# Homogeneous Lorentzian 3-manifolds with a parallel null vector field 

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#### Abstract

Lorentzian three-manifolds admitting a parallel null vector field have been intensively studied, since they possess several interesting geometrical properties which do not have a Riemannian counterpart. We completely classify homogeneous structures on Lorentzian three-manifolds which admit a parallel null vector field. This leads to the full classification of locally homogeneous examples within this class of manifolds.


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## 1 Introduction

The existence of parallel vector fields has strong and interesting consequences on the geometry of a manifold. If a Riemannian manifold $(M, g)$ admits such a vector field, then $(M, g)$ is locally reducible. The same property remains true for a pseudo-Riemannian manifold admitting a parallel non-null vector field. However, in the pseudo-Riemannian framework, a peculiar phenomenon arises: it can exist a parallel null vector field.

The geometry of a Lorentzian three-manifolds admitting a parallel null vector field has been studied in the fundamental paper [3]. These manifolds possess several interesting geometrical properties which do not have analogues in Riemannian settings. They are described in terms of a suitable system of local coordinates $(t, x, y)$ and form a quite large class, depending on an arbitrary two-variables function $f(x, y)$. The Levi-Civita connection and curvature of $\left(M, g_{f}\right)$ are completely described and several geometric consequences are deduced. In particular, locally symmetric examples are classified in terms of the defining function $f$, as well as curvature homogeneous examples with diagonal Ricci operator. It is then natural to ask the following

QUESTION: When $\left(M, g_{f}\right)$ is locally homogeneous?
The purpose of this paper is to answer the question above. It is worthwhile to remark that, differently from the Riemannian case, scalar curvature invariants do not help

[^0]here to determine the locally homogeneous examples. In fact, all scalar curvature invariants of a Lorentzian three-space ( $M, g_{f}$ ) admitting a parallel null vector field vanish identically ([3], p.844).

In order to determine the locally homogeneous spaces of the form $\left(M, g_{f}\right)$, we shall make use of the notion of homogeneous structure. Homogeneous pseudo-Riemannian structures were introduced by Gadea and Oubiña in [5], to obtain a characterization of reductive homogeneous pseudo-Riemannian manifolds, similar to the one given for homogeneous Riemannian manifolds by Ambrose and Singer [1] (see also [8]). In dimension three, as a consequence of the characterization proved by the second author in [2], a homogeneous Lorentzian space is necessarily reductive. Hence, the existence of a homogeneous structure on a Lorentzian three-space is a necessary and sufficient condition for local homogeneity. Recent complementary results about Lorentzian manifolds can be found in [6], [7].

In Section 2, we shall recall the description of Lorentzian three-manifolds $\left(M, g_{f}\right)$ admitting a parallel null vector field and the definition and basic properties of homogeneous (pseudo-Riemannian) structures. Then, in Section 3 we shall write down and solve the system of partial differential equations determining a homogeneous structure on $\left(M, g_{f}\right)$ in terms of its local components and we shall give the complete classification of locally homogeneous spaces $\left(M, g_{f}\right)$.

## 2 Preliminaries

We start with a short description of Lorentzian three-manifolds admitting a parallel null vector field, referring to [3] for more details and further results. Such a manifold $(M, g)$ admits a system of canonical local coordinates $(t, x, y)$, adapted to a parallel plane field in such a way that $\frac{\partial}{\partial_{t}}$ is the parallel null vector field, and there exists a smooth function $f=f(x, y)$, such that the Lorentzian metric tensor $g=g_{f}$ is described in local coordinates as follows:

$$
g=\left(\begin{array}{lll}
0 & 0 & 1  \tag{2.1}\\
0 & \varepsilon & 0 \\
1 & 0 & f
\end{array}\right)
$$

where $\varepsilon= \pm 1$. In the sequel, we shall denote by $\left(M, g_{f}\right)$ this Lorentzian manifold. In [3], a general description was provided for the wider class of Lorentzian three-manifolds admitting a parallel degenerate line field, that is, having a null vector field $u$ such that $\nabla u=\omega \otimes u$ for a suitable 1-form $\omega$. In this more general case, the function $f$ occurring in (2.1) also depends on $t$. Restricting ourselves to the case when $f=f(x, y)$, we have that the Levi-Civita connection $\nabla$ of $\left(M, g_{f}\right)$ is completely determined by

$$
\begin{equation*}
\nabla_{\partial_{x}} \partial_{y}=\frac{1}{2} f_{x} \partial_{t}, \quad \nabla_{\partial_{y}} \partial_{y}=\frac{1}{2} f_{y} \partial_{t}-\frac{\varepsilon}{2} f_{x} \partial_{x} \tag{2.2}
\end{equation*}
$$

where we put $\partial_{t}=\frac{\partial}{\partial_{t}}, \partial_{x}=\frac{\partial}{\partial_{x}}$ and $\partial_{y}=\frac{\partial}{\partial_{y}}$.
Next, the only non-vanishing local components of the curvature tensor are described by

$$
\begin{equation*}
R\left(\partial_{x}, \partial_{y}\right) \partial_{x}=-\frac{1}{2} f_{x x} \partial_{t}, \quad R\left(\partial_{x}, \partial_{y}\right) \partial_{y}=\frac{\varepsilon}{2} f_{x x} \partial_{x} \tag{2.3}
\end{equation*}
$$

Since $M$ is three-dimensional, its curvature is completely determined by the Ricci tensor $\varrho$. By (2.3) it easily follows that the local components of the Ricci tensor are given by

$$
\varrho=\left(\begin{array}{ccc}
0 & 0 & 0  \tag{2.4}\\
0 & 0 & 0 \\
0 & 0 & -\frac{\varepsilon}{2} f_{x x}
\end{array}\right) .
$$

In particular, (2.4) yields that $\left(M, g_{f}\right)$ is flat if and only if $f_{x x}=0$.
We now recall the definition and basic properties of homogeneous pseudo-Riemannian structures.

Definition 2.1. [5] Let $(M, g)$ be a pseudo-Riemannian manifold. A homogeneous pseudo-Riemannian structure is a tensor field $T$ of type $(1,2)$ on $M$, such that the connection $\tilde{\nabla}=\nabla-T$ satisfies

$$
\tilde{\nabla} g=0, \quad \tilde{\nabla} R=0, \quad \tilde{\nabla} T=0
$$

More explicitly, $T$ is the solution of the following system of equations (known as Ambrose-Singer equations):

$$
\begin{align*}
& g\left(T_{X} Y, Z\right)+g\left(Y, T_{X} Z\right)=0  \tag{2.5}\\
& \left(\nabla_{X} R\right)_{Y Z}=\left[T_{X}, R_{Y Z}\right]-R_{T_{X} Y Z}-R_{Y T_{X} Z}  \tag{2.6}\\
& \left(\nabla_{X} T\right)_{Y}=\left[T_{X}, T_{Y}\right]-T_{T_{X} Y} \tag{2.7}
\end{align*}
$$

for all vector fields $X, Y, Z$. The geometric meaning of the existence of a homogeneous pseudo-Riemannian structure is explained by the following

Theorem 2.2. [5] A connected, simply connected and complete pseudo-Riemannian manifold $(M, g)$ admits a homogeneous pseudo-Riemannian structure if and only if it is a reductive homogeneous pseudo-Riemannian manifold.

If at least one of the hypotheses of connectedness, simply connectedness or completeness is lacking in Theorem 2.2, then the existence of a solution of (2.5)-(2.7) implies that $(M, g)$ is locally isometric to a reductive homogeneous space (and so, locally homogeneous).

A three-dimensional homogeneous Lorentzian manifold is necessarily reductive. This was proved in [4] and also follows independently from the classification the second author gave in [2]. Therefore, the existence of a homogeneous structure on a Lorentzian three-manifold is a necessary and sufficient condition for its local homogeneity.

## 3 Homogeneous structures on ( $M, g_{f}$ )

In the sequel, we shall assume $\left(M, g_{f}\right)$ non-flat, that is, $f_{x x} \neq 0$. A homogeneous Lorentzian structure $T$ on $\left(M, g_{f}\right)$ is uniquely determined by its local components $T_{i j}^{k}$
with respect to the coordinate vector fields $\left\{\partial_{1}=\partial_{t}, \partial_{2}=\partial_{x}, \partial_{3}=\partial_{y}\right\}$. Functions $T_{i j}^{k}$ are defined by

$$
T\left(\partial_{i}, \partial_{j}\right)=\sum_{k=1}^{3} T_{i j}^{k} \partial_{k}
$$

for all indices $i, j, k$. Writing the Ambrose-Singer equations (2.5)-(2.7) for the coordinate vector fields $\partial_{i}$, we obtain the equivalent system of equations

$$
\begin{align*}
& T_{i j}^{r} g_{r k}+T_{i k}^{r} g_{r j}=0  \tag{3.1}\\
& \nabla_{i} \varrho_{j k}=-T_{i j}^{r} \varrho_{r k}-T_{i k}^{s} \varrho_{j s},  \tag{3.2}\\
& \left(\nabla_{i} T\right)_{\partial_{j}}=T_{\partial_{i}} T_{\partial_{j}}-T_{\partial_{j}} T_{\partial_{i}}-T_{T_{\partial_{i}} \partial_{j}}
\end{align*}
$$

for all indices $i, j, k$. Note that in (3.2) we took into account the fact that the curvature is completely determined by the Ricci tensor $\varrho$. Using (2.1), (2.2) and (2.4), from (3.1) and (3.2) we get

$$
\left\{\begin{array}{l}
T_{\partial_{t}} \partial_{t}=0,  \tag{3.4}\\
T_{\partial_{i}} \partial_{t}=-T_{i 3}^{3} \partial_{t}, \quad i=2,3 \\
T_{\partial_{i}} \partial_{x}=-\varepsilon T_{i 3}^{2} \partial_{t}, \quad i=1,2,3 \\
T_{\partial_{t}} \partial_{y}=T_{13}^{2} \partial_{x}, \\
T_{\partial_{i}} \partial_{y}=-f T_{i 3}^{3} \partial_{t}+T_{i 3}^{2} \partial_{x}+T_{i 3}^{3} \partial_{y}, \quad i=2,3, \\
T_{23}^{3}=-\frac{f_{x x x}}{2 f_{x x}}, \\
T_{33}^{3}=-\frac{f_{x x y}}{2 f_{x x}} .
\end{array}\right.
$$

From (3.3), we obtain the following system

$$
\left\{\begin{array}{l}
\partial_{t}\left(T_{13}^{2}\right)=\partial_{x}\left(T_{13}^{2}\right)=\partial_{y}\left(T_{13}^{2}\right)=0,  \tag{3.5}\\
\partial_{t}\left(T_{23}^{3}\right)=\partial_{x}\left(T_{23}^{3}\right)=\partial_{y}\left(T_{23}^{3}\right)=0, \\
\partial_{t}\left(T_{23}^{2}\right)=T_{13}^{2}\left(T_{23}^{3}+\varepsilon T_{13}^{2}\right), \\
\partial_{x}\left(T_{23}^{2}\right)=\varepsilon T_{23}^{2} T_{13}^{2}, \\
\partial_{y}\left(T_{23}^{2}\right)-\frac{f_{x}}{2} T_{13}^{2}-\varepsilon \frac{f_{x}}{2} T_{23}^{3}=T_{23}^{3} T_{33}^{2}-T_{33}^{3} T_{23}^{2}+\varepsilon T_{33}^{2} T_{13}^{2}, \\
\partial_{t}\left(T_{33}^{2}\right)=T_{13}^{2}\left(T_{33}^{3}-T_{23}^{2}\right), \\
\partial_{x}\left(T_{33}^{2}\right)-\frac{f_{x}}{2} T_{13}^{2}=T_{33}^{3} T_{23}^{2}-2 T_{23}^{3} T_{33}^{2}+f T_{23}^{3} T_{13}^{2}-\left(T_{23}^{2}\right)^{2}, \\
\partial_{y}\left(T_{33}^{2}\right)-\frac{f_{y}}{2} T_{13}^{2}+\varepsilon \frac{f_{x}}{2}\left(T_{23}^{2}-T_{33}^{3}\right)=f T_{13}^{2} T_{33}^{3}-T_{33}^{2} T_{23}^{2}-T_{33}^{3} T_{33}^{2}, \\
\partial_{t}\left(T_{33}^{3}\right)=-T_{13}^{2} T_{23}^{3}, \\
\partial_{x}\left(T_{33}^{3}\right)=-T_{23}^{3}\left(T_{23}^{2}+T_{33}^{3}\right), \\
\partial_{y}\left(T_{33}^{3}\right)+\varepsilon \frac{f_{x}}{2} T_{23}^{3}=-T_{33}^{2} T_{23}^{3}-\left(T_{33}^{3}\right)^{2}
\end{array}\right.
$$

We note that $T_{13}^{2}$ and $T_{23}^{3}$ are real constants; hence, by the sixth equation in (3.4), $\frac{f_{x x x}}{f_{x x}}$ must be constant. Now, put

$$
\begin{equation*}
T_{13}^{2}=\alpha, \quad T_{23}^{3}=\beta=-\frac{f_{x x x}}{2 f_{x x}}, \quad T_{23}^{2}=U, \quad T_{33}^{2}=V, \quad T_{33}^{3}=W=-\frac{f_{x x y}}{2 f_{x x}} \tag{3.6}
\end{equation*}
$$

where $\alpha, \beta$ are real constants and $U, V, W$ are real-valued smooth functions on $M$. Then, (3.5) becomes

$$
\left\{\begin{array}{l}
\partial_{t} U=\alpha(\beta+\varepsilon \alpha)  \tag{3.7}\\
\partial_{x} U=\varepsilon \alpha U \\
\partial_{y} U=\frac{f_{x}}{2}(\alpha+\varepsilon \beta)+V(\varepsilon \alpha+\beta)-U W \\
\partial_{t} V=\alpha(W-U) \\
\partial_{x} V=\alpha \frac{f_{x}}{2}+U(W-U)+\beta(\alpha f-2 V) \\
\partial_{y} V=\varepsilon \frac{f_{x}}{2}(W-U)+\alpha \frac{f_{y}}{2}+(\alpha f-V) W-U V \\
\partial_{t} W=-\alpha \beta \\
\partial_{x} W=-\beta(U+W) \\
\partial_{y} W=-\varepsilon \beta \frac{f_{x}}{2}-\beta V-W^{2}
\end{array}\right.
$$

Since $W$ depends only on $x$ and $y$, it follows from (3.7) that $\alpha \beta=0$, and (3.7) becomes

$$
\left\{\begin{array}{l}
\partial_{t} U=\varepsilon \alpha^{2}  \tag{3.8}\\
\partial_{x} U=\varepsilon \alpha U \\
\partial_{y} U=\frac{f_{x}}{2}(\alpha+\varepsilon \beta)+V(\varepsilon \alpha+\beta)-U W \\
\partial_{t} V=\alpha(W-U), \\
\partial_{x} V=\alpha \frac{f_{x}}{2}+U(W-U)-2 \beta V \\
\partial_{y} V=\varepsilon \frac{f_{x}}{2}(W-U)+\alpha \frac{f_{y}}{2}+(\alpha f-V) W-U V \\
\partial_{t} W=0 \\
\partial_{x} W=-\beta(U+W) \\
\partial_{y} W=-\varepsilon \beta \frac{f_{x}}{2}-\beta V-W^{2}
\end{array}\right.
$$

Next, we derive the first and the second equation of (3.8) for $x$ and $t$, respectively, obtaining

$$
\partial_{x} \partial_{t} U=0, \quad \partial_{t} \partial_{x} U=\alpha^{3}
$$

Therefore, $\alpha=0$ and (3.8) reduces to

$$
\left\{\begin{array}{l}
\partial_{t} U=0,  \tag{3.9}\\
\partial_{x} U=0 \\
\partial_{y} U=\frac{\varepsilon \beta}{2} f_{x}+\beta V-U W, \\
\partial_{t} V=0, \\
\partial_{x} V=U(W-U)-2 \beta V \\
\partial_{y} V=\frac{\varepsilon}{2} f_{x}(W-U)-V(W+U), \\
\partial_{t} W=0 \\
\partial_{x} W=-\beta(U+W), \\
\partial_{y} W=-\frac{\varepsilon \beta}{2} f_{x}-\beta V-W^{2}
\end{array}\right.
$$

We now calculate $\partial_{x} \partial_{y} U$ by the third equation of (3.9) and, taking into account the second equation and the expression of $\partial_{x} V$ and $\partial_{x} W$, we get

$$
\begin{equation*}
\beta\left(\frac{\varepsilon}{2} f_{x x}+2 U W-2 \beta V\right)=0 \tag{3.10}
\end{equation*}
$$

Moreover, deriving the last two equations of (3.9) respect to $y$ and $x$, respectively and using the expression of $\partial_{y} U, \partial_{x} V, \partial_{x} W$ and $\partial_{y} W$, by the identity $\partial_{x} \partial_{y} W=\partial_{y} \partial_{x} W$ we get

$$
\begin{equation*}
\beta\left(-\frac{\varepsilon}{2} f_{x x}+U^{2}+2 \beta V+W^{2}\right)=0 \tag{3.11}
\end{equation*}
$$

Summing up (3.10) and (3.11), then we obtain

$$
\beta(U+W)^{2}=0 .
$$

This leads to distinguish two cases:
First Case: $\beta=0$.
In this case, $f_{x x}$ only depends on $y$ and the system (3.9) reduces to

$$
\left\{\begin{array}{lll}
\partial_{t} U=0, & \partial_{x} U=0, & \partial_{y} U=-U W  \tag{3.12}\\
\partial_{t} V=0, & \partial_{x} V=U(W-U), & \partial_{y} V=\frac{\varepsilon}{2} f_{x}(W-U)-V(W+U), \\
\partial_{t} W=0, & \partial_{x} W=0, & \partial_{y} W=-W^{2}
\end{array}\right.
$$

By the last equation of (3.12), we can distinguish two subcases:
I a): If $W=0$ : we have that $f_{x x}$ is constant. Thus,

$$
\begin{equation*}
f(x, y)=\frac{\theta}{2} x^{2}+F(y) x+G(y) \tag{3.13}
\end{equation*}
$$

where $\theta \neq 0$ is a real constant and $F, G$ are smooth functions. Moreover, $U=u_{o}$, with $u_{o}$ a real constant, and the system (3.12) becomes

$$
\left\{\begin{array}{l}
\partial_{x} V=-u_{o}^{2},  \tag{3.14}\\
\partial_{y} V=-\frac{\varepsilon u_{o}}{2} f_{x}-u_{o} V
\end{array}\right.
$$

Derive the first equation of (3.14) with respect to $y$ and the second respect to $x$. Since $\partial_{y} \partial_{x} V=\partial_{y} \partial_{y} V$ we get

$$
u_{o}\left(u_{o}^{2}-\frac{\varepsilon}{2} f_{x x}\right)=0
$$

If $u_{o}=0$, we obtain $V=v_{o}$, with $v_{o}$ a real constant, $U=W=0$ and $f(x, y)$ determined by (3.13).
If $u_{o} \neq 0$, then $f_{x x}=2 \varepsilon u_{o}^{2}$ and so, in (3.13) we have $\theta=2 \varepsilon u_{o}^{2}$. The system (3.14) implies that $V(x, y)=-u_{o}^{2} x+H(y)$, with

$$
\begin{equation*}
H^{\prime}(y)+u_{o} H(y)+\frac{\varepsilon u_{o}}{2} F(y)=0 \tag{3.15}
\end{equation*}
$$

where $H$ is a smooth function.
I b): If $W \neq 0$, the last equation of (3.12) gives

$$
\begin{equation*}
W(y)=\frac{1}{y+k}, \tag{3.16}
\end{equation*}
$$

where $k$ is a real constant. On the other hand, $W=-\frac{f_{x x y}}{2 f_{x x}}$, thus

$$
\begin{equation*}
f(x, y)=\frac{s}{2(y+k)^{2}} x^{2}+P(y) x+Q(y) \tag{3.17}
\end{equation*}
$$

with $s \neq 0$ a real constant (since $f_{x x} \neq 0$ ). Next, from the third equation of (3.12) and by (3.16), if $U \neq 0$, we deduce

$$
\begin{equation*}
U=\frac{r}{y+k} \tag{3.18}
\end{equation*}
$$

for a real constant $r \neq 0$. On the other hand, the case $U=0$ can not occur. In fact, if $U=0$ then by (3.12) we get

$$
\left\{\begin{aligned}
\partial_{t} V & =0 \\
\partial_{x} V & =0 \\
\partial_{y} V & =\left(\frac{\varepsilon}{2} f_{x}-V\right) W
\end{aligned}\right.
$$

Deriving the last equation of the above system by $x$ and using the fact that $V=V(y)$ and $W=W(y)$, we obtain $f_{x x} W=0$, which is a contradiction, since $(M, g)$ is not flat and $W \neq 0$. Now, from (3.12), we calculate

$$
\begin{aligned}
\partial_{y} \partial_{x} V & =2 U W(U-W) \\
\partial_{x} \partial_{y} V & =(U-W)\left(U(U+W)-\frac{\varepsilon}{2} f_{x x}\right)
\end{aligned}
$$

which, together with (3.16), (3.17) and (3.18), imply

$$
(r-1)\left(2 r^{2}-2 r-\varepsilon s\right)=0
$$

So, if $r=1$ then $U(y)=W(y)=\frac{1}{y+k}$ and from (3.12) we easily get $V(y)=\frac{\sigma}{(y+k)^{2}}$, for a real constant $\sigma$, and $f$ is given by (3.17).
If $r=\frac{1 \pm \sqrt{1+2 \varepsilon s}}{2}$, with $1+2 \varepsilon s \geq 0$ and $s \neq 0$, because $W(y)=\frac{1}{y+k}$, (3.12) becomes

$$
\left\{\begin{array}{l}
\partial_{t} V=0  \tag{3.19}\\
\partial_{x} V=-\frac{\varepsilon s}{2(y+k)^{2}} \\
\partial_{y} V=\frac{\varepsilon(1 \mp \sqrt{1+2 \varepsilon s})}{4(y+k)} f_{x}-\frac{(3 \pm \sqrt{1+2 \varepsilon s})}{2(y+k)} V
\end{array}\right.
$$

The second equation in (3.19) gives

$$
\begin{equation*}
V(x, y)=-\frac{\varepsilon s}{2(y+k)^{2}} x+L(y) \tag{3.20}
\end{equation*}
$$

where $L=L(y)$ is a smooth function. Next, deriving the above expression of $V$ by $y$, we find

$$
\begin{equation*}
\partial_{y} V=\frac{\varepsilon s}{(y+k)^{3}} x+L^{\prime}(y) \tag{3.21}
\end{equation*}
$$

We compare (3.21) together with the third equation in (3.19). Taking into account the expression of $f_{x}$ deduced by (3.17), we obtain

$$
\begin{equation*}
L^{\prime}(y)+\frac{(3 \pm \sqrt{1+2 \varepsilon s})}{2(y+k)} L(y)=\frac{\varepsilon(1 \mp \sqrt{1+2 \varepsilon s})}{4(y+k)} P(y) \tag{3.22}
\end{equation*}
$$

Then, $V$ is given by (3.20), under the condition that (3.22) is satisfied.
Second Case: $\beta \neq 0$.
In this case $W=-U=-\frac{f_{x x y}}{2 f_{x x}}$ and the system (3.9) becomes

$$
\left\{\begin{array}{l}
\partial_{y} U=\frac{\varepsilon \beta}{2} f_{x}+\beta V+U^{2}  \tag{3.23}\\
\partial_{t} V=0 \\
\partial_{x} V=-2 U^{2}-2 \beta V \\
\partial_{y} V=-\varepsilon f_{x} U
\end{array}\right.
$$

Hence, $U$ only depends on $y$. The first equation of (3.23), derived with respect to the $x$, gives

$$
\begin{equation*}
\partial_{x} V=-\frac{\varepsilon}{2} f_{x x} \tag{3.24}
\end{equation*}
$$

We compare the above formula with the third equation of (3.23), getting

$$
V=\frac{1}{\beta}\left(\frac{\varepsilon}{4} f_{x x}-U^{2}\right)
$$

Now, taking into account the first and the third equation of (3.23), and making use of the (3.24), we get

$$
\begin{equation*}
2 U_{y}=\frac{\varepsilon}{2}\left(f_{x x}+2 \beta f_{x}\right) \tag{3.25}
\end{equation*}
$$

Since $\frac{f_{x x x}}{2 f_{x x}}=-\beta$, we can integrate with respect to the $x$ and we obtain

$$
f_{x x}= \pm e^{M(y)-2 \beta x}
$$

where $M=M(y)$ is a smooth function. Again integrating with respect to the $x$, we find

$$
\begin{equation*}
f_{x}(x, y)=\mp \frac{e^{M(y)-2 \beta x}}{2 \beta}+N(y) \tag{3.26}
\end{equation*}
$$

with $N$ a smooth function only depending on y. Taking into account (3.25), since $U=\frac{f_{x x y}}{2 f_{x x}}$, it holds:

$$
\begin{equation*}
M^{\prime \prime}(y)=\varepsilon \beta N(y) \tag{3.27}
\end{equation*}
$$

Thus, integrating again (3.26) with respect to the $x$, we obtain

$$
\begin{equation*}
f(x, y)= \pm \frac{e^{M(y)-2 \beta x}}{4 \beta^{2}}+\frac{\varepsilon}{\beta} M^{\prime \prime}(y) x+R(y) \tag{3.28}
\end{equation*}
$$

where $R=R(y)$ is a smooth function. Therefore, calculations above lead to the following

Theorem 3.1. Let $\left(M, g_{f}\right)$ be a non-flat Lorentzian three-space admitting a parallel null vector field. $\left(M, g_{f}\right)$ is locally homogeneous if and only for its defining function $f$, one of the following statements holds true:
(i) (locally symmetric case) there exist a real constant $\theta \neq 0$ and two one-variable smooth functions $F$ and $G$, such that

$$
f(x, y)=\frac{\theta}{2} x^{2}+F(y) x+G(y)
$$

(ii) there exist a real constant $s \neq 0$ and two one-variable smooth functions $P$ and $Q$, such that

$$
f(x, y)=\frac{s}{2(y+k)^{2}} x^{2}+P(y) x+Q(y)
$$

(iii) there exist two one-variable smooth functions $M$ and $R$, such that

$$
f(x, y)= \pm \frac{e^{M(y)-2 \beta x}}{4 \beta^{2}}+\frac{\varepsilon}{\beta} M^{\prime \prime}(y) x+R(y)
$$

Notice that in order to prove Theorem 3.1, we completely solved the system (3.1)(3.3). Thus, we also obtained the complete classification of homogeneous structures on $\left(M, g_{f}\right)$, summarized in the following

Proposition 3.2. Let $\left(M, g_{f}\right)$ be a non-flat locally homogeneous Lorentzian threespace admitting a parallel null vector field. All and the ones homogeneous structures on $\left(M, g_{f}\right)$ are determined (through their local components $T_{i j}^{k}$ with respect to the coordinate vector fields) by (3.6), where $\alpha=0, \beta$ is a real constant and $U, V, W$ are smooth real-valued functions for which one of the following statements holds true:
(i) If $f$ satisfies (3.13), then $\beta=0, W=0$ and

- either $U=0$ and $V=v_{0}$ is constant, or
- $U=u_{0} \neq 0$ is constant and $V(x, y)=-u_{o}^{2} x+H(y)$, for a smooth function $H$ satisfying (3.15). In this case, $\theta=2 \varepsilon u_{0}^{2}$ in (3.13).
(ii) If $f$ satisfies (3.17), then
- either there exist two real constants $k$ and $\sigma$, such that $U(y)=W(y)=\frac{1}{y+k}$ and $V(y)=\frac{\sigma}{(y+k)^{2}}$, or
- there exist two real constants $k$ and $s \neq 0$, such that $U=\frac{1 \pm \sqrt{1+2 \varepsilon s}}{2(y+k)}$, $W(y)=\frac{1}{y+k}$ and $V(x, y)=-\frac{\varepsilon s}{2(y+k)^{2}} x+L(y)$, for a smooth function $L$ satisfying (3.22).
(iii) If $f$ satisfies (3.28), then $\beta \neq 0, U(y)=-W(y)=\frac{M^{\prime}(y)}{2}$ and $V(x, y)=$ $\frac{1}{4 \beta}\left( \pm \varepsilon e^{M(y)-2 \beta x}-M^{\prime}(y)^{2}\right)$, for a smooth function $M$.

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