

Remarks on the relativistic self-dual Maxwell-Chern-Simons-Higgs system *

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Abstract

In this note we present some recent developments on the topological multi-vortex solutions of the self-dual Maxwell-Chern-Simons-Higgs system in \mathbb{R}^2 . We find that all the topological solutions are admissible in the sense defined in [2]. We also discover that the convergence of the topological solution to the solutions of the self-dual Chern-Simons equations can be improved to be strong.

1 Introduction

We are concerned with the semilinear elliptic system in \mathbb{R}^2 :

$$\Delta u = 2q^2(e^u - 1) - 2q\kappa A_0 + 4\pi \sum_{j=1}^m \delta(z - z_j) \quad (1)$$

$$\Delta A_0 = \kappa q(1 - e^u) + (\kappa^2 + 2q^2 e^u)A_0, \quad (2)$$

where $q > 0$ is the charge of electron, and $\kappa > 0$ is the Chern-Simons coupling constant. Each point of the prescribed set $\mathbf{Z} = \{z_1, \dots, z_m\}$ is called a vortex point. We consider two different physically-meaningful sets of boundary conditions related to finiteness of total energy ([5], [2]). One is the topological boundary condition,

$$\lim_{|z| \rightarrow \infty} u = 0, \quad \lim_{|z| \rightarrow \infty} A_0 = 0, \quad (3)$$

and the other is the nontopological boundary condition,

$$\lim_{|z| \rightarrow \infty} u = -\infty, \quad \lim_{|z| \rightarrow \infty} A_0 = -\frac{q}{\kappa}, \quad (4)$$

The system (1)-(2) arises from a Jaffe-Taubes[4] reduction of the Bogomolnyi equations of the self-dual Maxwell-Chern-Simons-Higgs system[5]. One major motivation for introducing the system was unification of two apparently separate

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systems, namely, the Abelian Higgs system and the Chern-Simons system. The Abelian Higgs system reduces to

$$\Delta u = 2q^2(e^u - 1) + 4\pi \sum_{j=1}^m \delta(z - z_j), \quad (5)$$

while the Chern-Simons system reduces to

$$\Delta u = 4l^2 e^u (e^u - 1) + 4\pi \sum_{j=1}^m \delta(z - z_j), \quad (6)$$

where l is a physical parameter related to the Chern-Simons constant. The Abelian Higgs system (5) has finite-energy solutions under only the topological boundary conditions (3), while the Chern-Simons system has finite-energy solutions under either boundary conditions (3) or (4). For the precise meaning of the finite energy of these systems, we refer to [4].

The mathematical analysis of the topological multi-vortex solutions for the Abelian Higgs system and the Chern-Simons system is studied in [4] and [7], [6]. We note that recently existence of non-topological multi-vortex solutions of (6) is established in [1]. The existence of topological multi-vortex solutions of (1)-(2) was established by a simple variational argument in [2]. Moreover, the existence of so-called admissible solutions was established by an iteration scheme in [2]. We recall the definition of admissible topological solution of (1)-(3).

Definition A topological solution pair (u, A_0) of (1)-(2) is called admissible if it satisfies one of the following inequalities.

- (i) $A_0 \leq 0$
- (ii) $u \leq 0$
- (iii) $A_0 \geq \frac{q}{\kappa}(e^u - 1)$
- (iv) $v \leq u_a^q$, where u_a^q is the solution the (topological) Abelian Higgs equation (5).

We remark that the conditions (i)-(iv) are shown to be equivalent to each other [2]. In [2] we also considered the two convergence problems of the *admissible* topological solutions of (1)-(2). One is the Abelian Higgs limit, namely the problem of identifying the behavior of the solution $u^{\kappa, q}$ of (1)-(2) as $\kappa \rightarrow 0$ with q kept fixed. The other is the Chern-Simons limit, the similar problem as both κ and q go to infinity with the ratio $l = q^2/\kappa$ kept fixed. In the Abelian Higgs limit we proved in [2] that admissible topological solution of (1)-(2) converges strongly to the solution of the Abelian Higgs equation, while in the Chern-Simons limit we could just show that our solution is weakly consistent to the Chern-Simons equation. For the precise statements of these results see [2]. The natural open questions raised were

1. Is any topological solution of (1)-(2), which is smooth except at the points z_1, \dots, z_m , admissible?
2. Can we strengthen the sense of convergence in the Chern-Simons limit problem?

Question 1 is concerned with the physical validity of the model system, (1)-(2), since $u \leq 0$ is equivalent to the condition $|\phi|^2 = e^u \leq 1$ for the Higgs field ϕ . Question 2, combined with the already established strong convergence in the Abelian Higgs limit, is concerned with the rigorous verification of the physical argument that the Maxwell-Chern-Simons-Higgs model is a unification of the Abelian Higgs model and the Chern-Simons model[5]. In [3] we answer these two questions in the affirmative. In the next section we state our results and their implications.

For further discussion, we introduce the background function u_0 defined by

$$u_0 = \sum_{j=1}^m \ln \left(\frac{|z - z_j|^2}{1 + |z - z_j|^2} \right),$$

and we set $u = v + u_0$ to remove the singular inhomogeneous term in (1). Then (1) and (2) become

$$\Delta v = 2q^2(e^{v+u_0} - 1) - 2q\kappa A_0 + g, \quad (7)$$

$$\Delta A_0 = \kappa q(1 - e^{v+u_0}) + (\kappa^2 + 2q^2 e^{v+u_0})A_0 \quad (8)$$

with the topological boundary condition

$$\lim_{|z| \rightarrow \infty} v = 0, \quad \lim_{|z| \rightarrow \infty} A_0 = 0, \quad (9)$$

where

$$g = \sum_{j=1}^m \frac{4}{(1 + |z - z_j|^2)^2}.$$

In this setting the Chern-Simons equation (6) becomes

$$\Delta v = 4l^2 e^{v+u_0} (e^{v+u_0} - 1) + g. \quad (10)$$

2 Main Results

The following result is proved in [3].

Theorem 1 *Suppose $\mathbf{Z} = \{z_1, \dots, z_m\} \subset \mathbb{R}^2$ is given as before. Then any topological solution (u, A_0) in $C^2(\mathbb{R}^2 \setminus \mathbf{Z})$ is admissible.*

One immediate consequence of Theorem 1 and the argument of the construction is that the solution constructed in Section 3 is maximal. On the other hand,

due to the monotonicity of the minimizing functional \mathcal{F} ,

$$\begin{aligned} \mathcal{F}(v) = & \int \left[\frac{1}{2} |\Delta v|^2 - (\Delta g - \kappa^2 g)v + 2q^4 (e^{v+u_0} - 1)^2 \right. \\ & \left. + \frac{1}{2} \kappa^2 |\nabla v|^2 + 2q^2 e^{v+u_0} |\nabla(v + u_0)|^2 \right] dx, \end{aligned} \quad (11)$$

as established in Section 4 of [2], the solution constructed in Section 2 of [2] by the variational method is minimal. Thus we have, as a result, constructed the maximal and the minimal solutions of the system. As remarked in [3], the strong convergence in the Chern-Simons limit for admissible solutions in the periodic boundary condition can be extended to the case of our solutions of the system (7)-(8). In particular, due to Theorem 1 we can remove the condition of admissibility, and obtain:

Theorem 2 *Let $(v^{\kappa,q}, A_0^{\kappa,q})$ be any topological solution of (7)-(8). Then $v^{\kappa,q} \rightarrow v_{cs}^l$, and $\frac{\kappa}{q} A_0^{\kappa,q} \rightarrow e^{v_{cs}^l + u_0} - 1$, both in $H^1(\mathbb{R}^2)$ as $\kappa \rightarrow \infty$ with $\frac{q^2}{\kappa} = l$ kept fixed, where v_{cs}^l denotes a topological solution of (10).*

We now have further open problems to consider for the system (1)- (2):

1. Prove uniqueness, or multiplicity of topological solution.
2. Prove existence of non-topological multi-vortex solutions.

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