Research Article

Global Existence of Solutions to the Fowler Equation in a Neighbourhood of Travelling-Waves

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We investigate a fractional diffusion/anti-diffusion equation proposed by Andrew C. Fowler to describe the dynamics of sand dunes sheared by a fluid flow. In this paper, we prove the global-intime well-posedness in the neighbourhood of travelling-waves solutions of the Fowler equation.

1. Introduction

The study of mechanisms that allow the formation of structures such as sand dunes and ripples at the bottom of a fluid flow plays a crucial role in the understanding of coastal dynamics. The modeling of these phenomena is particularly complex since we must not only solve the Navier-Stokes or Saint-Venant equations with equation for sediment transport, but also take into account the evolution of the bottom. Instead of solving the whole system fluid flow, free surface and free bottom, nonlocal models of fluid flow interacting with the bottom were introduced in [1, 2]. Among these models, we are interested in the following nonlocal conservation law [1, 3]:

$$\partial_t u(t,x) + \partial_x \left(\frac{u^2}{2}\right)(t,x) + \mathcal{O}[u(t,\cdot)](x) - \partial_{xx}^2 u(t,x) = 0, \quad t \in (0,T), \ x \in \mathbb{R},$$

$$u(0,x) = u_0(x), \quad x \in \mathbb{R},$$
(1.1)

where *T* is any given positive time, u = u(t, x) represents the dune height (see Figure 1), and \mathcal{I} is a nonlocal operator defined as follows: for any Schwartz function $\varphi \in \mathcal{S}(\mathbb{R})$ and any

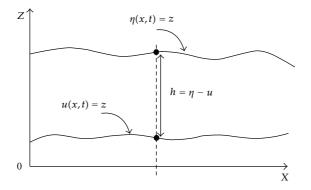


Figure 1: Domain considered for the Fowler model: *h* is the depth water, η the free surface, and *u* the sea bottom.

 $x \in \mathbb{R}$,

$$\mathcal{O}[\varphi](x) := \int_0^{+\infty} |\xi|^{-1/3} \varphi''(x-\xi) d\xi.$$
(1.2)

Equation (1.1) is valid for a river flow over an erodible bottom u(t, x) with slow variation and describes both accretion and erosion phenomena [4]. See [4, 5] for numerical results on this equation.

The nonlocal term \mathcal{I} can be seen as a fractional power of order 2/3 of the Laplacian with the bad sign. Indeed, it has been proved [4]

$$\mathcal{F}(\mathcal{O}[\varphi] - \varphi'')(\xi) = \psi_{\mathcal{O}}(\xi) \mathcal{F}\varphi(\xi), \tag{1.3}$$

where

$$\psi_{\mathcal{I}}(\xi) = 4\pi^2 \xi^2 - a_{\mathcal{I}} |\xi|^{4/3} + i b_{\mathcal{I}} \xi |\xi|^{1/3}, \qquad (1.4)$$

with $a_{\mathcal{I}}, b_{\mathcal{I}}$ positive constants, \mathcal{F} denotes the Fourier transform defined in (1.7), and Γ denotes the Euler function. One simple way to establish this fact is the derivation of a new formula for the operator \mathcal{I} , see Proposition 2.5.

Remark 1.1. For causal functions (i.e., $\varphi(x) = 0$ for x < 0), this operator is up to a multiplicative constant, the Riemann-Liouville fractional derivative operator which is defined as follows [6]:

$$\frac{1}{\Gamma(2/3)} \int_0^{+\infty} \frac{\varphi''(x-\xi)}{|\xi|^{1/3}} d\xi = \frac{d^{-2/3}}{dx^{-2/3}} \varphi''(x) = \frac{d^{4/3}}{dx^{4/3}} \varphi(x).$$
(1.5)

Therefore, the Fowler model has two antagonistic terms: a usual diffusion and a nonlocal fractional anti-diffusive term of lower order. This remarkable feature enabled to apply this model for signal processing. Indeed, the diffusion is used to reduce the noise whereas the nonlocal anti-diffusion is used to enhance the contrast [7].

Recently, some results regarding this equation have been obtained, namely, existence of travelling-waves $u_{\phi}(t, x) = \phi(x - ct)$ where $\phi \in C_b^1(\mathbb{R})$ and $c \in \mathbb{R}$ represents wave

velocity, the global well-posedness for L^2 -initial data, the failure of the maximum principle, and the local-in-time well-posedness in a subspace of C_b^1 [4, 8]. Notice that the travellingwaves are not necessarily of solitary type (see [8]) and therefore may not belong to $L^2(\mathbb{R})$, the space where a global well-posedness result is available. In [8], the authors prove local well-posedness in a subspace of $C_b^1(\mathbb{R})$ but fail to obtain global existence.

To prove the existence of travelling-waves solutions of the Fowler equation, the authors used the implicit function theorem on suitable Banach spaces [8]. Much work has been devoted to investigate existence, uniqueness, and regularity of travelling-waves for integral differential equations, see for instance [9] and references therein.

An interesting topic is to know if the shape of this travelling-wave is maintained when it is perturbed. This raises the question of the stability of travelling-waves. But before interesting ourselves in this problem, we have to show first the global existence of perturbations around these travelling-waves. Hence in this paper, we prove the global wellposedness in an L^2 -neighbourhood of a regular travelling-wave, namely $u = u_{\phi} + v$. To prove this result, we consider the following Cauchy problem:

$$\partial_t v(t,x) + \partial_x \left(\frac{v^2}{2} + u_{\phi} v\right)(t,x) + \mathcal{O}[v(t,\cdot)](x) - \partial_{xx}^2 v(t,x) = 0, \quad t \in (0,T), \ x \in \mathbb{R},$$

$$v(0,x) = v_0(x), \quad x \in \mathbb{R},$$
(1.6)

where $v_0 \in L^2(\mathbb{R})$ is an initial perturbation and *T* is any given positive time.

To prove the existence and uniqueness results, we begin by introducing the notion of *mild* solution (see Definition 2.1) based on Duhamel's formula (2.1), in which the kernel *K* of $\mathcal{O} - \partial_{xx}^2$ appears. The use of this formula allows to prove the local-in-time existence with the help of a contracting fixed point theorem. The global existence is obtained thanks to an energy estimate (4.68). This approach is classical: we refer for instance to [4, 10].

The plan of this paper is organised as follows. In the next section, we define the notion of mild solution to (1.6) and we give some properties on the kernel *K* of $\mathcal{O} - \partial_{xx}^2$ that will be needed in the sequel. Sections 3 and 4 are, respectively, devoted to the proof of the uniqueness and the existence of a mild solution for (1.6). Section 5 contains the proof of the regularity of the solution.

Notations

(i) The norm of a measurable function $f \in L^p(\mathbb{R})$ is written $||f||_{L^p(\mathbb{R})}^p = \int_{\mathbb{R}} |f(x)|^p dx$ for $1 \le p < \infty$.

(ii) We denote by \mathcal{F} the Fourier transform of f which is defined by the following: for all $\xi \in \mathbb{R}$

$$\mathcal{F}f(\xi) := \int_{\mathbb{R}} e^{-2i\pi x\xi} f(x) dx, \qquad (1.7)$$

and \mathcal{F}^{-1} denotes the inverse of Fourier transform.

- (iii) The Schwartz space of rapidly decreasing functions on \mathbb{R} is denoted by $\mathcal{S}(\mathbb{R})$.
- (iv) We write $C^k(\mathbb{R}) = \{ f : \mathbb{R} \to \mathbb{C}; f, f', \dots, f^{(k)} \text{ are continous on } \mathbb{R} \}.$

(v) We denote by $C_b(\mathbb{R})$ the space of all bounded continuous real-valued functions on \mathbb{R} with the norm $\|\cdot\|_{L^{\infty}} = \sup_{\mathbb{R}} |f|$.

(vi) We write for any T > 0,

$$C^{1,2}((0,T] \times \mathbb{R}) := \left\{ u \in C((0,T] \times \mathbb{R}); \partial_t u, \partial_x u, \partial_{xx}^2 u \in C((0,T] \times \mathbb{R}) \right\}.$$
 (1.8)

(vii) We denote by $\mathfrak{D}(U)$ the space of test functions on U and $\mathfrak{D}'(U)$ denotes the distribution space.

Here is our main result.

Theorem 1.2. Let T > 0 and $v_0 \in L^2(\mathbb{R})$. There exists a unique mild solution $v \in L^{\infty}((0,T); L^2(\mathbb{R}))$ of (1.6) (see Definition 2.1) which satisfies

$$v \in C([0,T]; L^2(\mathbb{R})), \quad v(0,\cdot) = v_0 \text{ almost everywhere.}$$
 (1.9)

Moreover, if $\phi \in C_{h}^{2}(\mathbb{R})$ *then* $v \in C^{1,2}((0,T] \times \mathbb{R})$ *and satisfies*

$$\partial_t v + \partial_x \left(\frac{v^2}{2} + u_\phi v \right) + \mathcal{O}[v] - \partial_{xx}^2 v = 0, \qquad (1.10)$$

on $(0,T] \times \mathbb{R}$, in the classical sense or equivalently, $u = u_{\phi} + v$ is a classical solution of (1.1).

2. Duhamel Formula and Main Properties of K

Definition 2.1. Let T > 0 and $v_0 \in L^2(\mathbb{R})$. We say that $v \in L^{\infty}((0,T); L^2(\mathbb{R}))$ is a mild solution to (1.6) if for any $t \in (0,T)$:

$$v(t,\cdot) = K(t,\cdot) * v_0 - \int_0^t \partial_x K(t-s,\cdot) * \left(\frac{v^2}{2} + u_\phi v\right)(s,\cdot) \, ds, \tag{2.1}$$

where $K(t, x) = \mathcal{F}^{-1}(e^{-t\psi_{\mathcal{I}}(\cdot)})(x)$ is the kernel of the operator $\mathcal{O} - \partial_{xx'}^2$ and $\psi_{\mathcal{O}}$ is defined in (1.4).

The expression (2.1) is the Duhamel formula and is obtained using the spatial Fourier transform.

Proposition 2.2 (main properties of *K*, [4]). The kernel K satisfies

- (1) $\forall t > 0, K(t, \cdot) \in L^1(\mathbb{R}) \text{ and } K \in C^{\infty}((0, \infty) \times \mathbb{R}),$
- $(2) \ \forall s,t>0, \ K(s,\cdot)*K(t,\cdot)=K(s+t,\cdot), \ \forall u_0\in L^2(\mathbb{R}), \ \lim_{t\to 0}K(t,\cdot)*u_0=u_0 \ in \ L^2(\mathbb{R}),$
- (3) $\forall T > 0, \exists K_0 \text{ such that } \forall t \in (0, T], ||\partial_x K(t, \cdot)||_{L^2(\mathbb{R})} \leq K_0 t^{-3/4},$
- (4) $\forall T > 0, \exists K_1 \text{ such that } \forall t \in (0, T], ||\partial_x K(t, \cdot)||_{L^1(\mathbb{R})} \leq K_1 t^{-1/2}.$

Remark 2.3. An interesting property for the kernel *K* is the non-positivity (see Figure 2), and the main consequence of this feature is the failure of maximum principle [4]. We use again this property to show that the constant solutions of the Fowler equation are unstable [11].

Remark 2.4. Using Plancherel formula, we have for any $v_0 \in L^2(\mathbb{R})$ and any $t \in (0,T]$

$$\|K(t,\cdot) * v_0\|_{L^2(\mathbb{R})} \le e^{\alpha_0 t} \|v_0\|_{L^2(\mathbb{R})},$$
(2.2)

where $\alpha_0 = -\min \operatorname{Re}(\psi_{\mathcal{D}}) > 0$.

Proposition 2.5 (integral formula for \mathcal{O}). *For all* $\varphi \in \mathcal{S}(\mathbb{R})$ *and all* $x \in \mathbb{R}$ *,*

$$\mathcal{O}[\varphi](x) = \frac{4}{9} \int_{-\infty}^{0} \frac{\varphi(x+z) - \varphi(x) - \varphi'(x)z}{|z|^{7/3}} dz.$$
(2.3)

Proof. The proof is based on simple integrating by parts. The regularity and the rapidly decreasing of φ ensure the validity of the computations that follow. We have

$$\begin{split} \int_{0}^{+\infty} \varphi''(x-\xi) |\xi|^{-1/3} d\xi &= \int_{0}^{+\infty} \frac{d}{d\xi} \big(\varphi'(x) - \varphi'(x-\xi) \big) |\xi|^{-1/3} d\xi \\ &= \frac{1}{3} \int_{0}^{+\infty} |\xi|^{-4/3} \big(\varphi'(x) - \varphi'(x-\xi) \big) d\xi \\ &= \frac{1}{3} \int_{0}^{+\infty} |\xi|^{-4/3} \frac{d}{d\xi} \big(\varphi'(x)\xi + \varphi(x-\xi) - \varphi(x) \big) d\xi \\ &= \frac{4}{9} \int_{0}^{+\infty} \frac{\varphi(x-\xi) - \varphi(x) + \varphi'(x)\xi}{|\xi|^{7/3}} d\xi \\ &= \frac{4}{9} \int_{-\infty}^{0} \frac{\varphi(x+\xi) - \varphi(x) - \varphi'(x)\xi}{|\xi|^{7/3}} d\xi. \end{split}$$
(2.4)

There is no boundary term at infinity (resp., at zero) because φ is a rapidly decreasing function on \mathbb{R} (resp., φ is smooth).

Remark 2.6. Using integral formula (2.3), [4, 8] proved that

$$\mathcal{F}(\mathcal{I}[\varphi])(\xi) = 4\pi^2 \Gamma\left(\frac{2}{3}\right) |\xi|^{4/3} \left(-\frac{1}{2} + i\frac{\sqrt{3}}{2} \operatorname{sgn}(\xi)\right) \mathcal{F}\varphi(\xi).$$
(2.5)

Notice that $\mathcal{F}(\mathcal{O}[\varphi])(\xi) = 4\pi^2 \Gamma(2/3)(i\xi)^{4/3}$ which is consistent with Remark 1.1: up to a multiplicative constant $\mathcal{O}[\varphi]$ is $d^{4/3}\varphi/dx^{4/3}$.

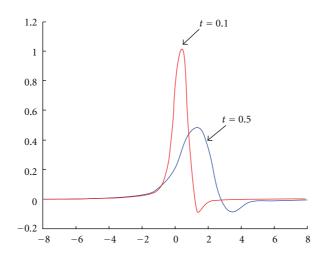


Figure 2: Evolution of the kernel *K* for t = 0.1 (red) and t = 0.5 s (blue).

Proposition 2.7. Let $s \in \mathbb{R}$ and $\varphi \in H^{s}(\mathbb{R})$. Then $\mathcal{O}[\varphi] \in H^{s-4/3}(\mathbb{R})$ and we have

$$\left\|\mathcal{O}[\varphi]\right\|_{H^{s-4/3}(\mathbb{R})} \leq 4\pi^2 \Gamma\left(\frac{2}{3}\right) \left\|\varphi\right\|_{H^s(\mathbb{R})}.$$
(2.6)

Proof. For all $s \in \mathbb{R}$ and all $\varphi \in H^s(\mathbb{R})$, we have, using Remark 2.6

$$\begin{split} \|\mathcal{D}[\varphi]\|_{H^{s-4/3}(\mathbb{R})} &= \left(\int_{\mathbb{R}} \left(1 + |\xi|^{2} \right)^{s-4/3} |\mathcal{F}(\mathcal{D}[\varphi])(\xi)|^{2} d\xi \right)^{1/2} \\ &= 4\pi^{2} \Gamma\left(\frac{2}{3}\right) \left(\int_{\mathbb{R}} \left(1 + |\xi|^{2} \right)^{s-4/3} \left| \frac{1}{2} - i \operatorname{sgn}(\xi) \frac{\sqrt{3}}{2} \right| |\xi|^{8/3} |\mathcal{F}(\varphi)(\xi)|^{2} d\xi \right)^{1/2} \\ &= 4\pi^{2} \Gamma\left(\frac{2}{3}\right) \left(\int_{\mathbb{R}} \left(\frac{|\xi|^{2}}{1 + |\xi|^{2}} \right)^{4/3} \left(1 + |\xi|^{2} \right)^{s} |\mathcal{F}(\varphi)(\xi)|^{2} d\xi \right)^{1/2} \\ &\leq 4\pi^{2} \Gamma\left(\frac{2}{3}\right) \left[\int_{\mathbb{R}} \left(1 + |\xi|^{2} \right)^{s} |\mathcal{F}(\varphi)(\xi)|^{2} d\xi \right]^{1/2} \\ &= 4\pi^{2} \Gamma\left(\frac{2}{3}\right) \|\varphi\|_{H^{s}(\mathbb{R})}. \end{split}$$

$$(2.7)$$

Remark 2.8. From the previous proposition, we deduce that for all $s \in \mathbb{R}$ and all $\varphi \in H^s(\mathbb{R})$, $\mathcal{O}[\varphi] \in H^{s-4/3}(\mathbb{R})$. In particular, using the Sobolev embedding $H^{2/3}(\mathbb{R}) \hookrightarrow C_b(\mathbb{R}) \cap L^2(\mathbb{R})$, we deduce that $\mathcal{O}: H^2(\mathbb{R}) \to C_b(\mathbb{R}) \cap L^2(\mathbb{R})$ is a bounded linear operator.

Proposition 2.9 (Duhamel formula (2.1) is well defined). Let T > 0, $v_0 \in L^2(\mathbb{R})$ and $w \in L^{\infty}((0,T); L^1(\mathbb{R})) + L^{\infty}((0,T); L^2(\mathbb{R}))$. Then, the function

$$v: t \in (0,T] \longrightarrow K(t,\cdot) * v_0 - \int_0^t \partial_x K(t-s,\cdot) * w(s,\cdot) ds$$
(2.8)

is well defined and belongs to $C([0,T]; L^2(\mathbb{R}))$ (being extended at t = 0 by the value $v(0, \cdot) = v_0$).

Proof. From Proposition 2.2, it is easy to see that v is well defined and that for any $t \in (0, T]$, $v(t, \cdot) \in L^2(\mathbb{R})$. Indeed, $\forall t > 0$, $\partial_x K(t, \cdot) \in L^1(\mathbb{R}) \cap L^2(\mathbb{R})$ so by Young inequalities $\partial_x K(t, \cdot) * w(t, \cdot)$ exists, and using the estimates on the gradient (item 3 and 4 of Proposition 2.2) we deduce that v is well defined and $v(t, \cdot) \in L^2(\mathbb{R})$.

Let us prove the continuity of v. By the second item of Proposition 2.2, we have that the function $t \in (0,T] \rightarrow K(t, \cdot) * v_0$ is continuous and it is extended continuously up to t = 0 by the value $v(0, \cdot) = v_0$. We define the function

$$F: t \in [0,T] \longrightarrow \int_0^t \partial_x K(t-s,\cdot) * w(s,\cdot) ds.$$
(2.9)

Now, we are going to prove that *F* is uniformly continuous. For any h > 0, Young inequalities imply

$$\begin{aligned} \|F(t+h,\cdot) - F(t,\cdot)\|_{L^{2}} &\leq \int_{0}^{t} \|\partial_{x}K(t+h-s,\cdot) - \partial_{x}K(t-s,\cdot)\|_{L^{i}} \, ds \|w\|_{L^{\infty}((0,T);L^{j})} \\ &+ \int_{t}^{t+h} \|\partial_{x}K(t+h-s,\cdot)\|_{L^{i}} \, ds \|w\|_{L^{\infty}((0,T);L^{j})}, \end{aligned}$$
(2.10)

where $i, j \in \mathbb{N}^*$ are such that i + j = 3. Since $\partial_x K(t, \cdot) = \mathcal{F}^{-1}(\xi \to 2i\pi\xi e^{-t\psi_0(\xi)})$, the dominated convergence theorem implies that

$$\|\partial_x K(t-s+h,\cdot) - \partial_x K(t-s,\cdot)\|_{L^i(\mathbb{R})} \longrightarrow 0, \quad \text{as } h \longrightarrow 0.$$
(2.11)

Moreover, using the estimates on the gradient (items 3 and 4 of Proposition 2.2), we have the following inequality:

$$\int_{t}^{t+h} \|\partial_x K(t-s+h,\cdot)\|_{L^j(\mathbb{R})} ds \le c_j h^{\alpha_j},$$
(2.12)

where c_i is a positive constant and $\alpha_i = 1/2$ if j = 1, 1/4 if j = 2.

From (2.10), we obtain that $||F(t+h,\cdot)-F(t,\cdot)||_{L^2(\mathbb{R})} \to 0$, as $h \to 0$. Hence, the function F is continuous and this completes the proof of the continuity of v.

Remark 2.10. Using the semi group property of the kernel *K*, we have for all $t_0 \in (0, T)$ and all $t \in [0, T - t_0]$, [4]

$$v(t+t_0,\cdot) = K(t,\cdot) * v(t_0,\cdot) - \int_0^t \partial_x K(t-s,\cdot) * w(t_0+s,\cdot) ds.$$
(2.13)

3. Uniqueness of a Solution

Let us first establish the following Lemma.

Lemma 3.1. Let T > 0 and $v_0 \in L^2(\mathbb{R})$. For i = 1, 2, let $w_i \in L^{\infty}((0, T); L^1(\mathbb{R})) \cup L^{\infty}((0, T); L^2(\mathbb{R}))$ and define v_i as in Proposition 2.9 by

$$v_i(t,\cdot) = K(t,\cdot) * v_0 - \int_0^t \partial_x K(t-s,\cdot) * w_i(s,\cdot) ds.$$
(3.1)

Then,

$$\|v_{1} - v_{2}\|_{C([0,T];L^{2}(\mathbb{R}))} \leq \begin{cases} 4K_{0}T^{1/4}\|w_{1} - w_{2}\|_{L^{\infty}((0,T);L^{1}(\mathbb{R}))} & \text{if } w_{i} \in L^{\infty}((0,T);L^{1}(\mathbb{R})), \\ 2K_{1}\sqrt{T}\|w_{1} - w_{2}\|_{L^{\infty}((0,T);L^{2}(\mathbb{R}))} & \text{if } w_{i} \in L^{\infty}((0,T);L^{2}(\mathbb{R})). \end{cases}$$
(3.2)

Proof. For all $t \in [0, T]$, we have

$$v_1(t,\cdot) - v_2(t,\cdot) = -\int_0^t \partial_x K(t-s,\cdot) * (w_1 - w_2)(s,\cdot) ds.$$
(3.3)

Hence with the help of Young inequalities, we get

$$\|v_{1}(t,\cdot) - v_{2}(t,\cdot)\|_{L^{2}(\mathbb{R})} \leq \begin{cases} \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{2}(\mathbb{R})}\|(w_{1}-w_{2})(s,\cdot)\|_{L^{1}(\mathbb{R})}ds \\ & \text{if } w_{i} \in L^{\infty}((0,T);L^{1}(\mathbb{R})), \end{cases} \\ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}(\mathbb{R})}\|(w_{1}-w_{2})(s,\cdot)\|_{L^{2}(\mathbb{R})}ds \\ & \text{if } w_{i} \in L^{\infty}((0,T);L^{2}(\mathbb{R})). \end{cases}$$
(3.4)

It then follows that

$$\|v_{1}(t,\cdot) - v_{2}(t,\cdot)\|_{L^{2}(\mathbb{R})} \leq \begin{cases} \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{2}(\mathbb{R})} \ ds \|w_{1} - w_{2}\|_{L^{\infty}((0,T);L^{1}(\mathbb{R}))} \\ & \text{if } w_{i} \in L^{\infty}\big((0,T);L^{1}(\mathbb{R})\big), \\ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}(\mathbb{R})} \ ds \|w_{1} - w_{2}\|_{L^{\infty}((0,T);L^{2}(\mathbb{R}))} \\ & \text{if } w_{i} \in L^{\infty}\big((0,T);L^{2}(\mathbb{R})\big). \end{cases}$$
(3.5)

Using again the estimates of the gradient of *K* (see Proposition 2.2), we conclude the proof of this Lemma. \Box

Proposition 3.2. Let T > 0 and $v_0 \in L^2(\mathbb{R})$. There exists at most one $v \in L^{\infty}((0,T); L^2(\mathbb{R}))$ which is a mild solution to (1.6).

Proof. Let $v_1, v_2 \in L^{\infty}((0,T); L^2(\mathbb{R}))$ be two mild solutions to (1.6) and $t \in [0,T]$. Using the previous Lemma, we get

$$\|v_1 - v_2\|_{C([0,t];L^2(\mathbb{R}))} \le 2K_0 t^{1/4} \|v_1^2 - v_2^2\|_{L^{\infty}((0,t);L^1(\mathbb{R}))} + 2K_1 \sqrt{t} \|u_{\phi} v_1 - u_{\phi} v_2\|_{L^{\infty}((0,t);L^2(\mathbb{R}))}.$$
 (3.6)

Since,

$$\left\| v_1^2 - v_2^2 \right\|_{L^{\infty}((0,t);L^1(\mathbb{R}))} \le M \| v_1 - v_2 \|_{C([0,t];L^2(\mathbb{R}))}, \tag{3.7}$$

with $M = \|v_1\|_{C([0,T];L^2(\mathbb{R}))} + \|v_2\|_{C([0,T];L^2(\mathbb{R}))}$, then

$$\|v_1 - v_2\|_{C([0,t];L^2(\mathbb{R}))} \le \left(2MK_0t^{1/4} + 2K_1t^{1/2}\|u_{\phi}\|_{C_b^1(\mathbb{R})}\right)\|v_1 - v_2\|_{C([0,t];L^2(\mathbb{R}))}.$$
(3.8)

Therefore, $v_1 = v_2$ on [0, t] for any $t \in (0, T]$ satisfying $2MK_0t^{1/4} + 2K_1t^{1/2}||u_{\phi}||_{C_b^1(\mathbb{R})} < 1$. Since v_1 and v_2 are continuous with values in $L^2(\mathbb{R})$, we have that $v_1 = v_2$ on $[0, T_*]$ where T_* is the positive solution of the following:

$$2MK_0 t^{1/4} + 2K_1 t^{1/2} \| u_{\phi} \|_{C_b^1(\mathbb{R})} = 1,$$
(3.9)

that is, $T_* = ((-2MK_0 + \sqrt{4M^2K_0^2 + 8K_1 \|u_{\phi}\|_{C_b^1(\mathbb{R})}})/4K_1 \|u_{\phi}\|_{C_b^1(\mathbb{R})})^4$. To prove that $v_1 = v_2$ on [0, T], let us define

$$t_0 := \sup\{t \in [0, T] \text{ s.t } v_1 = v_2[0, t]\},$$
(3.10)

and we assume that $t_0 < T$. By continuity of v_1 and v_2 , we have that $v_1(t_0, \cdot) = v_2(t_0, \cdot)$. Using the semigroup property, see Remark 2.10, we deduce that $v_1(t_0 + \cdot, \cdot) = v_2(t_0 + \cdot, \cdot)$ are mild solutions to (1.6) with the same initial data $v_1(t_0, \cdot) = v_2(t_0, \cdot)$ which implies, from the first step of this proof, that $v_1(t, \cdot) = v_2(t, \cdot)$ for $t \in [t_0, T_* + t_0]$. Finally, we get a contradiction with the definition of t_0 and we infer that $t_0 = T$. This completes the proof of this proposition. \Box

4. Global-In-Time Existence of a Mild Solution

Proposition 4.1 (local-in-time existence). Let $v_0 \in L^2(\mathbb{R})$. There exists $T_* > 0$ that only depends on $\|v_0\|_{L^2(\mathbb{R})}$ and $\|u_{\phi}\|_{C_b^1(\mathbb{R})}$ such that (1.6) admits a unique mild solution $v \in C([0, T_*]; L^2(\mathbb{R})) \cap C((0, T_*]; H^1(\mathbb{R}))$. Moreover, v satisfies

$$\sup_{t \in (0,T_*]} t^{1/2} \|\partial_x v(t, \cdot)\|_{L^2(\mathbb{R})} < +\infty.$$
(4.1)

Proof. For $v \in C([0,T]; L^2(\mathbb{R})) \cap C((0,T]; H^1(\mathbb{R}))$, we consider the following norm:

$$\||v\|\| := \|v\|_{C([0,T];L^2(\mathbb{R}))} + \sup_{t \in (0,T]} t^{1/2} \|\partial_x v(t,\cdot)\|_{L^2(\mathbb{R})},$$
(4.2)

and we define the affine space

$$X := \left\{ v \in C\left([0,T]; L^{2}(\mathbb{R})\right) \cap C\left((0,T]; H^{1}(\mathbb{R})\right) \text{ s.t. } v(0,\cdot) = v_{0}, |||v||| < +\infty \right\}.$$
(4.3)

It is readily seen that X endowed with the distance induced by the norm $||| \cdot |||$ is a complete metric space. For $v \in X$, we define the function

$$\Theta v: t \in [0,T] \longrightarrow K(t,\cdot) * v_0 - \frac{1}{2} \int_0^t \partial_x K(t-s,\cdot) * v^2(s,\cdot) ds - \int_0^t \partial_x K(t-s,\cdot) * u_\phi v(s,\cdot) ds.$$

$$(4.4)$$

From Proposition 2.9, $\Theta v \in C([0,T]; L^2(\mathbb{R}))$ and satisfies $\Theta v(0, \cdot) = v_0$.

Step 1 ($\Theta v \in X$). Since

$$\partial_{x}(K(t,\cdot)*v_{0}) = \partial_{x}K(t,\cdot)*v_{0} = \mathcal{F}^{-1}\Big(\xi \longmapsto 2i\pi\xi e^{-t\psi_{\mathcal{I}}(\xi)}\mathcal{F}v_{0}(\xi)\Big), \tag{4.5}$$

the dominated convergence theorem implies that for any $t_0 > 0$,

$$\int_{\mathbb{R}} 4\pi^2 |\xi|^2 \left| e^{-t\psi_{\mathcal{I}}(\xi)} - e^{-t_0\psi_{\mathcal{I}}(\xi)} \right|^2 |\mathcal{F}v_0(\xi)|^2 d\xi \longrightarrow 0, \quad \text{as } t \longrightarrow t_0.$$

$$\tag{4.6}$$

Therefore, the function $t > 0 \rightarrow (\xi \mapsto 2i\pi\xi e^{-t\psi_2(\xi)} \mathcal{F}v_0(\xi)) \in L^2(\mathbb{R})$ is continuous and since \mathcal{F} is an isometry of L^2 , we deduce that $t > 0 \rightarrow \partial_x K(t, \cdot) * v_0 \in L^2(\mathbb{R})$ is continuous. We have then established that $t > 0 \rightarrow K(t, \cdot) * v_0 \in H^1(\mathbb{R})$ is continuous. Moreover, from Proposition 2.2, we have

$$\|\partial_x K(t,\cdot) * v_0\|_{L^2(\mathbb{R})} \le K_1 t^{-1/2} \|v_0\|_{L^2(\mathbb{R})}.$$
(4.7)

Let w denote the function

$$w(t,\cdot) = \frac{1}{2} \int_0^t \partial_x K(t-s,\cdot) * v^2(s,\cdot) ds + \int_0^t \partial_x K(t-s,\cdot) * u_{\phi} v(s,\cdot) ds.$$
(4.8)

Let us now prove that $w \in C((0,T]; H^1(\mathbb{R}))$. We first have

$$\partial_x w(t,\cdot) = \int_0^t \partial_x K(t-s,\cdot) * v \partial_x v(s,\cdot) ds + \int_0^t \partial_x K(t-s,\cdot) * \partial_x (u_\phi v)(s,\cdot) ds.$$
(4.9)

Using Young inequalities and Proposition 2.2, we get

$$\begin{split} \|\partial_{x}w(t,\cdot)\|_{L^{2}(\mathbb{R})} &\leq \int_{0}^{t} \|\partial_{x}K(t-s,\cdot) * v\partial_{x}v(s,\cdot)\|_{L^{2}(\mathbb{R})} \, ds \\ &+ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot) * \partial_{x}(u_{\phi}v)(s,\cdot)\|_{L^{2}(\mathbb{R})} \, ds \\ &\leq \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{2}(\mathbb{R})} \|v\partial_{x}v(s,\cdot)\|_{L^{1}(\mathbb{R})} \, ds \\ &+ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}(\mathbb{R})} \|\partial_{x}(u_{\phi}v)(s,\cdot)\|_{L^{2}(\mathbb{R})} \, ds \\ &\leq \|v\|_{C([0,T];L^{2}(\mathbb{R}))} \int_{0}^{t} K_{0}(t-s)^{-3/4} s^{-1/2} ds \sup_{s\in(0,T]} s^{1/2} \|\partial_{x}v(s,\cdot)\|_{L^{2}(\mathbb{R})} \\ &+ \int_{0}^{t} K_{1}(t-s)^{-1/2} s^{-1/2} ds \sup_{s\in(0,T]} s^{1/2} \|\partial_{x}(u_{\phi}v)(s,\cdot)\|_{L^{2}(\mathbb{R})}. \end{split}$$

We then obtain

$$\begin{aligned} \|\partial_{x}w(t,\cdot)\|_{L^{2}(\mathbb{R})} &\leq K_{0}I\|v\|_{C([0,T];L^{2}(\mathbb{R}))}T^{-1/4}\sup_{s\in(0,T]}s^{1/2}\|\partial_{x}v(s,\cdot)\|_{L^{2}(\mathbb{R})} \\ &+ K_{1}J\sup_{s\in(0,T]}s^{1/2}\|\partial_{x}(u_{\phi}v)(s,\cdot)\|_{L^{2}(\mathbb{R})}, \end{aligned}$$

$$(4.11)$$

where I = B(1/2, 1/4) and $J = B(1/2, 1/2) = \pi$, *B* being the beta function defined by

$$B(x,y) := \int_0^1 t^{x-1} (1-t)^{y-1} dt.$$
(4.12)

As $|||v||| < \infty$ then

$$\sup_{s \in (0,T]} s^{1/2} \|\partial_x v(s, \cdot)\|_{L^2(\mathbb{R})} < \infty, \qquad \sup_{s \in (0,T]} s^{1/2} \|\partial_x (u_{\phi} v)(s, \cdot)\|_{L^2(\mathbb{R})} < \infty.$$
(4.13)

We then deduce that $\partial_x w(t, \cdot)$ is in L^2 and so $\partial_x \Theta v(t, \cdot) \in L^2(\mathbb{R})$ for all $t \in (0, T]$.

Let us now prove that $\partial_x w$ is continuous on (0, T] with values in L^2 . For $\delta > 0$ and $t \in (0, T]$, we define

$$(\partial_x w)_{\delta}(t, \cdot) := \int_0^t \partial_x K(t-s, \cdot) * \mathbf{1}_{\{s>\delta\}}(v\partial_x v)(s, \cdot) ds + \int_0^t \partial_x K(t-s, \cdot) * \mathbf{1}_{\{s>\delta\}} \partial_x (u_{\phi} v)(s, \cdot) ds.$$

$$(4.14)$$

Since $1_{\{s>\delta\}}(v\partial_x v)(s,\cdot) \in L^{\infty}([0,T];L^1(\mathbb{R}))$ and $1_{\{s>\delta\}}\partial_x(u_{\phi}v)(s,\cdot) \in L^{\infty}([0,T];L^2(\mathbb{R}))$ then Proposition 2.9 implies that $(\partial_x w)_{\delta} : [0,T] \to L^2(\mathbb{R})$ is continuous. Moreover, we have for any $t \in (0,T]$ and $\delta \leq t$,

$$\begin{aligned} \|\partial_{x}w(t,\cdot) - (\partial_{x}w)_{\delta}(t,\cdot)\|_{L^{2}} &\leq K_{0} \int_{0}^{\delta} (t-s)^{-3/4} s^{-1/2} \, ds \|v\|_{C([0,T];L^{2})} \sup_{s\in(0,T]} s^{1/2} \|\partial_{x}v(s,\cdot)\|_{L^{2}} \\ &+ K_{1} \int_{0}^{\delta} (t-s)^{-1/2} s^{-1/2} ds \sup_{s\in(0,T]} s^{1/2} \|\partial_{x}(u_{\phi}v)(s,\cdot)\|_{L^{2}}. \end{aligned}$$

$$(4.15)$$

It then follows that

$$\sup_{t \in (0,T]} \|\partial_x w(t,\cdot) - (\partial_x w)_{\delta}(t,\cdot)\|_{L^2(\mathbb{R})} \longrightarrow 0 \quad \text{as } \delta \longrightarrow 0.$$
(4.16)

We next infer that $\partial_x w \in C((0,T]; L^2(\mathbb{R}))$ because it is a local uniform limit of continuous functions. Hence, we have established that $\Theta v \in C([0,T]; L^2(\mathbb{R})) \cap C((0,T]; H^1(\mathbb{R}))$. To prove that $\Theta v \in X$, it remains to show that $||\Theta v|| < +\infty$. Using (4.7) and (4.11), we have

$$\sup_{t \in (0,T]} t^{1/2} \|\partial_x \Theta v(t, \cdot)\|_{L^2} \le K_1 \|v_0\|_{L^2} + K_0 I T^{1/4} \sup_{s \in (0,T]} s^{1/2} \|\partial_x v(s, \cdot)\|_{L^2} \|v\|_{C([0,T];L^2)} + K_1 J T^{1/2} \sup_{s \in (0,T]} s^{1/2} \|\partial_x (u_{\phi} v)(s, \cdot)\|_{L^2}.$$

$$(4.17)$$

Finally, we have $\Theta : X \to X$.

Step 2. We begin by considering a ball of X of radius R centered at the origin

$$B_R := \{ v \in X / |||v||| \le R \}, \tag{4.18}$$

where $R > ||v_0||_{L^2(\mathbb{R})} + K_1 ||v_0||_{L^2(\mathbb{R})}$. Take $v \in B_R$ and let us now prove that Θ maps B_R into itself. We have

$$\|\Theta v(t,\cdot)\|_{L^{2}(\mathbb{R})} \leq \|K(t,\cdot) * v_{0}\|_{L^{2}(\mathbb{R})} + \int_{0}^{t} \left\|\partial_{x}K(t-s,\cdot) * \left(\frac{v^{2}}{2} + u_{\phi}v\right)(s,\cdot)\right\|_{L^{2}(\mathbb{R})} ds.$$
(4.19)

By Remark 2.4, we get

$$\|K(t,\cdot) * v_0\|_{L^2(\mathbb{R})} \le e^{\alpha_0 T} \|v_0\|_{L^2(\mathbb{R})},$$
(4.20)

where $\alpha_0 = -\min \operatorname{Re}(\psi_2) > 0$. Moreover, since $||v^2||_{L^{\infty}((0,T);L^1(\mathbb{R}))} = ||v||^2_{L^{\infty}((0,T);L^2(\mathbb{R}))}$ and with the help of Proposition 2.2, we get

$$\|\Theta v(t,\cdot)\|_{L^{2}(\mathbb{R})} \leq e^{\alpha_{0}T} \|v_{0}\|_{L^{2}(\mathbb{R})} + 2K_{0}T^{1/4}R^{2} + 2K_{1}T^{1/2} \|u_{\phi}\|_{C^{1}_{b}(\mathbb{R})}R.$$
(4.21)

Using (4.17) and (4.21), we deduce that

$$\begin{split} |||\Theta v||| &\leq e^{\alpha_0 T} ||v_0||_{L^2(\mathbb{R})} + K_1 ||v_0||_{L^2(\mathbb{R})} + (2+I)K_0 T^{1/4} R^2 + (2+J)RK_1 T^{1/2} ||u_{\phi}||_{C_b^1(\mathbb{R})} \\ &+ K_1 J ||u_{\phi}||_{C_b^1(\mathbb{R})} RT. \end{split}$$
(4.22)

Therefore, for T > 0 sufficiently small such that

$$e^{a_0 T} \|v_0\|_{L^2(\mathbb{R})} + K_1 \|v_0\|_{L^2(\mathbb{R})} + (2+I)K_0 T^{1/4} R^2 + (2+J)RK_1 T^{1/2} \|u_{\phi}\|_{C_b^1(\mathbb{R})} + K_1 J \|u_{\phi}\|_{C_b^1(\mathbb{R})} RT \le R,$$
(4.23)

we get that $|||\Theta v||| \le R$.

To finish with, we are going to prove that Θ is a contraction. For $v, w \in B_R$, we have for any $t \in (0, T)$

$$\begin{split} \|\Theta v(t,\cdot) - \Theta w(t,\cdot)\|_{L^{2}(\mathbb{R})} &\leq \frac{1}{2} \int_{0}^{t} \|\partial_{x} K(t-s,\cdot)\|_{L^{2}(\mathbb{R})} \left\| \left(v^{2} - w^{2} \right)(s,\cdot) \right\|_{L^{1}(\mathbb{R})} ds \\ &+ \int_{0}^{t} \|\partial_{x} K(t-s,\cdot)\|_{L^{1}(\mathbb{R})} \left\| u_{\phi}(v-w)(s,\cdot) \right\|_{L^{2}(\mathbb{R})} ds \\ &\leq 2K_{0} t^{1/4} \left\| v^{2} - w^{2} \right\|_{C([0,T];L^{1}(\mathbb{R}))} \\ &+ 2K_{1} t^{1/2} \left\| u_{\phi} \right\|_{C^{1}_{b}(\mathbb{R})} \|v-w\|_{C([0,T];L^{2}(\mathbb{R}))}, \end{split}$$
(4.24)

and since,

$$\left\| v^{2} - w^{2} \right\|_{C([0,T];L^{1}(\mathbb{R}))} \leq \left(\| v \|_{C([0,T];L^{2}(\mathbb{R}))} + \| w \|_{C([0,T];L^{2}(\mathbb{R}))} \right) \| v - w \|_{C([0,T];L^{2}(\mathbb{R}))}$$

$$\leq 2R \| v - w \|_{C([0,T];L^{2}(\mathbb{R}))},$$

$$(4.25)$$

we get

$$\|\Theta v(t,\cdot) - \Theta w(t,\cdot)\|_{L^{2}(\mathbb{R})} \leq \left(4RK_{0}t^{1/4} + 2K_{1}t^{1/2} \|u_{\phi}\|_{C_{b}^{1}(\mathbb{R})}\right)\|v - w\|_{C([0,T];L^{2}(\mathbb{R}))}.$$
(4.26)

Moreover

$$\begin{aligned} \|\partial_{x}(\Theta v - \Theta w)(t, \cdot)\|_{L^{2}(\mathbb{R})} &\leq \frac{1}{2} \int_{0}^{t} \left\|\partial_{x}K(t - s, \cdot) * \partial_{x}\left(v^{2} - w^{2}\right)(s, \cdot)\right\|_{L^{2}(\mathbb{R})} ds \\ &+ \int_{0}^{t} \left\|\partial_{x}K(t - s, \cdot) * \partial_{x}\left(u_{\phi}(v - w)\right)(s, \cdot)\right\|_{L^{2}(\mathbb{R})} ds \\ &\leq K_{0}It^{-1/4} \sup_{s \in (0,T]} s^{1/2} \left\|(v\partial_{x}v - w\partial_{x}w)(s, \cdot)\right\|_{L^{1}(\mathbb{R})} \\ &+ K_{1}J \sup_{s \in (0,T]} s^{1/2} \left\|\partial_{x}\left(u_{\phi}(v - w)\right)(s, \cdot)\right\|_{L^{2}(\mathbb{R})}. \end{aligned}$$

$$(4.27)$$

And since

$$\|(v\partial_{x}v - w\partial_{x}w)(t, \cdot)\|_{L^{1}} \leq \|\partial_{x}w(t, \cdot)\|_{L^{2}}\|(v - w)(t, \cdot)\|_{L^{2}} + \|v(t, \cdot)\|_{L^{2}}\|\partial_{x}(v - w)(t, \cdot)\|_{L^{2}},$$
(4.28)

then

$$t^{1/2} \| (v\partial_x v - w\partial_x w)(t, \cdot) \|_{L^1} \le \| (v - w)(t, \cdot) \|_{L^2} \| w \| \| + \| v \| |t^{1/2} \| \partial_x (v - w)(t, \cdot) \|_{L^2}$$

$$\le 2R \| \| v - w \| \|.$$
(4.29)

Therefore, we obtain

$$\begin{aligned} \|\partial_{x}(\Theta v - \Theta w)(t, \cdot)\|_{L^{2}(\mathbb{R})} &\leq 2K_{0}It^{-1/4}R|\|v - w\|| + K_{1}J\|u_{\phi}\|_{C_{b}^{1}(\mathbb{R})}T^{1/2}|\|v - w\|| \\ &+ K_{1}J\|u_{\phi}\|_{C_{b}^{1}(\mathbb{R})}|\|v - w\||. \end{aligned}$$

$$(4.30)$$

Finally, using (4.26) and (4.30), we get

$$|||\Theta v - \Theta w||| \le \left[(2+I)2RK_0 T^{1/4} + (2+J) \| u_\phi \|_{C_b^1(\mathbb{R})} K_1 T^{1/2} + K_1 J T \| u_\phi \|_{C_b^1(\mathbb{R})} \right] |||v - w|||.$$
(4.31)

Step 3 (conclusion). For any $T_* > 0$ sufficiently small such that (4.23) holds true and

$$(2+I)2RK_0T_*^{1/4} + (2+J) \|u_{\phi}\|_{C_b^1(\mathbb{R})} K_1T_*^{1/2} + K_1JT_* \|u_{\phi}\|_{C_b^1(\mathbb{R})} < 1,$$
(4.32)

 Θ is a contraction from B_R into itself. The Banach fixed point theorem then implies that Θ admits a unique fixed point $v \in C([0, T_*]; L^2(\mathbb{R})) \cap C((0, T_*]; H^1(\mathbb{R}))$ which is a mild solution to (1.6).

Lemma 4.2 (regularity H^2 of $v(t, \cdot)$). Let $v_0 \in L^2(\mathbb{R})$ and $\phi \in C_b^2(\mathbb{R})$. There exists $T'_* > 0$ that only depends on $\|v_0\|_{L^2(\mathbb{R})}$ and $\|u_{\phi}\|_{C_b^2(\mathbb{R})}$ such that (1.6) admits a unique mild solution $v \in C([0, T'_*]; L^2(\mathbb{R})) \cap C((0, T'_*]; H^2(\mathbb{R}))$. Moreover, v satisfies

$$\sup_{t \in (0,T'_{*}]} t^{1/2} \|\partial_{x} v(t,\cdot)\|_{L^{2}(\mathbb{R})} < +\infty, \qquad \sup_{t \in (0,T'_{*}]} t \|\partial^{2}_{xx} v(t,\cdot)\|_{L^{2}(\mathbb{R})} < +\infty.$$
(4.33)

Proof. To prove this result, we use again a contracting fixed point theorem. But this time, it is the gradient of the solution *v* which is searched as a fixed point.

From Proposition 4.1, there exists $T_* > 0$ which depends on $\|v_0\|_{L^2(\mathbb{R})}$ and $\|u_{\phi}\|_{C^1(\mathbb{R})}$ such that $v \in C([0,T_*];L^2(\mathbb{R})) \cap C((0,T_*];H^1(\mathbb{R}))$ is a mild solution to (1.6). Since $v \in C((0,T_*];H^1(\mathbb{R}))$, we can consider the gradient of $v(t, \cdot)$ for any $t \in (0,T_*]$. Let then $t_0 \in (0,T_*)$ and $T'_* \in (0,T_*-t_0]$. We consider the same complete metric space *X* defined in the proof of Proposition 4.1 and we take the norm $||| \cdot |||$ defined in (4.2):

$$X := \left\{ w \in C\left([0, T'_*]; L^2(\mathbb{R}) \right) \cap C\left((0, T'_*]; H^1(\mathbb{R}) \right) \text{ s.t. } w(0, \cdot) = w_0, \ |||w||| < +\infty \right\},$$
(4.34)

with the initial data $w_0 = \partial_x v(t_0, \cdot)$.

We now wish to apply the fixed point theorem at the following function:

$$\Theta w : t \in [0, T'_{*}] \longrightarrow K(t, \cdot) * w_{0} - \int_{0}^{t} \partial_{x} K(t - s, \cdot) * (\overline{v}w)(s, \cdot) ds$$
$$- \int_{0}^{t} \partial_{x} K(t - s, \cdot) * (\partial_{x}(u_{\phi})\overline{v})(s, \cdot) ds$$
$$- \int_{0}^{t} \partial_{x} K(t - s, \cdot) * (u_{\phi}w)(s, \cdot) ds, \qquad (4.35)$$

where $\overline{v}(t, \cdot) := v(t_0 + t, \cdot)$. First, we leave the reader to verify that Θ maps X into itself. The proof is similar to the one given in Proposition 4.1.

For any $w \in X$, we have from Young inequalities and Remark 2.4

$$\begin{split} \|\Theta w(t,\cdot)\|_{L^{2}(\mathbb{R})} &\leq e^{\alpha_{0}T'_{*}} \|w_{0}\|_{L^{2}(\mathbb{R})} + \|\overline{v}\|_{C([t_{0},T'_{*}];H^{1}(\mathbb{R}))} \|\|w\|\| \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{2}(\mathbb{R})} ds \\ &+ \|u_{\phi}\|_{C^{1}_{b}(\mathbb{R})} \|\overline{v}\|_{C([t_{0},T'_{*}];H^{1}(\mathbb{R}))} \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}(\mathbb{R})} ds \\ &+ \|u_{\phi}\|_{C^{1}_{b}(\mathbb{R})} \|w\|\| \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}(\mathbb{R})} ds, \end{split}$$
(4.36)

and from Proposition 2.2, we get

$$\begin{aligned} \|\Theta w(t,\cdot)\|_{L^{2}} &\leq e^{\alpha_{0}T_{*}^{\prime}} \|w_{0}\|_{L^{2}} + 4K_{0}T_{*}^{\prime 1/4} \|\overline{v}\|_{C([t_{0},T_{*}^{\prime}];H^{1})} \|w\|| \\ &+ 2K_{1}T_{*}^{\prime 1/2} \|u_{\phi}\|_{C_{b}^{1}} \|\overline{v}\|_{C([t_{0},T_{*}^{\prime}];H^{1})} + 2K_{1}T_{*}^{\prime 1/2} \|u_{\phi}\|_{C_{b}^{1}} \|w\||. \end{aligned}$$

$$(4.37)$$

Differentiating $\Theta v(t, \cdot)$ with respect to the space variable, we obtain

$$\partial_{x}\Theta v(t,\cdot) = \partial_{x}K(t,\cdot) * w_{0} - \int_{0}^{t} \partial_{x}K(t-s,\cdot) * \partial_{x}(\overline{v}w)(s,\cdot)ds - \int_{0}^{t} \partial_{x}K(t-s,\cdot) * \partial_{x}(\partial_{x}(u_{\phi})\overline{v})(s,\cdot)ds - \int_{0}^{t} \partial_{x}K(t-s,\cdot) * \partial_{x}(u_{\phi}w)(s,\cdot)ds,$$

$$(4.38)$$

and developing, we get

$$\partial_{x}\Theta v(t,\cdot) = \partial_{x}K(t,\cdot) * w_{0} - \int_{0}^{t} \partial_{x}K(t-s,\cdot) * [w\partial_{x}\overline{v} + \overline{v}\partial_{x}w](s,\cdot)ds$$
$$-\int_{0}^{t} \partial_{x}K(t-s,\cdot) * \left[\partial_{x}^{2}(u_{\phi})\overline{v} + \partial_{x}(u_{\phi})\partial_{x}\overline{v}\right](s,\cdot)ds \qquad (4.39)$$
$$-\int_{0}^{t} \partial_{x}K(t-s,\cdot) * \left[\partial_{x}(u_{\phi})w + u_{\phi}\partial_{x}w\right](s,\cdot)ds.$$

Now, from Young inequalities, we have

$$\begin{aligned} \|\partial_{x}\Theta v(t,\cdot)\|_{L^{2}} &\leq \|\partial_{x}K(t,\cdot)\|_{L^{1}}\|w_{0}\|_{L^{2}} + \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{2}}\|w\partial_{x}\overline{v}(s,\cdot)\|_{L^{1}}ds \\ &+ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{2}}\|\overline{v}\partial_{x}w(s,\cdot)\|_{L^{1}}ds \\ &+ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}}\Big[\left\|\partial_{x}^{2}(u_{\phi})\overline{v}(s,\cdot)\right\|_{L^{2}} + \left\|\partial_{x}(u_{\phi})\partial_{x}\overline{v}(s,\cdot)\right\|_{L^{2}}\Big]ds \\ &+ \int_{0}^{t} \|\partial_{x}K(t-s,\cdot)\|_{L^{1}}\Big[\left\|\partial_{x}(u_{\phi})w(s,\cdot)\right\|_{L^{2}} + \left\|u_{\phi}\partial_{x}w(s,\cdot)\right\|_{L^{2}}\Big]ds. \end{aligned}$$
(4.40)

Finally, from Proposition 2.2, we obtain

$$\begin{split} \|\partial_{x}\Theta v(t,\cdot)\|_{L^{2}} &\leq t^{-1/2}K_{1}\|w_{0}\|_{L^{2}} + 4t^{1/4}K_{0}\|\overline{v}\|_{C([t_{0};T'_{*}];H^{1})}\|\|w\|\| \\ &+ \int_{0}^{t}K_{0}(t-s)^{-3/4}s^{-1/2}ds\|\overline{v}\|_{C([t_{0};T'_{*}];H^{1})}\sup_{s\in(0,T'_{*}]}s^{1/2}\|\partial_{x}w(s,\cdot)\|_{L^{2}} \\ &+ 4K_{1}t^{1/2}\|u_{\phi}\|_{C^{2}_{b}}\|\overline{v}\|_{C([t_{0};T'_{*}];H^{1})} + 2K_{1}t^{1/2}\|u_{\phi}\|_{C^{1}_{b}}\|\|w\|\| \\ &+ \int_{0}^{t}K_{1}(t-s)^{-1/2}s^{-1/2}ds\|u_{\phi}\|_{C^{2}_{b}}\sup_{s\in(0,T'_{*}]}s^{1/2}\|\partial_{x}w(s,\cdot)\|_{L^{2}}. \end{split}$$

$$(4.41)$$

In other words, we have for all $t \in (0, T_*']$

$$t^{1/2} \|\partial_{x} \Theta v(t, \cdot)\|_{L^{2}} \leq K_{1} \|w_{0}\|_{L^{2}} + 4 T_{*}^{'3/4} K_{0} \|\overline{v}\|_{C([t_{0};T_{*}'];H^{1})} \|\|w\|\| + K_{0} I T_{*}^{'1/4} \|\overline{v}\|_{C([t_{0};T_{*}'];H^{1})} \|\|w\|\| + 4K_{1} T_{*}^{'} \|u_{\phi}\|_{C_{b}^{2}} \|\overline{v}\|_{C([t_{0};T_{*}'];H^{1})}$$

$$+ 2K_{1} \pi T_{*}^{'} \|8_{\phi}\|_{C_{b}^{1}} \|\|w\|\| + K_{1} T_{*}^{'1/2} \|u_{\phi}\|_{C_{b}^{2}} \|\|w\|\|,$$

$$(4.42)$$

where I = B(1/2, 1/4). Hence, using (4.37) and (4.42), we get

$$\begin{split} |||\Theta w||| &\leq e^{\alpha_0 T'_*} ||w_0||_{L^2(\mathbb{R})} + K_1 ||w_0||_{L^2(\mathbb{R})} + 2K_1 ||u_\phi||_{C^2_b(\mathbb{R})} ||\overline{v}||_{C([t_0;T'_*];H^1(\mathbb{R}))} \left(2T'_* + T'^{1/2}_*\right) \\ &+ C ||w|| \left(T'^{1/4}_* + T'^{1/2}_* + T'^{3/4}_* + T'_*\right), \end{split}$$

$$(4.43)$$

for some positive constant *C* which depends on $K_0, K_1, \|\overline{v}\|_{C([t_0;T'_*];H^1(\mathbb{R}))}$ and $\|u_{\phi}\|_{C^2_{\iota}(\mathbb{R})}$. We next leave reader to verify that: for any $w_1, w_2 \in X$,

$$\||\Theta w_1 - \Theta w_2|\| \le C' \left(T_*^{'1/4} + T_*^{'1/2} + T_*^{'3/4} + T_*' \right) |\|w_1 - w_2\||, \tag{4.44}$$

,

where C' is a positive constant which depends on $K_0, K_1, \|\overline{v}\|_{C([t_0;T'_*];H^1(\mathbb{R}))}$ and $\|u_{\phi}\|_{C^2_{L}(\mathbb{R})}$. Then, if $T'_* > 0$ satisfies

$$e^{\alpha_{0}T_{*}'}\|w_{0}\|_{L^{2}(\mathbb{R})} + K_{1}\|w_{0}\|_{L^{2}(\mathbb{R})} + 2K_{1}\|u_{\phi}\|_{C_{b}^{1}(\mathbb{R})}\|\overline{v}\|_{C([t_{0};T'];H^{1}(\mathbb{R}))}\left(2T_{*}' + T_{*}^{'1/2}\right) + CR\left(T_{*}^{'1/4} + T_{*}^{'1/2} + T_{*}^{'3/4} + T_{*}'\right) \leq R,$$

$$C'\left(T_{*}^{'1/4} + T_{*}^{'1/2} + T_{*}^{'3/4} + T_{*}\right) < 1,$$

$$(4.45)$$

 Θ : $B_R(X) \rightarrow B_R(X)$ is a contraction, where $B_R(X)$ is ball of X of radius R centered at the origin. Using a contracting point fixed theorem, there exists a unique fixed point, which we denote by w. But it is easy to see that $\Theta \partial_x \overline{v} = \partial_x \overline{v}$ taking into account the space derived from the Duhamel formulation (2.1). Thanks to a uniqueness argument, we deduce that $w = \partial_x \overline{v}$ and thus that $v \in C((0, T'_*]; H^2(\mathbb{R}))$, which completes the proof of this lemma.

Let us now prove the global-in-time existence of mild solution *v*.

Proposition 4.3 (global-in-time existence). Let $v_0 \in L^2(\mathbb{R})$, $\phi \in C_b^2(\mathbb{R})$ and T > 0. Then, there exists a (unique) mild solution $v \in C([0,T]; L^2(\mathbb{R})) \cap C((0,T]; H^2(\mathbb{R}))$ to (1.6). Moreover, v satisfies the PDE (1.6) in the distribution sense.

Proof.

Step 1 (*v* is a distribution solution). Taking the Fourier transform with respect to the space variable in (2.1), we get for all $t \in [0, T]$ and all $\xi \in \mathbb{R}$,

$$\begin{aligned} \boldsymbol{\mathcal{F}}(\boldsymbol{\upsilon}(t,\cdot))(\boldsymbol{\xi}) &= e^{-t\psi_{\mathcal{I}}(\boldsymbol{\xi})} \boldsymbol{\mathcal{F}}\boldsymbol{\upsilon}_{0}(\boldsymbol{\xi}) - \int_{0}^{t} i\pi\boldsymbol{\xi}e^{-(t-s)\psi_{\mathcal{I}}(\boldsymbol{\xi})} \boldsymbol{\mathcal{F}}\left(\boldsymbol{\upsilon}^{2}(s,\cdot)\right)(\boldsymbol{\xi})ds \\ &- \int_{0}^{t} 2i\pi\boldsymbol{\xi}e^{-(t-s)\psi_{\mathcal{I}}(\boldsymbol{\xi})} \boldsymbol{\mathcal{F}}\left(\boldsymbol{u}_{\phi}\boldsymbol{\upsilon}(s,\cdot)\right)(\boldsymbol{\xi})ds. \end{aligned}$$

$$(4.46)$$

Define

$$G(t,\xi) = -\int_0^t 2i\pi\xi e^{-(t-s)\psi_0(\xi)} \mathcal{F}\left(\frac{v^2}{2} + u_\phi v\right)(s,\cdot)(\xi) ds.$$

$$(4.47)$$

Classical results on ODE imply that G is differentiable with respect to the time with

$$\partial_{t}G(t,\xi) + \psi_{\mathcal{I}}(\xi)G(t,\xi) = -i\pi\xi \mathcal{F}\left(v^{2}(t,\cdot)\right)(\xi) - 2i\pi\xi \mathcal{F}\left(\left(u_{\phi}v\right)(t,\cdot)\right)(\xi)$$

$$= -\mathcal{F}\left(\partial_{x}\left(\frac{v^{2}}{2}\right)(t,\cdot)\right)(\xi) - \mathcal{F}\left(\partial_{x}\left(u_{\phi}v\right)(t,\cdot)\right)(\xi).$$
(4.48)

Let us now prove that all terms in (4.48) are continuous with values in L^2 . Since, $v \in C((0,T]; H^1(\mathbb{R}))$ then $\partial_x(v^2)$, $\partial_x(u_\phi v) \in C((0,T]; L^2(\mathbb{R}))$. We thus deduce that $\mathcal{F}(\partial_x(v^2/2))$ and $\mathcal{F}(\partial_x(u_\phi v))$ are continuous with values in $L^2(\mathbb{R})$. Moreover, (4.46) implies that

$$\psi_{\mathcal{D}}G(t,\cdot) = \psi_{\mathcal{D}}(\mathcal{F}(v(t,\cdot)) - e^{-t\psi_{\mathcal{D}}}\mathcal{F}v_0), \qquad (4.49)$$

and so $\psi_{\mathcal{D}}G(t, \cdot)$ is continuous with values in L^2 . Indeed,

$$\begin{split} \int_{\mathbb{R}} \left| \psi_{\mathcal{I}}(\xi) G(t,\xi) \right|^2 d\xi &= \int_{-1}^{1} \left| \psi_{\mathcal{I}}(\xi) G(t,\xi) \right|^2 d\xi + \int_{\mathbb{R} \setminus (-1,1)} \left| \psi(\xi) G(t,\xi) \right|^2 d\xi \\ &\leq \sup_{\xi \in [-1,1]} \left| \psi_{\mathcal{I}}(\xi) \right|^2 \| G(t,\cdot) \|_{L^2(\mathbb{R})}^2 + C \int_{\mathbb{R} \setminus (-1,1)} \left| \xi^2 G(t,\xi) \right|^2 d\xi \\ &\leq \sup_{\xi \in [-1,1]} \left| \psi_{\mathcal{I}}(\xi) \right|^2 \| G(t,\cdot) \|_{L^2(\mathbb{R})}^2 \\ &+ C \int_{\mathbb{R} \setminus (-1,1)} \left| \mathcal{F} \left(\partial_{xx}^2 v(t,\cdot) \right) - \xi^2 e^{-t\psi_{\mathcal{I}}(\xi)} \mathcal{F} v_0 \right|^2 d\xi \end{split}$$

$$\leq \sup_{\xi \in [-1,1]} |\psi_{\mathcal{I}}(\xi)|^{2} ||G(t,\cdot)||_{L^{2}(\mathbb{R})}^{2} + \widetilde{C} ||v(t,\cdot)||_{H^{2}(\mathbb{R})}^{2} + \widetilde{C} ||v_{0}||_{L^{2}(\mathbb{R})}^{2} + C ||v(t,\cdot)||_{H^{2}} ||v_{0}||_{L^{2}} < \infty,$$
(4.50)

because $\psi_{\mathcal{D}}$ behaves at infinity as $|\cdot|^2$. *C*, \tilde{C} are two positive constants. Hence, we have that the function $t \to \psi_{\mathcal{D}}G(t, \cdot) \in L^2(\mathbb{R}, \mathbb{C})$ is continuous. Finally, we have proved that all the terms in (4.48) are continuous with values in L^2 . Therefore, from (4.48), we get that $G \in C^1([0,T]; L^2(\mathbb{R}, \mathbb{C}))$ and then

$$\frac{d}{dt}(G(t,\cdot)) + \psi_{\mathcal{I}}G(t,\cdot) = -\mathcal{F}\left(\partial_x \left(\frac{v^2}{2}\right)(t,\cdot)\right) - \mathcal{F}\left(\partial_x \left(u_{\phi}v\right)(t,\cdot)\right). \tag{4.51}$$

Moreover, $t \in [0, T] \rightarrow e^{-t\psi_{\mathcal{I}}} \mathcal{F} v_0 \in L^2(\mathbb{R}, \mathbb{C})$ is C^1 with

$$\frac{d}{dt}(e^{-t\psi_{\mathcal{I}}}\mathcal{F}v_{0}) + \psi_{\mathcal{I}}e^{-t\psi_{\mathcal{I}}}\mathcal{F}v_{0} = 0.$$
(4.52)

From (4.46), we infer that $\mathcal{F}v$ is C^1 on [0,T] with values in L^2 with

$$\frac{d}{dt}\mathcal{F}(v(t,\cdot)) = -\psi_{\mathcal{I}}\mathcal{F}(v(t,\cdot)) - \mathcal{F}\left(\partial_{x}\left(\frac{v^{2}}{2}\right)(t,\cdot)\right) - \mathcal{F}\left(\partial_{x}\left(u_{\phi}v\right)(t,\cdot)\right).$$
(4.53)

Since \mathcal{F} is an isometry of L^2 , we deduce that $v \in C^1([0,T]; L^2(\mathbb{R}))$ and by (1.3), we get

$$\frac{d}{dt}(v(t,\cdot)) = -\mathcal{F}^{-1}(\psi_{\mathcal{I}}\mathcal{F}(v(t,\cdot))) - \partial_{x}\left(\frac{v^{2}}{2}\right)(t,\cdot) - \partial_{x}\left(u_{\phi}v\right)(t,\cdot)$$

$$= -\mathcal{I}[v(t,\cdot)] + \partial_{xx}^{2}v(t,\cdot) - \partial_{x}\left(\frac{v^{2}}{2}\right)(t,\cdot) - \partial_{x}\left(u_{\phi}v\right)(t,\cdot).$$
(4.54)

We are now going to prove that v satisfies the PDE (1.6) in the distribution sense. Let us note

$$w(t,\cdot) \coloneqq -\mathcal{O}[v(t,\cdot)] + \partial_{xx}^2 v(t,\cdot) - \partial_x \left(\frac{v^2}{2}\right)(t,\cdot) - \partial_x \left(u_\phi v\right)(t,\cdot), \tag{4.55}$$

and let us show that

$$\partial_t v = w \quad \text{in } \mathfrak{D}'((0,T) \times \mathbb{R}).$$
 (4.56)

By definition, we have for any $\varphi \in \mathfrak{D}(0,T)$ and $\psi \in \mathfrak{D}(\mathbb{R})$:

$$\langle \partial_t v, \varphi \psi \rangle = -\int_0^T \int_{\mathbb{R}} v(t, x) \frac{d\varphi}{dt} \psi(x) dt \, dx$$

$$= -\int_0^T \left(\int_{\mathbb{R}} v(t, x) \psi(x) dx \right) \frac{d\varphi(t)}{dt} dt.$$

$$(4.57)$$

Therefore, it is enough to prove that

$$\int_{0}^{T} \left(\int_{\mathbb{R}} w(t, x) \psi(x) dx \right) \varphi(t) dt = -\int_{0}^{T} \left(\int_{\mathbb{R}} v(t, x) \psi(x) dx \right) \varphi'(t) dt,$$
(4.58)

that is,

$$\frac{d}{dt} \int_{\mathbb{R}} v(t, x) \psi(x) dx = \int_{\mathbb{R}} w(t, x) \psi(x) dx, \qquad (4.59)$$

in the sense of $\mathfrak{D}'(0, T)$. But by (4.54), we have that the function

$$t \in (0,T) \longmapsto \int_{\mathbb{R}} v(t,x)\psi(x)dx \in \mathbb{R}$$
(4.60)

is C^1 and

$$\frac{d}{dt} \int_{\mathbb{R}} v(t, x) \psi(x) dx = \int_{\mathbb{R}} w(t, x) \psi(x) dx$$
(4.61)

in the classical sense, which proves that the mild solution v is a distribution solution of (1.6). *Step 2* (priori estimate). By Step 1, we have

$$\partial_t v + \partial_x \left(\frac{v^2}{2} + u_{\phi} v \right) + \mathcal{O}[v] - \partial_{xx}^2 v = 0$$
(4.62)

in the distribution sense. Therefore, multiplying this equality by v and integrating with respect to the space variable, we get:

$$\int_{\mathbb{R}} v_t v \, dx + \int_{\mathbb{R}} (\mathcal{O}[v] - v_{xx}) v \, dx + \int_{\mathbb{R}} (u_\phi v)_x v \, dx = 0, \qquad (4.63)$$

because the nonlinear term is zero. Indeed, integrating by parts, we have

$$\int_{\mathbb{R}} \partial_x \left(\frac{v^2}{2}\right) v \, dx = -\int_{\mathbb{R}} \frac{v^2}{2} \partial_x v \, dx = -\frac{1}{2} \int_{\mathbb{R}} \partial_x \left(\frac{v^2}{2}\right) v \, dx. \tag{4.64}$$

There is no boundary term from the infinity because for all $t \in (0,T]$, $v(t, \cdot) \in H^2(\mathbb{R})$. Using (1.3) and the fact that $\int_{\mathbb{R}} (\mathcal{I}[v] - \partial_{xx}^2 v) v \, dx$ is real, we get

$$\int_{\mathbb{R}} \left(\mathcal{O}[v] - \partial_{xx}^2 v \right) v \, dx = \int_{\mathbb{R}} \mathcal{F}^{-1}(\psi_{\mathcal{O}} \mathcal{F} v) v \, dx = \int_{\mathbb{R}} \psi_{\mathcal{O}} |\mathcal{F} v|^2 d\xi = \int_{\mathbb{R}} \operatorname{Re}(\psi_{\mathcal{O}}) |\mathcal{F} v|^2 d\xi.$$
(4.65)

Moreover, since $u_{\phi}v \in H^1(\mathbb{R})$ we have

$$\int_{\mathbb{R}} \left(u_{\phi} v \right)_{x} v \, dx = -\int_{\mathbb{R}} u_{\phi} v v_{x} \, dx = -\int_{\mathbb{R}} u_{\phi} \left(\frac{v^{2}}{2} \right)_{x} dx = \int_{\mathbb{R}} \left(\partial_{x} u_{\phi} \right) \frac{v^{2}}{2} \, dx. \tag{4.66}$$

Using (4.63), (4.65), and (4.66), we obtain

$$\frac{1}{2}\frac{d}{dt}\|v(t,\cdot)\|_{L^2}^2 \le (\alpha_0 + C_{\phi})\|v(t,\cdot)\|_{L^2}^2, \tag{4.67}$$

where $\alpha_0 = -\min \operatorname{Re}(\psi_2) > 0$ and $C_{\phi} = (1/2)||u_{\phi}||_{C_b^1}$. Finally, we get for all $t \in [0,T]$ the following estimate:

$$\|v(t,\cdot)\|_{L^2(\mathbb{R})} \le e^{(\alpha_0 + C_{\phi})t} \|v_0\|_{L^2(\mathbb{R})}.$$
(4.68)

Step 3 (global-in-time existence). Up to this point, we know thanks to Proposition 4.1 and Lemma 4.2 that there exists $T_* = T_*(||v_0||_{L^2(\mathbb{R})}, ||u_{\phi}||_{C^2_b(\mathbb{R})}) > 0$ such that $v \in C([0, T_*]; L^2(\mathbb{R})) \cap C((0, T_*]; H^2(\mathbb{R}))$ is a mild solution of (1.6) on $(0, T_*]$. Let us define

$$t_0 := \sup\{t > 0/\text{there exists a mild solution of (1.6) on (0, t) with initial condition } v_0\}.$$
(4.69)

To prove the global-in-time existence of a mild solution, we have to prove that $t_0 \ge T$, where T is any positive constant. Assume by contradiction that $t_0 < T$. With again the help of Proposition 4.1, there exists $T'_* > 0$ such that for any initial data w_0 that satisfy

$$\|w_0\|_{L^2(\mathbb{R})} \le e^{(\alpha_0 + C_\phi)t_0} \|v_0\|_{L^2(\mathbb{R})},\tag{4.70}$$

it exists a mild solution w on $(0, T'_*]$. Using (4.68), we have that $w_0 := v(t_0 - T'_*/2, \cdot)$ satisfies (4.70). Therefore, using an argument of uniqueness, we deduce that $v(t_0 - T'_*/2 + t, \cdot) = w(t, \cdot)$ for all $t \in [0, T'_*/2)$. To finish with, we define \tilde{v} by $\tilde{v} = v$ on $[0, t_0)$ and $\tilde{v}(t_0 - T'_*/2 + t, \cdot) = w(t, \cdot)$ for $t \in [T'_*/2, T'_*]$. Hence, \tilde{v} is a mild solution on $[0, t_0 + T'_*/2]$ with initial datum v_0 , which gives us a contradiction.

5. Regularity of the Solution

This section is devoted to the proof of the existence of classical solutions v to (1.6).

Proposition 5.1 (Solution in the classical sense). Let $v_0 \in L^2(\mathbb{R})$, $\phi \in C_b^2(\mathbb{R})$ and T > 0. The unique mild solution $v \in C([0,T];L^2(\mathbb{R})) \cap C((0,T];H^2(\mathbb{R}))$ of (1.6) belongs to $C^{1,2}((0,T] \times \mathbb{R})$ and satisfies

$$\partial_t v + \partial_x \left(\frac{v^2}{2} + u_{\phi} v \right) + \mathcal{O}[v] - \partial_{xx}^2 v = 0, \tag{5.1}$$

on $(0,T] \times \mathbb{R}$ in the classical sense.

Proof.

Step 1 (C^2 -regularity in space). Let us take any $t_0 \in (0, T]$ as initial time and let $T' \in (0, T - t_0]$. Differentiating the Duhamel formulation (2.1) two times with respect to the space, we get for any $t \in [0, T']$,

$$\partial_{xx}^2 v(t+t_0,\cdot) = K(t,\cdot) * \partial_{xx}^2 v(t_0,\cdot) - \int_0^t \partial_x K(t-s,\cdot) * (u_1+u_2)(t_0+s,\cdot) ds,$$
(5.2)

where $u_1 := (\partial_x v)^2 + v \partial_{xx}^2 v$ and $u_2 := v \partial_x^2 u_{\phi} + 2 \partial_x u_{\phi} \partial_x v + u_{\phi} \partial_{xx}^2 v$. Since $v \in C((0,T]; H^2(\mathbb{R}))$ then $u_2 \in C((0,T]; L^2(\mathbb{R}))$ and from the Sobolev embedding $H^2(\mathbb{R}) \hookrightarrow C_b^1(\mathbb{R})$, we get that $u_1 \in C((0,T]; L^1(\mathbb{R}) \cap L^2(\mathbb{R}))$. Let us now define the following functions:

$$F_i(t,x) := \int_0^t \partial_x K(t-s,\cdot) * u_i(t_0+s,\cdot)(x) ds, \quad \text{for } i = 1,2.$$
(5.3)

For all $x, y \in \mathbb{R}$, we have thanks to Cauchy-Schwartz inequality

$$\begin{aligned} \left| \partial_{x} K(t-s,\cdot) * u_{i}(t_{0}+s,\cdot)(x) - \partial_{x} K(t-s,\cdot) * u_{i}(t_{0}+s,\cdot)(y) \right| \\ &\leq \int_{\mathbb{R}} \left| \partial_{x} K(t-s,z) \right| \left| u_{i}(t_{0}+s,x-z) - u_{i}(t_{0}+s,y-z) \right| dz \\ &\leq \left\| \mathcal{T}_{(x-y)}(u_{i}(t_{0}+s,\cdot)) - u_{i}(s+t_{0},\cdot) \right\|_{L^{2}(\mathbb{R})} \left\| \partial_{x} K(t-s,\cdot) \right\|_{L^{2}(\mathbb{R})}, \end{aligned}$$
(5.4)

where $\mathcal{T}_{z}\varphi$ denotes the translated function $x \to \varphi(x + z)$.

Therefore, for all $x, y \in \mathbb{R}$ and all $t \in [0, T']$, we deduce that

$$\begin{aligned} \left| F_{i}(t,x) - F_{i}(t,y) \right| &\leq \int_{0}^{t} K_{0}(t-s)^{-3/4} \left\| \mathcal{T}_{(x-y)}(u_{i}(t_{0}+s,\cdot)) - u_{i}(t_{0}+s,\cdot) \right\|_{L^{2}(\mathbb{R})} ds \\ &\leq 4K_{0}T'^{1/4} \sup_{s \in [0,T']} \left\| \mathcal{T}_{(x-y)}(\overline{u}_{i}(s,\cdot)) - \overline{u}_{i}(s,\cdot) \right\|_{L^{2}(\mathbb{R})'} \end{aligned}$$
(5.5)

where $\overline{u}_i(s, \cdot) = u_i(t_0 + s, \cdot)$. Then, \overline{u}_i is uniformly continuous with values in L^2 as a continuous function on a compact set [0, T']. Therefore, for any $\epsilon > 0$, there exists a finite sequence $0 = s_0 < s_1 < \cdots < s_N = T'$ such that for any $s \in [0, T']$, there exists $j \in \{0, \dots, N-1\}$ such that

$$\left\|\overline{u}_{i}(s,\cdot) - \overline{u}_{i}(s_{j},\cdot)\right\|_{L^{2}(\mathbb{R})} \leq \epsilon.$$
(5.6)

Therefore, using (5.5) we have

$$\begin{aligned} |F_{i}(t,x) - F_{i}(t,y)| &\leq 4K_{0}T'^{1/4} \sup_{s \in [0,T']} \|\mathcal{T}_{(x-y)}(\overline{u}_{i}(s,\cdot)) - \mathcal{T}_{(x-y)}(\overline{u}_{i}(s_{j},\cdot))\|_{L^{2}} \\ &+ 4K_{0}T'^{1/4} \left\{ \|\mathcal{T}_{(x-y)}(\overline{u}_{i}(s_{j},\cdot)) - \overline{u}_{i}(s_{j},\cdot)\|_{L^{2}} + \sup_{s \in [0,T]} \|\overline{u}_{i}(s,\cdot) - \overline{u}_{i}(s_{j},\cdot)\|_{L^{2}} \right\}. \end{aligned}$$

$$(5.7)$$

And since $\|\mathcal{T}_{(x-y)}(\overline{u}_i(s,\cdot)) - \mathcal{T}_{(x-y)}(\overline{u}_i(s_j,\cdot))\|_{L^2(\mathbb{R})} = \|\overline{u}_i(s,\cdot) - \overline{u}_i(s_j,\cdot)\|_{L^2(\mathbb{R})}$, we get

$$\left|F_{i}(t,x)-F_{i}(t,y)\right| \leq 4K_{0}T'^{1/4}\left\{\left\|\boldsymbol{\mathcal{T}}_{(x-y)}(\overline{u}_{i}(s_{j},\cdot))-\overline{u}_{i}(s_{j},\cdot)\right\|_{L^{2}(\mathbb{R})}+2\varepsilon\right\}.$$
(5.8)

And since the translated function is continuous in $L^2(\mathbb{R})$, we have

$$\left\|\mathcal{T}_{(x-y)}(\overline{u}_i(s_j,\cdot)) - \overline{u}_i(s_j,\cdot)\right\|_{L^2(\mathbb{R})} \longrightarrow 0,$$
(5.9)

as $(x - y) \rightarrow 0$. Hence,

$$\lim_{(x-y)\to 0} \sup \left| F_i(t,x) - F_i(t,y) \right| \le 2\epsilon.$$
(5.10)

Taking the infimum with respect to $\epsilon > 0$, we infer that F_i is continuous with respect to the variable x. Moreover, arguing as the proof of Proposition 2.9, we get that $F_i \in C([0,T']; L^2(\mathbb{R}))$. From classical results, we then deduce that F_i is continuous with respect to the couple (t, x) on $[0, T'] \times \mathbb{R}$.

Moreover, since $v(t_0, \cdot) \in H^2(\mathbb{R})$, we can easily check that $(t, x) \to K(t, \cdot) * \partial_{xx}^2 v(t_0, \cdot)(x)$ is continuous on $(0, T] \times \mathbb{R}$. Finally, we get that $\partial_{xx}^2 v \in C([t_0, T] \times \mathbb{R})$ and since t_0 is arbitrary in (0, T], we conclude that $\partial_{xx}^2 v \in C((0, T] \times \mathbb{R})$.

Step 2 (C^1 -regularity in time). From Proposition 4.3, we know that the terms $\partial_t v$ and $-\partial_x(v^2/2 + u_\phi v) + \partial_{xx}^2 v - \mathcal{O}[v]$ have the same regularity. Moreover, by the Step 1 of this proposition, we have that $\partial_{xx}^2 v \in C((0,T] \times \mathbb{R})$, and from Sobolev embeddings and Remark 2.8, we deduce that $\partial_x(v^2/2 + u_\phi v)$ and $\mathcal{O}[v]$ belong to $C((0,T] \times \mathbb{R})$. Finally, we obtain that $\partial_t v \in C((0,T] \times \mathbb{R})$ and thus $v \in C^{1,2}((0,T] \times \mathbb{R})$. The proof of this proposition is now complete.

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References

- [1] A. C. Fowler, "Mathematics and the environment," lecture notes, http://www.maths.ox.ac.uk/ ~fowler/courses/mathenvo.html.
- [2] P.-Y. Lagrée, Asymptotic Methods in Fluid Mechanics: Survey and Recent Advances, vol. 523, Springer, New York, NY, USA, 2010, CISM International Centre for Mechanical Sciences Udine.
- [3] A. C. Fowler, "Dunes and Drumlins," in Geomorphological Fluid Mechanics, A. Provenzale and N. Balmforth, Eds., vol. 211, pp. 430–454, Springer, Berlin, Germany, 2001. [4] N. Alibaud, P. Azerad, and D. Isèbe, "A non-monotone nonlocal conservation law for dune mor-
- phodynamics," Differential and Integral Equations, vol. 23, no. 1-2, pp. 155-188, 2010.
- [5] P. Azerad and A. Bouharguane, "Finite difference approximationsfor a fractional diffusion/antidiffusion equation," http://arxiv.org/abs/1104.4861.
- [6] I. Podlubny, An Introduction to Fractional Derivatives, Fractional Differential Equations, to Methods Of Their Solution and Some of Their Applications, vol. 198 of Mathematics in Science and Engineering, Academic Press, San Diego, CA, 1999.
- [7] P. Azerad, A. Bouharguane, and J.-F. Crouzet, "Simultaneous denoising and enhancement of signals by a fractal conservation law," Communications in Nonlinear Science and Numerical Simulation, vol. 17, no. 2, pp. 867-881.
- [8] B. Alvarez-Samaniego and P. Azerad, "Existence of travelling-wave solutions and local well-posedness of the Fowler equation," Discrete and Continuous Dynamical Systems. Series B, vol. 12, no. 4, pp. 671-692, 2009.
- [9] X. Chen, "Existence, uniqueness, and asymptotic stability of traveling waves in nonlocal evolution equations," Advances in Differential Equations, vol. 2, no. 1, pp. 125-160, 1997.
- [10] J. Droniou, T. Gallouet, and J. Vovelle, "Global solution and smoothing effect for a non-local regularization of a hyperbolic equation," Journal of Evolution Equations, vol. 3, no. 3, pp. 499–521, 2003.
- [11] A. Bouharguane, "On the instability of a nonlocal conservation law," Discrete and Continuous *Dynamical Systems. Series S*, vol. 5, no. 3, 2012.



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