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Strong maximum principle for a sublinear elliptic problem at resonance

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Abstract. We examine the semilinear resonant problem

$$-\Delta u = \lambda_1 u + \lambda g(u)$$
 in Ω , $u \ge 0$ in Ω , $u_{|\partial\Omega} = 0$,

where $\Omega \subset \mathbb{R}^N$ is a smooth, bounded domain, λ_1 is the first eigenvalue of $-\Delta$ in Ω , $\lambda > 0$. Inspired by a previous result in literature involving power-type nonlinearities, we consider here a generic sublinear term g and single out conditions to ensure: the existence of solutions for all $\lambda > 0$; the validity of the strong maximum principle for sufficiently small λ . The proof rests upon variational arguments.

Keywords: resonant problem, existence, maximum principle.

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1 Introduction

Let $\Omega \subset \mathbb{R}^N$, $N \geq 1$, be a bounded domain of class C^2 , and let λ_1 be the first eigenvalue of $-\Delta$ in Ω with Dirichlet boundary conditions. The issue of the existence of solutions of the problem

$$\begin{cases}
-\Delta u = \lambda_1 u + u^{s-1} - \mu u^{r-1} & \text{in } \Omega \\
u \ge 0 & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(1.1)

 $s \in (1,2)$, $r \in (1,s)$, and $\mu > 0$, has been the subject of study of the recent [3]. As a distinctive feature, the right-hand side term $f(t) := \lambda_1 t + t^{s-1} - \mu t^{r-1}$ in (1.1) is not locally Lipschitz near 0, and moreover satisfies the sign property

$$f^{-1}((-\infty,0]) \supseteq (0,a]$$
, for some $a > 0$.

As a result, from the celebrated paper [13] (see also [8]), it is known that the strong maximum principle may fail to be valid in this context. By adopting minimax and perturbation

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techniques, the author of [3] showed instead that such a principle does hold as long as the perturbation parameter is chosen sufficiently large. More precisely, the main results in [3] state that problem (1.1) has non-zero solutions for the entire positive range of μ ; positive solutions for μ large enough.

The fact that, after a rescaling, (1.1) can be turned into the problem

$$\begin{cases}
-\Delta u = \lambda_1 u + \lambda (u^{s-1} - u^{r-1}) & \text{in } \Omega \\
u \ge 0 & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(1.2)

for a suitable $\lambda > 0$, raises the natural question whether, as explicitly expressed in [3, Remark 2.4], the same results mentioned above continue to hold when the powers in (1.2) are replaced by a generic nonlinear term g. And, if it is so, it would be interesting of course to identify some "minimal" structure conditions on g for the validity of such results. In the present paper we address these questions and consider the problem

$$\begin{cases}
-\Delta u = \lambda_1 u + \lambda g(u) & \text{in } \Omega \\
u \ge 0 & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(P_{\lambda})

where $g:[0,+\infty)\to\mathbb{R}$ is continuous, g(0)=0, and obeys the following conditions:

$$(g_1)$$
 there exists $q \in (1,2)$ such that $k_1 := \sup_{t>0} \frac{|g(t)|}{1+t^{q-1}} < +\infty;$

$$(g_2) \lim_{t\to 0^+} \frac{g(t)}{t} = -\infty;$$

$$(g_3)$$
 $\liminf_{t\to+\infty} G(t) > 0$;

$$(g_4) \lim_{t\to +\infty} (g(t)t - 2G(t)) = -\infty,$$

where, as usual,

$$G(t) := \int_0^t g(s)ds$$
, for all $t \ge 0$.

Problems like (P_{λ}) are being investigated since Landesman and Lazer's pioneering work [9], in which sufficient conditions, based on the interaction between the nonlinearity and the spectrum of the linear operator, were given for them to have a solution. Noteworthy contributions following that work can be found in [2,5,12] and also in [6,7,10,11,14] (see the related references as well) in which several classes of elliptic problems at resonance are investigated via variational and topological methods.

Coming back to (P_{λ}) , our approach develops along the same line of reasoning as [3]. We prove initially that (P_{λ}) has at least a non-zero solution for all $\lambda > 0$. This is accomplished by considering a sequence of problems near resonance whose solutions are shown to converge to a solution of the original problem. In this regard, assumption (g_4) comes into play to prove the boundedness of the sequence of approximating solutions. Then, by exploiting the classical decomposition of $H_0^1(\Omega)$ into the first eigenspace and its orthogonal complement, we show

that, for sufficiently small λ , the set of solutions to (P_{λ}) is contained in the interior of the positive cone of $C_0^1(\overline{\Omega})$. It still remains an open question to investigate the uniqueness of positive solutions to (P_{λ}) (in the one-dimensional case and for power-nonlinearities it has instead been established in [4]), as well as the existence of non-zero solutions compactly supported in Ω , in the spirit of [8].

Our main results, Theorems 2.3 and 2.4, are stated and proved in the coming section. Before going on, we arrange some notation and the variational framework for (P_{λ}) . We set

$$||u|| := \left(\int_{\Omega} |\nabla u|^2 dx\right)^{\frac{1}{2}}$$
, for all $u \in H_0^1(\Omega)$,

and denote by $\|\cdot\|_p$, $p \in [1, +\infty]$, the classical L^p -norm on Ω . We also set

$$c_p := \sup_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\|u\|_p}{\|u\|}$$

for each $p \ge 1$, with $p \le \frac{2N}{N-2}$ if $N \ge 3$, and denote by ϕ_1 the positive eigenfunction associated with λ_1 and normalized with respect to $\|\cdot\|_{\infty}$. We recall that the first two eigenvalues λ_1, λ_2 of $-\Delta$ in Ω admit the variational characterization

$$\lambda_1 = \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\|u\|^2}{\|u\|_2^2}, \quad \lambda_2 = \inf_{u \in \operatorname{span}\{\phi_1\}^{\perp} \setminus \{0\}} \frac{\|u\|^2}{\|u\|_2^2}.$$

Given a set $E \subset \mathbb{R}^N$, its Lebesgue measure will be denoted by the symbol |E|. Throughout this paper, the symbols C, C_1, C_2, \ldots represent generic positive constants whose exact value may change from occurrence to occurrence.

For all $\lambda > 0$, we denote by $I_{\lambda} : H_0^1(\Omega) \to \mathbb{R}$ the energy functional associated with (P_{λ}) ,

$$I_{\lambda}(u) := \frac{1}{2} \|u\|^2 - \frac{\lambda_1}{2} \|u_+\|_2^2 - \lambda \int_{\Omega} G(u_+) dx$$
, for all $u \in H_0^1(\Omega)$,

where $u_+ = \max\{u, 0\}$. By a weak solution to (P_λ) we mean any $u \in C^0(\overline{\Omega}) \cap H^1_0(\Omega)$ verifying

$$\int_{\Omega} (\nabla u \nabla v - \lambda_1 u v - \lambda g(u) v) \, dx = 0, \quad \text{for all } v \in H^1_0(\Omega).$$

2 Results

As already mentioned, we start by considering a sequence of approximating problems.

Lemma 2.1. For each $\lambda > 0$, there exists $\bar{n} \in \mathbb{N}$ such that the problem

$$\begin{cases}
-\Delta u = \left(\lambda_1 - \frac{1}{n}\right) u + \lambda g(u) & \text{in } \Omega \\
u \ge 0 & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(P_n)

admits a non-zero weak solution u_n , with positive energy, for all $n \geq \bar{n}$.

Proof. Fix $\lambda > 0$ and let $n \in \mathbb{N}$ with $n > \frac{1}{\lambda_1}$. Let us first show that the energy functional $I_n : H_0^1(\Omega) \to \mathbb{R}$ corresponding to (P_n) ,

$$I_n(u) := I_{\lambda}(u) + \frac{1}{2n} \|u_+\|_2^2 = \frac{1}{2} \|u\|^2 - \frac{1}{2} \left(\lambda_1 - \frac{1}{n}\right) \|u_+\|_2^2 - \lambda \int_{\Omega} G(u_+) dx, \qquad (2.1)$$

for all $u \in H_0^1(\Omega)$, has the mountain pass geometry for sufficiently large $n \in \mathbb{N}$.

Fix $k \in (2, 2^*)$ and set

$$M := \frac{k}{2} \sup_{t>0} \frac{\lambda_1 t^2 + 2\lambda G(t)}{t^k}.$$

By (g_1) and (g_2) one has $0 < M < +\infty$ and $\frac{\lambda_1}{2}t^2 + \lambda G(t) \le \frac{M}{k}t^k$, for all $t \ge 0$. Then, defining

$$R:=(Mc_k^k)^{\frac{1}{2-k}},$$

we easily obtain

$$\inf_{u \in S_{R}} I_{n}(u) \geq \inf_{\|u\| = R} \left(\frac{1}{2} \|u\|^{2} - \frac{M}{k} \|u\|_{k}^{k} \right)
\geq \inf_{u \in S_{R}} \left(\frac{1}{2} \|u\|^{2} - \frac{Mc_{k}^{k}}{k} \|u\|^{k} \right)
= \left(\frac{1}{2} - \frac{1}{k} \right) R^{2} > 0,$$
(2.2)

for any $n \in \mathbb{N}$, where $S_R := \{ u \in H_0^1(\Omega) : ||u|| = R \}$.

Now, let us show that there exist $u_1 \in H_0^1(\Omega)$, with $||u_1|| > R$, and $\bar{n} \in \mathbb{N}$, such that $I_n(u_1) < 0$ for all $n \ge \bar{n}$. Owing to (g_3) , there exist L, b > 0 such that

$$G(t) > L$$
, for all $t > b$.

If we denote by

$$E_{\gamma} := \{ x \in \Omega : \phi_1(x) < \gamma \},$$

with $\gamma > 0$, then there exists $\gamma_1 > 0$ such that

$$L > \frac{k_1(bq + b^q)|E_{\gamma}|}{q(|\Omega| - |E_{\gamma}|)}, \quad \text{for all } \gamma \in (0, \gamma_1).$$
 (2.3)

Fix $\bar{\gamma} \in \mathbb{R}$ satisfying

$$0<\bar{\gamma}<\min\left\{\gamma_1,\frac{b}{R}\right\}.$$

Since the function $\psi(t) := q\bar{\gamma}t + \bar{\gamma}^qt^q$ is continuous in $(0, +\infty)$ and $\psi\left(\frac{b}{\bar{\gamma}}\right) = bq + b^q$, thanks to (2.3), there exists $\bar{t} > \frac{b}{\bar{\gamma}}$ such that

$$L > \frac{k_1(q\bar{\gamma}\bar{t} + \bar{\gamma}^q\bar{t}^q)|E_{\bar{\gamma}}|}{q(|\Omega| - |E_{\bar{\gamma}}|)}.$$
(2.4)

With the aid of (g_1) and (2.4) we then obtain

$$\begin{split} \int_{\Omega} G(\bar{t}\phi_1) dx &= \int_{E_{\bar{\gamma}}} G(\bar{t}\phi_1) dx + \int_{\{\phi_1 \geq \bar{\gamma}\}} G(\bar{t}\phi_1) dx \\ &\geq -k_1 \int_{E_{\bar{\gamma}}} \left(\bar{t}\phi_1 + \frac{(\bar{t}\phi_1)^q}{q} \right) dx + \int_{\{\phi_1 \geq \bar{\gamma}\}} G(\bar{t}\phi_1) dx \\ &\geq -k_1 \left(\bar{t}\bar{\gamma} + \frac{\bar{t}^q \bar{\gamma}^q}{q} \right) |E_{\bar{\gamma}}| + L(|\Omega| - |E_{\bar{\gamma}}|) \\ &> 0. \end{split}$$

As a result, there exists $\bar{n} \in \mathbb{N}$, with $\bar{n} > \frac{1}{\lambda_1}$, such that

$$I_n(\bar{t}\phi_1) = \frac{\bar{t}^2}{2n} \|\phi_1\|_2^2 - \lambda \int_{\Omega} G(\bar{t}\phi_1) dx < 0$$

for all $n \geq \bar{n}$. Therefore, the functional I_n satisfies the geometric conditions required by the mountain pass theorem for all $n \geq \bar{n}$.

Moreover, by (g_1) and Sobolev embeddings, one has

$$I_{n}(u) \geq \frac{1}{2n\lambda_{1}} \|u\|^{2} - \lambda k_{1} \left(\int_{\Omega} |u| dx + \frac{1}{q} \int_{\Omega} |u|^{q} dx \right)$$

$$\geq \frac{1}{2n\lambda_{1}} \|u\|^{2} - \lambda c_{1} k_{1} \|u\| - \frac{\lambda c_{q} k_{1}}{q} \|u\|^{q},$$

and thus $I_n(u) \to +\infty$ as $||u|| \to +\infty$. This fact, in addition to standard arguments (see for instance Example 38.25 of [15]), ensures that I_n satisfies the Palais–Smale condition. Then, by invoking the classical mountain pass theorem, I_n admits a critical point $u_n \in H_0^1(\Omega) \setminus \{0\}$ for all $n \ge \bar{n}$, and, by (2.2), one also has

$$I_n(u_n) = \inf_{\gamma \in \Gamma} \max_{t \in [0,1]} I_n(\gamma(t)) \ge \left(\frac{1}{2} - \frac{1}{k}\right) R^2, \tag{2.5}$$

where $\Gamma := \{ \gamma \in C^0([0,1], H_0^1(\Omega)) : \gamma(0) = 0, \ \gamma(1) = u_1 \}$. This concludes the proof.

Lemma 2.2. Let $\lambda > 0$, $\bar{n} \in \mathbb{N}$ and let u_n , with $n \geq \bar{n}$, be as in Lemma 2.1. Then, the sequence $\{u_n\}_{n\geq \bar{n}}$ is bounded in $H_0^1(\Omega)$.

Proof. Let $n \in \mathbb{N}$, $n \ge \overline{n}$. By standard regularity theory, $u_n \in C^{1,\alpha}(\overline{\Omega})$, for some $\alpha \in (0,1)$. For any $n \in \mathbb{N}$, $n \ge \overline{n}$ there exist, uniquely determined, $t_n \in \mathbb{R}$ and $w_n \in \text{span}\{\phi_1\}^{\perp}$ such that

$$u_n = t_n \phi_1 + w_n.$$

It is straightforward to verify that $w_n \in C^{1,\alpha}(\overline{\Omega})$ is a weak solution to

$$\begin{cases}
-\Delta u = \left(\lambda_1 - \frac{1}{n}\right) u + \lambda g(t_n \phi_1 + u) - \frac{t_n}{n} \phi_1 & \text{in } \Omega \\
u \ge 0 & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(2.6)

and therefore, also by (g_1) , one has

$$||w_{n}||^{2} \leq \left(\frac{\lambda_{1} - \frac{1}{n}}{\lambda_{2}}\right) ||w_{n}||^{2} + \lambda \int_{\Omega} g(t_{n}\phi_{1} + w_{n})w_{n}dx$$

$$\leq \left(\frac{\lambda_{1} - \frac{1}{n}}{\lambda_{2}}\right) ||w_{n}||^{2} + \lambda k_{1} ||w_{n}||_{1} + \lambda k_{1} t_{n}^{q-1} ||\phi_{1}||_{\infty}^{q-1} ||w_{n}||_{1} + \lambda k_{1} ||w_{n}||_{q}^{q}.$$

$$(2.7)$$

From (2.7), it follows that

$$||w_n|| \le C \left((1 + t_n^{q-1}) + ||w_n||^{q-1} \right),$$
 (2.8)

for some C>0. We claim that the sequence $\{t_n\}_{n\geq \bar{n}}$ is bounded in \mathbb{R} . Arguing by contradiction, assume that, up to a subsequence, $t_n\to +\infty$ as $n\to +\infty$. Without loss of generality, we can assume that $t_n\geq 1$ for all $n\geq \bar{n}$ and, since

$$y^{q-1} \le C_1 + \frac{1}{2C}y \le C_1t_n^{q-1} + \frac{1}{2C}y$$
, for all $y > 0$,

from (2.8) we deduce

$$||w_n|| \le 2Ct_n^{q-1} + C||w_n||^{q-1} \le 2Ct_n^{q-1} + CC_1t_n^{q-1} + \frac{1}{2}||w_n||,$$

and then

$$||w_n|| \le C_2 t_n^{q-1}.$$

Therefore, fixing $p > \max\left\{\frac{N}{2}, \frac{q}{q-1}\right\}$, we obtain

$$||w_{n}||_{\infty} \leq C_{3} \left(||w_{n}||_{p} + ||g(t_{n}\phi_{1} + w_{n})||_{p} + \frac{t_{n}}{n} ||\phi_{1}||_{p} \right)$$

$$\leq C_{4} \left(||w_{n}||_{\infty}^{\frac{p-1}{p}} ||w_{n}||_{1}^{\frac{1}{p}} + 1 + t_{n}^{q-1} + ||w_{n}||_{\infty}^{q-1-\frac{q}{p}} ||w_{n}||_{q}^{\frac{q}{p}} + \frac{t_{n}}{n} \right)$$

$$\leq C_{5} \left(||w_{n}||_{\infty}^{\frac{p-1}{p}} t_{n}^{\frac{q-1}{p}} + t_{n}^{q-1} + ||w_{n}||_{\infty}^{q-1-\frac{q}{p}} t_{n}^{\frac{q(q-1)}{p}} + \frac{t_{n}}{n} \right).$$

Dividing the first and the last side of the previous inequality by t_n and bearing in mind that $y^m \le 1 + y$, for all $m \in [0, 1]$ and y > 0, we get

$$\begin{split} \left\| \frac{w_n}{t_n} \right\|_{\infty} &\leq C_5 \left(\left\| \frac{w_n}{t_n} \right\|_{\infty}^{\frac{p-1}{p}} t_n^{\frac{q-2}{p}} + t_n^{q-2} + \left\| \frac{w_n}{t_n} \right\|_{\infty}^{q-1-\frac{q}{p}} t_n^{(q-2)\left(1+\frac{q}{p}\right)} + \frac{1}{n} \right) \\ &\leq C_5 \left(t_n^{q-2} + \left(t_n^{\frac{q-2}{p}} + t_n^{(q-2)\left(1+\frac{q}{p}\right)} \right) \left(1 + \left\| \frac{w_n}{t_n} \right\|_{\infty} \right) + \frac{1}{n} \right) \\ &\leq C_5 \left(t_n^{\frac{q-2}{p}} + 2t_n^{\frac{q-2}{p}} \left(1 + \left\| \frac{w_n}{t_n} \right\|_{\infty} \right) + \frac{1}{n} \right). \end{split}$$

It follows that

$$\left(1 - 2C_5 t_n^{\frac{q-2}{p}}\right) \left\| \frac{w_n}{t_n} \right\|_{\infty} \le 3C_5 t_n^{\frac{q-2}{p}} + \frac{C_5}{n},$$

and, as a consequence,

$$\lim_{n\to+\infty}\left\|\frac{w_n}{t_n}\right\|_{\infty}=0,$$

i.e.,

$$\frac{u_n}{t_n} \to \phi_1$$
 uniformly in $\overline{\Omega}$.

So, fixing $\gamma \in (0, \|\phi_1\|_{\infty})$, we can find $E \subset \Omega$, with |E| > 0, and $\tilde{n} \in \mathbb{N}$, $\tilde{n} \geq \bar{n}$, such that

$$u_n(x) \ge \gamma t_n$$
, for all $n \ge \tilde{n}$ and $x \in E$.

At this point, set

$$\delta := \sup_{t>0} (g(t)t - 2G(t)) \in [0, +\infty),$$

and let $\bar{t} > 0$ such that

$$g(t)t - 2G(t) \le -\frac{(\delta+1)|\Omega|}{|E|}$$
, for all $t \ge \overline{t}$,

and $n^* \ge \tilde{n}$ such that $t_n \ge \frac{\tilde{t}}{\gamma}$ for all $n \ge n^*$. Then, for all $n \ge n^*$, taking also (2.5) into account, we obtain

$$0 < \int_{\Omega} (g(u_n)u_n - 2G(u_n))dx$$

$$= \int_{\Omega \setminus E} (g(u_n)u_n - 2G(u_n))dx + \int_{E} (g(u_n)u_n - 2G(u_n))dx$$

$$\leq \delta |\Omega| - (\delta + 1)|\Omega| < 0,$$

a contradiction. Therefore, the sequence $\{t_n\}_{n\geq \bar{n}}$ is bounded in \mathbb{R} and (2.8) yields the boundedness of $\{w_n\}_{n\geq \bar{n}}$ in $H^1_0(\Omega)$, as well. As a consequence, we get the boundedness of $\{u_n\}_{n\geq \bar{n}}$ in $H^1_0(\Omega)$, as desired.

Collecting the results of the previous lemmas, it is now easy to derive our first existence result.

Theorem 2.3. For all $\lambda > 0$, problem (P_{λ}) has at least one non-zero solution.

Proof. Let $\{u_n\}$ be the sequence of solutions to (P_n) in Lemma 2.1. By Lemma 2.2 there exists $u^* \in H_0^1(\Omega)$ such that, up to a subsequence,

$$u_n \to u^*$$
 in $H_0^1(\Omega)$, $u_n \to u^*$ in $L^p(\Omega)$, for all $p \in [1,2^*)$.

Fixing $v \in H_0^1(\Omega)$ and taking the limit as $n \to +\infty$ in the identity $I'_n(u_n)(v) = 0$, we get $I'_{\lambda}(u^*)(v) = 0$, i.e. u^* is a weak solution to (P_{λ}) . To justify that $u^* \neq 0$, observe that, by (2.5) one has

$$0 < \left(\frac{1}{2} - \frac{1}{k}\right) R^{2}$$

$$\leq \lambda \int_{\Omega} (g(u_{n})u_{n}dx - 2G(u_{n})) dx$$

$$\leq \lambda k_{1} \left(\|u_{n}\|_{1} + \|u_{n}\|_{q}^{q}\right) + 2\lambda k_{1} \left(\|u_{n}\|_{1} + \frac{1}{q} \|u_{n}\|_{q}^{q}\right),$$

and so, letting $n \to +\infty$, the conclusion is achieved.

We now show that, when λ approaches zero, every non-zero solution to (P_{λ}) is actually positive. To this aim, for all $\lambda > 0$, set

$$S_{\lambda} := \{ u \in H_0^1(\Omega) \setminus \{0\} : u \text{ is a solution to } (P_{\lambda}) \},$$

and denote by \mathcal{P} the interior of the positive cone of $C_0^1(\overline{\Omega})$, i.e.

$$\mathcal{P} := \left\{ u \in C_0^1(\overline{\Omega}) : u > 0 \text{ in } \Omega, \ \frac{\partial u}{\partial \nu} < 0 \text{ on } \partial \Omega \right\},\,$$

 ν being the unit outer normal to $\partial\Omega$. Our second result reads as follows:

Theorem 2.4. There exists $\Lambda^* > 0$ such that for each $\lambda \in (0, \Lambda^*)$, $S_{\lambda} \subset \mathcal{P}$.

Proof. We first observe that, by the regularity theory of elliptic equations, for all $\lambda > 0$ and $u_{\lambda} \in S_{\lambda}$, one has $u_{\lambda} \in C^{1,\alpha}(\overline{\Omega})$, for some $\alpha \in (0,1)$.

If $u_{\lambda} \in S_{\lambda}$, it is straightforward to check that $v_{\lambda} := \lambda^{-1}u_{\lambda}$ is a solution to the problem

$$\begin{cases}
-\Delta u = \lambda_1 u + g(\lambda u) & \text{in } \Omega \\
u \ge 0 & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(\tilde{P}_{λ})

clearly equivalent to (P_{λ}) . Note that (g_2) ensures the existence of some a > 0 such that g(t) < 0 for all $t \in (0, a)$, and moreover it must hold

$$||v_{\lambda}||_{\infty} \ge \frac{a}{\lambda},\tag{2.9}$$

otherwise we would get $g(u_{\lambda}) < 0$ in $\Omega \setminus u_{\lambda}^{-1}(0)$, and so

$$\|u_{\lambda}\|^2 - \lambda_1 \|u_{\lambda}\|_2^2 = \lambda \int_{\Omega} g(u_{\lambda}) u_{\lambda} dx < 0,$$

against the definition of λ_1 . From now on, we will then focus on (\tilde{P}_{λ}) . We split the proof in several steps.

Step 1. We show that there exist two constants C^* , $\Lambda_0 > 0$ such that, for any $\lambda \in (0, \Lambda_0]$ and for any $v_{\lambda} \in S_{\lambda}$,

$$||v_{\lambda}|| \ge \frac{C^*}{\lambda}.\tag{2.10}$$

Fix $\beta > \max\{\frac{N}{2}, \frac{1}{q-1}\}$. By [1, Theorem 8.2] and the embedding $W^{2,\beta}(\Omega) \hookrightarrow C^1(\overline{\Omega})$, one has $v_{\lambda} \in W^{2,\beta}(\Omega)$ and there exists a constant $C_0 > 0$, independent of λ , such that

$$\|v_{\lambda}\|_{C^{1}(\overline{\Omega})} \le C_{0}\left((\lambda_{1}+1)\|v_{\lambda}\|_{\beta}+\|g(\lambda v_{\lambda})\|_{\beta}\right).$$
 (2.11)

So, by (g_1) and Hölder's inequality, we get

$$\begin{split} \int_{\Omega} |g(\lambda v_{\lambda})|^{\beta} dx &\leq k_{1}^{\beta} \int_{\Omega} \left(1 + (\lambda v_{\lambda})^{q-1}\right)^{\beta} dx \\ &\leq 2^{\beta-1} k_{1}^{\beta} \left(|\Omega| + \lambda^{\beta(q-1)} \left\|v_{\lambda}\right\|_{\infty}^{\beta(q-1)-1} \left\|v_{\lambda}\right\|_{1}\right), \end{split}$$

and therefore

$$\begin{split} \|v_{\lambda}\|_{\infty} &\leq C_{0} \left((\lambda_{1}+1) \|v_{\lambda}\|_{\infty}^{\frac{\beta-1}{\beta}} \|v_{\lambda}\|_{1}^{\frac{1}{\beta}} \\ &+ 2^{\frac{\beta-1}{\beta}} k_{1} \left(|\Omega|^{\frac{1}{\beta}} + \lambda^{q-1} \|v_{\lambda}\|_{\infty}^{q-1-\frac{1}{\beta}} \|v_{\lambda}\|_{1}^{\frac{1}{\beta}} \right) \right). \end{split}$$

Now, dividing by $\|v_{\lambda}\|_{\infty}^{\frac{\beta-1}{\beta}}$ both sides of the previous inequality and taking (2.9) into account, we obtain,

$$\left(\frac{a}{\lambda}\right)^{\frac{1}{\beta}} \leq \|v_{\lambda}\|_{\infty}^{\frac{1}{\beta}} \leq C_{1} \left(\|v_{\lambda}\|_{1}^{\frac{1}{\beta}} + \|v_{\lambda}\|_{\infty}^{\frac{1-\beta}{\beta}} + \lambda^{q-1} \|v_{\lambda}\|_{\infty}^{q-2} \|v_{\lambda}\|_{1}^{\frac{1}{\beta}}\right)
\leq C_{1} \left(\|v_{\lambda}\|_{1}^{\frac{1}{\beta}} + a^{\frac{1-\beta}{\beta}} \lambda^{\frac{\beta-1}{\beta}} + a^{q-2} \lambda \|v_{\lambda}\|_{1}^{\frac{1}{\beta}}\right)
\leq C_{2} \left((1+\lambda) \|v_{\lambda}\|_{\beta}^{\frac{1}{\beta}} + \lambda^{\frac{\beta-1}{\beta}}\right).$$
(2.12)

Now, if $0 < \lambda \le \min\{1, a(2C_2)^{-\beta}\} := \Lambda_0$, one has

$$\|v_{\lambda}\|^{\frac{1}{\beta}} \geq \frac{1}{2C_2} \left(\frac{a}{\lambda}\right)^{\frac{1}{\beta}} - \frac{1}{2} \geq \frac{1}{4C_2} \left(\frac{a}{\lambda}\right)^{\frac{1}{\beta}}$$

and hence (2.10) is fulfilled with $C^* = a(4C_2)^{-\beta}$. Since of course $||v_{\lambda}|| \to +\infty$ as $\lambda \to 0^+$, by (2.12) we can determine $C_3 > 0$ and $\Lambda_1 \in (0, \Lambda_0]$ such that $||v_{\lambda}|| \ge 1$ and

$$\|v_{\lambda}\|_{\infty} \le C_3 \|v_{\lambda}\| \tag{2.13}$$

for any $\lambda \in (0, \Lambda_1]$. For the rest of the proof, we assume $\lambda \in (0, \Lambda_1]$.

Step 2. We now show that, writing v_{λ} as

$$v_{\lambda} = t_{\lambda}\phi_1 + w_{\lambda}$$

with $t_{\lambda} \in \mathbb{R}$ and $w_{\lambda} \in \operatorname{span}\{\phi_1\}^{\perp}$, then it holds

$$\|w_{\lambda}\|_{C^{1}(\overline{\Omega})} \leq \tilde{C} \|v_{\lambda}\|^{\frac{q}{2}}, \tag{2.14}$$

for some $\tilde{C} > 0$. By the same arguments as [3], it is easily seen that $t_{\lambda} > 0$ and that w_{λ} is a weak solution to

$$\begin{cases}
-\Delta u = \lambda_1 u + g(\lambda v_\lambda) & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega.
\end{cases}$$
(2.15)

The relation $I'_{\lambda}(v_{\lambda})(\phi_1) = 0$ and the definition of ϕ_1 imply that

$$\int_{\Omega} \nabla v_{\lambda} \nabla \phi_{1} dx - \lambda_{1} \int_{\Omega} v_{\lambda} \phi_{1} dx - \int_{\Omega} g(\lambda v_{\lambda}) \phi_{1} dx = -\int_{\Omega} g(\lambda v_{\lambda}) \phi_{1} dx = 0,$$

and therefore

$$\int_{\Omega} g(\lambda v_{\lambda}) w_{\lambda} dx = \int_{\Omega} g(\lambda v_{\lambda}) (v_{\lambda} - t_{\lambda} \phi_{1}) dx = \int_{\Omega} g(\lambda v_{\lambda}) v_{\lambda} dx.$$

So, we get

$$||w_{\lambda}||^{2} = \lambda_{1} ||w_{\lambda}||_{2}^{2} + \int_{\Omega} g(\lambda v_{\lambda}) w_{\lambda} dx$$

$$\leq \frac{\lambda_{1}}{\lambda_{2}} ||w_{\lambda}||^{2} + \int_{\Omega} g(\lambda v_{\lambda}) v_{\lambda} dx$$

$$\leq \frac{\lambda_{1}}{\lambda_{2}} ||w_{\lambda}||^{2} + k_{1} \left(||v_{\lambda}||_{1} + \lambda^{q-1} ||v_{\lambda}||_{q}^{q} \right)$$

$$\leq \frac{\lambda_{1}}{\lambda_{2}} ||w_{\lambda}||^{2} + C_{4} ||v_{\lambda}||^{q},$$

from which we deduce the estimate

$$||w_{\lambda}||^{2} \le C_{5} ||v_{\lambda}||^{q}, \tag{2.16}$$

being $C_5 = \frac{\lambda_2 C_4}{\lambda_2 - \lambda_1}$. By applying the same arguments as before to the function w_λ and bearing in mind also (2.13) and (2.16), we obtain

$$\begin{split} \|w_{\lambda}\|_{C^{1}(\overline{\Omega})} &\leq C_{6}\left((\lambda_{1}+1)\|w_{\lambda}\|_{\beta} + \|g(\lambda v_{\lambda})\|_{\beta}\right) \\ &\leq C_{6}\left((\lambda_{1}+1)\|w_{\lambda}\|_{\infty}^{\frac{\beta-1}{\beta}}\|w_{\lambda}\|_{1}^{\frac{1}{\beta}} + 2^{\frac{\beta-1}{\beta}}k_{1}\left(|\Omega|^{\frac{1}{\beta}} + \lambda^{q-1}\|v_{\lambda}\|_{\infty}^{q-1-\frac{1}{\beta}}\|v_{\lambda}\|_{1}^{\frac{1}{\beta}}\right)\right) \\ &\leq C_{7}\left(\|w_{\lambda}\|_{C^{1}(\overline{\Omega})}^{\frac{\beta-1}{\beta}}\|v_{\lambda}\|^{\frac{q}{2\beta}} + 1 + \lambda^{q-1}\|v_{\lambda}\|^{q-1}\right) \\ &\leq C_{7}\left(\|w_{\lambda}\|_{C^{1}(\overline{\Omega})}^{\frac{\beta-1}{\beta}}\|v_{\lambda}\|^{\frac{q}{2\beta}} + 2\|v_{\lambda}\|^{q-1}\right). \end{split}$$

So, either

$$\|w_{\lambda}\|_{C^{1}(\overline{\Omega})} \leq 2C_{7} \|w_{\lambda}\|_{C^{1}(\overline{\Omega})}^{\frac{\beta-1}{\beta}} \|v_{\lambda}\|_{2^{\frac{q}{2\beta}}}^{\frac{q}{2\beta}}$$

or

$$||w_{\lambda}||_{C^{1}(\overline{\Omega})} \leq 4C_{7} ||v_{\lambda}||^{q-1}.$$

In any case, we get

$$\|w_{\lambda}\|_{C^{1}(\overline{\Omega})} \leq \tilde{C} \|v_{\lambda}\|^{\frac{q}{2}}, \tag{2.17}$$

where $\tilde{C} = 4C_7$, as desired.

Step 3 (conclusion). Taking (2.10) and (2.16) into account, for $0 < \lambda \leq \min\{1, \Lambda_0, \Lambda_1, \Lambda_2\}$, where $\Lambda_2 := \left(\frac{1}{2C_5}\right)^{\frac{1}{2-q}}C^*$, we obtain

$$t_{\lambda}^{2} \geq \frac{\|v_{\lambda}\|^{2} - C_{5} \|v_{\lambda}\|^{q}}{\|\phi_{1}\|^{2}} \geq \frac{\|v_{\lambda}\|^{2}}{\|\phi_{1}\|^{2}} \left(1 - \frac{C_{5}C^{*q-2}}{\lambda^{q-2}}\right) \geq \frac{\|v_{\lambda}\|^{2}}{2 \|\phi_{1}\|^{2}} = C_{8} \|v_{\lambda}\|^{2}, \quad (2.18)$$

where $C_8 = \frac{1}{2\|\phi_1\|^2}$. For this range of λ , in view of (2.17), we then obtain

$$\left\| t_{\lambda}^{-1} v_{\lambda} - \phi_{1} \right\|_{C^{1}(\overline{\Omega})} = t_{\lambda}^{-1} \left\| w_{\lambda} \right\|_{C^{1}(\overline{\Omega})} \leq \tilde{C} C_{8}^{-\frac{1}{2}} \left\| v_{\lambda} \right\|^{\frac{q}{2} - 1} \leq C_{9} \lambda^{1 - \frac{q}{2}}$$

with $C_9 = \tilde{C}C_8^{-\frac{1}{2}}C^{*\frac{q}{2}-1}$. Since $\phi_1 \in \mathcal{P}$ and \mathcal{P} is an open subset of $C^1(\overline{\Omega})$, there exists $\delta > 0$ such that

$$\{u \in C^1(\overline{\Omega}) : \|u - \phi_1\|_{C^1(\overline{\Omega})} < \delta\} \subset \mathcal{P}.$$

So, setting $\Lambda_3 := \left(\frac{\delta}{C_9}\right)^{\frac{2}{2-q}}$, for all $0 < \lambda \leq \min\{1, \Lambda_0, \Lambda_1, \Lambda_2, \Lambda_3\} := \Lambda^*$, one has $t_\lambda^{-1} v_\lambda \in \mathcal{P}$ and hence $v_\lambda \in \mathcal{P}$. This concludes the proof.

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References

- [1] S. Agmon, The L^p approach to the Dirichlet problem I. Regularity theorems, *Ann. Scuola Norm. Sup. Pisa* **13**(1959), 405–448. MR125306
- [2] S. Ahmad, A. C. Lazer, J. L. Paul, Elementary critical point theory and perturbations of elliptic boundary value problems at resonance, *Indiana Univ. Math. J.* **25**(1976), 933–944. https://doi.org/10.1512/iumj.1976.25.25074
- [3] G. Anello, Existence results and strong maximum principle for a resonant sublinear elliptic problem, *Minimax Theory Appl.* **4**(2019), No. 2, 217–229. MR3973626
- [4] G. Anello, L. Vilasi, Uniqueness of positive and compacton-type solutions for a resonant quasilinear problem, *Topol. Methods Nonlinear Anal.* **49**(2016), No. 2, 565–575. https://doi.org/10.12775/tmna.2016.090
- [5] P. Bartolo, V. Benci, D. Fortunato, Abstract critical point theorems and applications to some nonlinear problems with strong resonance at infinity, *Nonlinear Anal.* 7(1983), 981–1012. https://doi.org/10.1016/0362-546X(83)90115-3
- [6] D. G. Costa, E. A. B. Silva, Existence of solution for a class of resonant elliptic problems, *J. Math. Anal. Appl.* **175**(1993), 411–424. https://doi.org/10.1006/jmaa.1993.1180
- [7] J. V. A. Gonçalves, O. H. Miyagaki, Three solutions for a strongly resonant elliptic problem, *Nonlinear Anal.* **24**(1995), No. 2, 265–272. https://doi.org/10.1016/0362-546X(94)E0016-A
- [8] Y. ILYASOV, Y. EGOROV, Hopf boundary maximum principle violation for semilinear elliptic equations, *Nonlinear Anal.* **72**(2010), 3346–3355. https://doi.org/10.1016/j.na.2009. 12.015
- [9] E. M. LANDESMAN, A. C. LAZER, Nonlinear perturbations of linear elliptic boundary value problems at resonance, *J. Math. Mech.* **19**(1970), 609–623. MR0267269
- [10] S. Liu, Multiple solutions for elliptic resonant problems, *Proc. Roy. Soc. Edinburgh Sect. A* **138**(2008), No. 6, 1281–1289. https://doi.org/10.1017/S0308210507000443
- [11] M. Schechter, Strong resonance problems for elliptic semilinear boundary value problems, *J. Operator Theory* **30**(1993), No. 2, 301–314. MR1305509
- [12] K. Thews, Nontrivial solutions of elliptic equations at resonance, *Proc. Roy. Soc. Edinburgh Sect. A* **85**(1980), 119–129. https://doi.org/10.1017/S0308210500011732
- [13] J. L. Vázquez, A strong maximum principle for some quasilinear elliptic equations, *Appl. Math. Optim.* **12**(1984), 191–202. https://doi.org/10.1007/BF01449041

- [14] X. P. Wu, C. L. Tang, Some existence theorems for elliptic resonant problems, *J. Math. Anal. Appl.* **264**(2001), 133–146. https://doi.org/10.1006/jmaa.2001.7660
- [15] E. Zeidler, Nonlinear functional analysis and its applications III, Springer, Berlin, 1985. https://doi.org/10.1007/978-1-4612-5020-3