta - \mathcal{J}(E), \text{ for each } E \in H^2. \tag{6.5}$$

To show the existence of a solution u^h of (3.2), we need to show that the operator equation (6.5) has a fixed point. In other words, we are looking for an $e(E)$ such that

$$e(E) = E$$

THEOREM 6.1. Suppose that the finite dimensional space S_h^0 satisfies inverse property and u is the unique solution of (2.7)-(2.9). Further, let the regularity conditions R_1 be satisfied. Then for some $\delta > 0$, there exists a solution $u^h \in S_h^0$ of (3.2) satisfying $\|u - u^h\|_{L^\infty(0,T_0; H^2(I))} \leq \delta$.

PROOF. Set $\chi = \zeta$ in (6.3) to get

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\zeta_x\|^2 + \tilde{\alpha} \|\zeta_{xx}\|^2 \\ & \leq \{ \|\eta_t\| + \rho \|\eta\| + K_1(\|\eta\|_{L^\infty} + \|\zeta\|_{L^\infty}) (\|u_{xx}\| + \|E_{xx}\|) \\ & \quad + |\eta_x(1)| (\|u_x\| + \|E_x\|) \} \|\zeta_{xx}\| + \rho \|\zeta_x\|^2 \\ & \quad + (\|u_x\| + \|E_x\|) |\zeta_x(1)| \|\zeta_{xx}\|. \end{aligned}$$

Using $|\zeta_x(1)| \leq \|\zeta_x\|^{1/2} \|\zeta_{xx}\|^{1/2}$, applying Young's inequality for the last term and $ab \leq \frac{a^2}{2\epsilon} + \frac{\epsilon}{2} b^2$, $a, b \geq 0$; $\epsilon > 0$ for the remaining terms, we get

$$\begin{aligned} & \frac{d}{dt} \|\zeta_x\|^2 + 2\tilde{\alpha} \|\zeta_{xx}\|^2 \\ & \leq K_{14}(\epsilon) \|\zeta_{xx}\|^2 + K(K_1, K_2, \rho; \epsilon) \{ \|\eta_t\|^2 + \|\eta\|_{L^\infty}^2 \\ & \quad + |\eta_x(1)| (1 + \|E\|_2^2) + K(K_1, K_2, \rho; \epsilon) (1 + \|E\|_2^2) \} \|\zeta_x\|^2. \end{aligned}$$

Choosing ϵ appropriately so that $2\tilde{\alpha} = K_{14}(\epsilon)$, integrating with respect to 't' and there after applying Gronwall's inequality, we get

$$\begin{aligned} \|\zeta\|_1^2(t) &\leq K(K_1, K_2; \rho) \exp[K(\rho, K_1; K_2)t(1 + \|E\|_{L^\infty(H^2)}^2)] \int_0^t \{\|\eta_\tau\|^2 \\ &+ (\|\eta\|_1^2 + |\eta_x(1)|^2)(1 + \|E\|_{L^\infty(H^2)}^2)\}. \end{aligned}$$

From the estimates (4.16), (4.17) and (4.19), it follows that

$$\|\zeta\|_{L^\infty(H^1)} \leq K_{15} \{h^{r+1} + (h^r + h^{2(r-1)})(1 + \|E\|_{L^\infty(H^2)}^2)\}, \tag{6.6}$$

where $K_{15} = K_{15}(K_1, K_2, K_4, K_8, \rho$ and $\|E\|_{L^\infty(H^2)})$. Thus we have

$$\begin{aligned} \|e\|_{L^\infty(H^2)} &\leq \|\eta\|_{L^\infty(H^2)} + \|\zeta\|_{L^\infty(H^2)} \\ &\leq \|\eta\|_{L^\infty(H^2)} + K_0 h^{-1} \|\zeta\|_{L^\infty(H^1)}. \end{aligned} \tag{6.7}$$

For $\|E\|_{L^\infty(H^2)} \leq \delta$ and from (4.16), (6.6), (6.7), we get

$$\|e\|_{L^\infty(H^2)} \leq K_{16} h^{r-1}, \text{ where } K_{16} = K_{16}(K_{15}, K_4, K_0; \delta).$$

Therefore, for sufficiently small h

$$\|e\|_{L^\infty(H^2)} \leq \delta.$$

Now, an application of Schauder's fixed point theorem guarantees the existence of an E such that $e = E$, which is a solution of the operator equation (6.5). The uniqueness of the approximate solution u^h is easy to prove. So we formalize the above in the form of a Theorem.

THEOREM 6.2. Let all the hypotheses of the Theorem 6.1 be satisfied and let $K > 0$. Then there exists one and only one solution $u^h \in S_h^0$ of (3.2) in the ball

$$\{\|u - u^h\|_{L^\infty(H^2)} \leq K\}, \text{ for sufficiently small } h \text{ and } r \geq 3.$$

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