

SOME SPECIAL SUBSPACES OF A FINSLER SPACE

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Abstract. In the present paper are studied such subspaces of a Finsler space for which the absolute differential of the tangent or normal vectors have special positions.

1. Introduction. The equation of a subspace F_m of a Finsler space F_n , the definitions of the tangent vectors B_α^i , the normal vectors N_μ^i , and the induced and intrinsic connection coefficients and curvature tensors are the same as in [6], [2] and [3]; so they are omitted. The induced connection coefficients and curvature tensors shall be denoted as usual by —.

Let us denote by $T_H(P)$ the subspace of the tangent space of F_n at $P(x, \dot{x}) = (X^i(u^\alpha), \dot{u}^\alpha)$ spanned by B_α^i and by $T_V(P)$ the subspace spanned by N_μ^i .

The object of the present paper is to study special subspaces which satisfy some of the following conditions at a fixed P for every displacement $(du^\alpha, d\dot{u}^\alpha)$ on the subspace F_m :

- 1) $DB_\alpha^i \in T_H \Leftrightarrow DN_\mu^i \in T_V \Leftrightarrow (DB_\alpha^i \in T_H) \wedge (DN_\mu^i \in T_V)$
- 1a) $DN_\mu^i = 0 \Rightarrow DB_\alpha^i \in T_H$
- 1b) $DB_\alpha^i \Rightarrow DN_\mu^i \in T_V$
- 2) $DB_\alpha^i \in T_V$
- 3) $DN_\mu^i \in T_H$
- 2a) = 3a) $(DB_\alpha^i \in T_V) \wedge (DN_\mu^i \in T_H)$

for every $\alpha = 1, 2, \dots, m$, $\mu = m + 1, \dots, n$.

Cases 1a) and 1b) are special cases of 1); 2a) = 3a) is a special case of 2) or 3).

For the case 1) the induced and intrinsic connection coefficients are the same and the normal curvature $\overset{\nu}{N}(u, \dot{u}) = 0$ for every curve $u^\alpha = u^\alpha(s)$ trough P . Theorem 1.1 gives equivalent conditions for F_m to satisfy the conditions of case 1) for a fixed u and every \dot{u} .

For case 2) the subspace F_m is Riemannian with

$${}^0\overline{R}_{\alpha\beta\gamma}{}^\delta = 0, \quad {}^0\overline{P}_{\alpha\beta\gamma}{}^\delta = 0, \quad {}^0\overline{S}_{\alpha\beta\gamma}{}^\delta = 0.$$

For case 3) we have

$${}^1\overline{R}_{\mu\beta\gamma}{}^\nu = 0, \quad {}^1\overline{P}_{\mu\beta\gamma}{}^\nu = 0, \quad {}^1\overline{S}_{\mu\beta\gamma}{}^\nu = 0.$$

2. Case 1). $DB_\alpha^i \in T_H$. For any subspace F_m of F_n we have

$$\begin{aligned} DB_\alpha^i &= (\overline{\Gamma}_{\alpha\beta}^{*\delta} du^\beta + \overline{A}_{\alpha\beta}{}^\delta \overline{D}l^\beta) B_\delta^i + (\overline{\theta}_{\alpha\beta}^{*\mu} du^\beta + \overline{A}_{\alpha\beta}{}^\mu) N_\mu^i, \\ DN_\mu^i &= (-\overline{\theta}_{\mu\beta}^{*\delta} du^\beta - \overline{A}_{\mu\beta}{}^\delta \overline{D}l^\beta) B_\delta^i + (\overline{\lambda}_{\mu\beta}^{*\nu} du^\beta + \overline{A}_{\mu\beta}{}^\nu \overline{D}l^\beta) N_\nu^i. \end{aligned}$$

In the case 1) these formulae become

$$(2.1) \quad DB_\alpha^i = (\overline{\Gamma}_{\alpha\beta}^{*\delta} du^\beta + \overline{A}_{\alpha\beta}{}^\delta \overline{D}l^\beta) B_\delta^i,$$

$$(2.2) \quad DN_\mu^i = (\overline{\lambda}_{\mu\beta}^{*\nu} du^\beta + \overline{A}_{\mu\beta}{}^\nu \overline{D}l^\beta) N_\nu^i.$$

In this case $\overline{\theta}_{\alpha\beta}^{*\mu} du^\beta + \overline{A}_{\alpha\beta}{}^\mu \overline{D}l^\beta = 0$, for every du^β and $\overline{D}l^\beta$, so

$$(2.3) \quad \overline{\theta}_{\alpha\beta}^{*\mu} du^\beta = 0, \quad \overline{A}_{\alpha\beta}{}^\mu \overline{D}l^\beta = 0$$

for all

$$\alpha, \beta = 1, 2, \dots, m \quad \mu = m + 1, \dots, n.$$

From (2.1), (2.3) and

$$\overline{\theta}_{\mu\alpha\beta}^* = -\overline{\theta}_{\mu\alpha\beta}^*, \quad \overline{A}_{\mu\alpha\beta} = -\overline{A}_{\mu\alpha\beta}$$

we obtain

$$(2.4) \quad \overline{\theta}_{\mu\alpha\beta}^* = 0 \quad \overline{A}_{\mu\alpha\beta} = 0$$

for all

$$\alpha, \beta = 1, 2, \dots, m \quad \mu = m + 1, \dots, n.$$

As for any subspace F_m we have

$$Dl^k = B_\alpha^k Dl^\alpha + \overline{H}_\beta^k du^\beta$$

and for case 1)

$$Dl^k = D(B_\alpha^k l^\alpha) = B_\alpha^k \overline{D}l^\alpha$$

we conclude that in this case

$$\overline{H}_\beta^k du^\beta = 0.$$

The above equation is true for any du^β so that in case 1)

$$(2.5) \quad \overline{H}_\beta^k = 0, \quad k = 1, \dots, n \quad \beta = 1, \dots, m.$$

From (2.5) it follows that the corresponding equations for $\overline{\Gamma}_{\alpha\beta}^{\delta}$ and $\overline{\lambda}_{\mu\beta}^{*\nu}$ reduce to

$$(2.6) \quad \overline{\Gamma}_{\alpha\gamma\beta}^* = g_{ir} B_\gamma^r (B_{\alpha\beta}^i + \Gamma_{j\ k}^{*i} B_{\alpha\beta}^{jk})$$

$$(2.7) \quad \overline{\lambda}_{\mu\beta}^{*\nu} = g_{ir} N_\nu^r (\partial_\beta N_\nu^i - \delta 0 t \partial_\delta N_\nu^i \overline{\Gamma}_{\alpha}^{*\delta} + \overline{\Gamma}_{j\ k}^{*i} N_\nu^j B_\beta^k).$$

Tensors $\overline{A}_{\alpha\beta\gamma}$ and $\overline{A}_{\mu\nu\gamma}$ are determined by

$$(2.8) \quad \overline{A}_{\alpha\beta\gamma} = A_{ijk} B_{\alpha\beta\gamma}^{ijk} = L(u, \dot{u}) 2^{-1} \dot{\partial}_\gamma g_{\alpha\beta}(u, \dot{u})$$

$$(2.9) \quad \overline{A}_{\mu\nu\gamma} = g_{ij} N_\nu^j L \dot{\partial}_\gamma N_\mu^i + A_{ijk} N_\mu^i N_\nu^j B_\gamma^k.$$

The normal curvature $\overset{\nu}{N}$ of a curve $u^\alpha = u^\alpha(s)$ of the subspace F_m in the direction of $\overset{\nu}{N}_i$ is given by

$$\overset{\nu}{N}(u, \dot{u}) = L^{-2}(u, \dot{u}) \overline{\theta}_{\alpha\beta}^{*\nu} \dot{u}^\alpha \dot{u}^\beta \quad (\dot{u}^\alpha = du^\alpha / du^s)$$

From (2.3) it follows that

$$(2.10) \quad \overset{\nu}{N}(u, \dot{u}) = 0$$

for every curve $u^\alpha = u^\alpha(s)$ through the point (u) .

From (2.1) and (2.2) we obtain

$$(2.11) \quad [\Delta D] B_\alpha^i = \overline{\Omega}_\alpha^\delta(d, \delta) B_\delta^i = \{2^{-1} {}^0 \overline{R}_{\alpha\beta\gamma}^\delta [du^\beta \delta u^\gamma] + {}^0 \overline{P}_{\alpha\beta\gamma}^\delta [du^\beta \overline{\Delta} l^\gamma] + 2^{-1} {}^0 \overline{S}_{\alpha\beta\gamma}^\delta [\overline{D} l^\beta \overline{\Delta} l^\gamma]\} B_\delta^i$$

$$(2.12) \quad [\Delta D] N_\mu^i = \overline{\Omega}_\mu^\nu(d, \delta) N_\nu^i = \{2^{-1} {}^1 \overline{R}_{\mu\beta\gamma}^\nu [du^\beta \delta u^\gamma] + {}^1 \overline{P}_{\mu\beta\gamma}^\nu [du^\beta \overline{\Delta} l^\gamma] + 2^{-1} {}^1 \overline{S}_{\mu\beta\gamma}^\nu [\overline{D} l^\beta \overline{\Delta} l^\gamma]\} N_\nu^i$$

It may be seen that in case 1)

$${}^0 \overline{R}_{\alpha\beta\gamma}^\mu, \quad {}^0 \overline{P}_{\alpha\beta\gamma}^\mu, \quad {}^0 \overline{S}_{\alpha\beta\gamma}^\mu$$

The definitions of curvature tensors given above and in the sequel are given in [6].

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Some vector field $\xi^i(x(u) B_\alpha \dot{u}^\alpha)$ defined on the subspace F_m , may be decomposed in the following way

$$\xi^i = B_\alpha^i \xi^\alpha + \overset{\nu}{N}_i \xi^\mu$$

Using the known formulae

$$[\Delta D]\xi^i = \{2^{-1}R_{j\ hk}^i[dx^h\delta x^k] + P_{j\ hk}^i[dx^h\Delta l^k] + 2^{-1}S_{j\ hk}^i[Dl^h\Delta l^k]\}\xi^j$$

$$dx^h = B_\alpha^h du^\alpha$$

and for case 1)

$$Dl^h = B_\alpha^h \bar{D}l^\alpha$$

we get

$$(2.14) \quad R_{j\ hk}^i \xi^j B_{\beta\ \gamma}^{h\ k} = {}^0\bar{R}_{\alpha\ \beta\ \gamma}^\varepsilon \xi^\alpha B_\varepsilon^i + {}^1\bar{R}_{\mu\ \beta\ \gamma}^\nu \xi^\mu N_\nu^i$$

The above formula is true for tensors P and S . Comparing the coefficients of ξ^α and ξ^μ we obtain

$$(2.15) \quad \begin{aligned} \text{a)} \quad & R_{j\ hk}^i B_{\alpha\ \beta\ \gamma}^{j\ h\ k} = {}^0\bar{R}_{\alpha\ \beta\ \gamma}^\varepsilon B_\varepsilon^i \\ \text{b)} \quad & R_{j\ hk}^i N_\mu^j B_{\beta\ \gamma}^{h\ k} = {}^1\bar{R}_{\mu\ \beta\ \gamma}^\nu N_\nu^i \\ \text{c)} \quad & P_{j\ hk}^i B_{\alpha\ \beta\ \gamma}^{j\ h\ k} = {}^0\bar{P}_{\alpha\ \beta\ \gamma}^\varepsilon B_\varepsilon^i \\ \text{d)} \quad & P_{j\ hk}^i N_\mu^j B_{\beta\ \gamma}^{h\ k} = {}^1\bar{P}_{\mu\ \beta\ \gamma}^\nu N_\nu^i \\ \text{e)} \quad & S_{j\ hk}^i B_{\alpha\ \beta\ \gamma}^{j\ h\ k} = {}^0\bar{S}_{\alpha\ \beta\ \gamma}^\varepsilon B_\varepsilon^i \\ \text{f)} \quad & S_{j\ hk}^i N_\mu^j B_{\beta\ \gamma}^{h\ k} = {}^1\bar{S}_{\mu\ \beta\ \gamma}^\nu N_\nu^i \end{aligned}$$

If we define the induced covariant differentiations $\overset{1}{\nabla}$ and $\overset{1}{\nabla}$ for some mixed tensor $T_{\alpha\nu}^{\beta\mu}$ in the form

$$\begin{aligned} T_{\alpha\nu}^{\beta\mu} \overset{1}{\nabla}_\gamma &= \partial_\gamma T_{\alpha\nu}^{\beta\mu} - \dot{\partial}_\varkappa T_{\alpha\nu}^{\beta\mu} \bar{\Gamma}_{\gamma}^{*\varkappa} - T_{\varkappa\nu}^{\beta\mu} \bar{\Gamma}_{\alpha\ \gamma}^{*\varkappa} + \\ & T_{\varkappa\nu}^{\beta\mu} \bar{\Gamma}_{\varkappa\ \gamma}^{\beta} - T_{\alpha\ \zeta}^{\beta\mu} \bar{\lambda}_{\nu\ \gamma}^{*\zeta} + T_{\alpha\nu}^{\beta\mu} \bar{\lambda}_{\xi\ \gamma}^{*\mu} \\ T_{\alpha\nu}^{\beta\mu} \overset{1}{\nabla}_\gamma &= L\partial_\gamma T_{\alpha\nu}^{\beta\mu} - T_{\varkappa\nu}^{\beta\mu} \bar{A}_{\alpha\ \gamma}^{\varkappa} + T_{\varkappa\nu}^{\beta\mu} \bar{A}_{\varkappa\ \gamma}^{\beta} - \\ & T_{\alpha\ \xi}^{\beta\mu} \bar{A}_{\nu\ \gamma}^{\xi} + T_{\alpha\nu}^{\beta\zeta} \bar{A}_{\zeta\ \gamma}^{\mu}. \end{aligned}$$

then the Bianchi identities ([6], (3.1)—(3.3)) for the case 1) reduce to

$$(2.16) \quad \begin{aligned} \text{a)} \quad & {}^0\bar{R}_{\alpha\ \beta\ \gamma}^\varepsilon \overset{1}{\nabla}_\delta + {}^0\bar{P}_{\alpha\ [\gamma|\delta|}^\varepsilon \overset{1}{\nabla}\beta] + {}^0\bar{R}_{\alpha\ \varkappa[\gamma}^\varepsilon A_{\beta]}^{\varkappa}{}_\delta + \\ & {}^0\bar{S}_{\alpha\ \delta\varkappa}^\varepsilon {}^0\bar{K}_{o\ \beta\ \gamma}^{\varkappa} - {}^0\bar{P}_{\alpha\ [\gamma|\varkappa}^\varepsilon \dot{\partial}_\delta \bar{\Gamma}_{\nu\ |\beta]}^{*\varkappa} u^\nu = -l_\delta \bar{A}_{\alpha\varkappa}^\varepsilon {}^0\bar{K}_{o\ \beta\ \gamma}^{\varkappa}, \\ \text{b)} \quad & {}^1\bar{R}_{\mu\ \beta\ \gamma}^\nu \overset{1}{\nabla}_\delta + {}^1\bar{P}_{\mu\ [\gamma|\delta|}^\nu \overset{1}{\nabla}\beta] + {}^1\bar{R}_{\mu\ \varkappa[\gamma}^\nu A_{\beta]}^{\varkappa}{}_\delta + \\ & {}^1\bar{S}_{\mu\ \delta\varkappa}^\nu {}^0\bar{K}_{o\ \beta\ \gamma}^{\varkappa} - {}^1\bar{P}_{\mu\ [\gamma|\varkappa}^\nu \dot{\partial}_\delta \bar{\Gamma}_{\nu\ |\beta]}^{*\varkappa} u^\nu = -l_\delta \bar{A}_{\mu\varkappa}^\nu {}^0\bar{K}_{o\ \beta\ \gamma}^{\varkappa}, \\ \text{c)} \quad & ({}^0\bar{R}_{\alpha\ \beta\ \gamma}^\varepsilon \overset{1}{\nabla}_\delta + {}^0\bar{P}_{\alpha\ \beta\varkappa}^\varepsilon {}^0\bar{K}_{o\ \gamma\ \delta}^{\varkappa}) + \text{cycl}(\beta\gamma\delta) = 0, \\ \text{d)} \quad & ({}^1\bar{R}_{\mu\ \beta\ \gamma}^\nu \overset{1}{\nabla}_\delta + {}^1\bar{P}_{\mu\ \beta\varkappa}^\nu {}^0\bar{K}_{o\ \gamma\ \delta}^{\varkappa}) + \text{cycl}(\beta\gamma\delta) = 0, \end{aligned}$$

$$\begin{aligned}
\text{e)} \quad & {}^0\overline{P}_{\alpha\beta[\gamma|\delta]}{}^1\overline{\Gamma}_{\beta]} + A_{\alpha}{}^{\varkappa}{}_{[\delta}{}^0\overline{P}|\alpha{}^{\varepsilon}{}_{\varkappa|\gamma]} + {}^0\overline{S}_{\alpha\gamma\beta}{}^1\overline{\Gamma}_{\beta]} + {}^0\overline{S}_{|\alpha{}^{\varepsilon}{}_{\iota[\gamma]}|\delta]} \dot{\delta} \overline{\Gamma}_{\varkappa\beta}^{*\iota} \dot{u}^{\varkappa} = \\
& L l_{[\delta} \dot{\delta}_{\gamma]} \overline{\Gamma}_{\alpha\beta}^{*\varepsilon}, \\
\text{f)} \quad & {}^1\overline{P}_{\mu\nu\beta[\gamma|\delta]}{}^1\overline{\Gamma}_{\beta]} + A_{\beta}{}^{\varkappa}{}_{[\delta}{}^1\overline{P}|\mu{}^{\nu}{}_{\varkappa|\gamma]} + {}^1\overline{S}_{\mu\nu\gamma\beta}{}^1\overline{\Gamma}_{\beta]} + {}^1\overline{S}_{|\mu{}^{\nu}{}_{\iota[\gamma]}|\delta]} \dot{\delta} \overline{\Gamma}_{\varkappa\beta}^{*\iota} \dot{u}^{\varkappa} = \\
& L l_{[\delta} \dot{\delta}_{\gamma]} \overline{\lambda}_{\mu\beta}^{*\nu}.
\end{aligned}$$

If we denote by D_i the absolute differential in F_n which corresponds to the displacement $(d_i u^\alpha, d_i \dot{u}^\alpha)$ ($i = 1, 2$) in F_m then from (2.11), (2.12), (2.15a), (2.15b) we have

$$\begin{aligned}
(2.17) \quad & ([D_2 D_1] R_j{}^i{}_{hk}) B_{\alpha\beta\gamma}{}^j{}^h{}^k = {}^0\overline{R}_{\alpha\beta\gamma}{}^{\varepsilon} \overline{\Omega}_{\varepsilon}{}^{\delta}(d_1, d_2) B_{\delta}{}^i - \\
& {}^0\overline{R}_{\varepsilon\beta\gamma}{}^{\delta} \overline{\Omega}_{\alpha}{}^{\varepsilon}(d_1, d_2) B_{\delta}{}^i - {}^0\overline{R}_{\alpha\varepsilon\gamma}{}^{\delta} \overline{\Omega}_{\beta}{}^{\varepsilon}(d_1, d_2) B_{\delta}{}^i - \\
& {}^0\overline{R}_{\alpha\beta\varepsilon}{}^{\delta} \overline{\Omega}_{\gamma}{}^{\varepsilon}(d_1, d_2) B_{\delta}{}^i,
\end{aligned}$$

$$\begin{aligned}
(2.18) \quad & ([D_2 D_1] R_j{}^i{}_{hk}) N_{\mu}{}^j{}^h{}^k B_{\beta\gamma}{}^i = {}^1\overline{R}_{\mu\beta\gamma}{}^{\nu} \overline{\Omega}_{\nu}{}^{\psi}(d_1, d_2) N_{\psi}{}^i - \\
& {}^1\overline{R}_{\psi\beta\gamma}{}^{\nu} \overline{\Omega}_{\mu}{}^{\psi}(d_1, d_2) N_{\nu}{}^j - {}^1\overline{R}_{\mu\varepsilon\gamma}{}^{\nu} \overline{\Omega}_{\beta}{}^{\varepsilon}(d_1, d_2) N_{\nu}{}^j - \\
& {}^1\overline{R}_{\mu\beta\varepsilon}{}^{\nu} \overline{\Omega}_{\gamma}{}^{\varepsilon}(d_1, d_2) N_{\varepsilon}{}^j.
\end{aligned}$$

Formulae of type (2.17), (2.18) are satisfied for tensors P and S and we may get them substituting the letter R with P and S .

If the space F_n satisfies the relation

$$(2.19) \quad [D_2 D_1] R_j{}^i{}_{hk} = 0$$

then from (2.17) and (2.18) we have

$$\begin{aligned}
(2.20) \quad & {}^0\overline{R}_{\alpha\beta\gamma}{}^{\varepsilon} \overline{\Omega}_{\varepsilon}{}^{\delta}(d_1, d_2) - {}^0\overline{R}_{\varepsilon\beta\gamma}{}^{\delta} \overline{\Omega}_{\alpha}{}^{\varepsilon}(d_1, d_2) - \\
& {}^0\overline{R}_{\alpha\varepsilon\gamma}{}^{\delta} \overline{\Omega}_{\beta}{}^{\varepsilon}(d_1, d_2) - {}^0\overline{R}_{\alpha\beta\varepsilon}{}^{\delta} \overline{\Omega}_{\gamma}{}^{\varepsilon}(d_1, d_2) = 0,
\end{aligned}$$

$$\begin{aligned}
(2.21) \quad & {}^1\overline{R}_{\mu\beta\gamma}{}^{\nu} \overline{\Omega}_{\nu}{}^{\psi}(d_1, d_2) - {}^1\overline{R}_{\psi\beta\gamma}{}^{\nu} \overline{\Omega}_{\mu}{}^{\psi}(d_1, d_2) - \\
& {}^1\overline{R}_{\mu\varepsilon\gamma}{}^{\nu} \overline{\Omega}_{\beta}{}^{\varepsilon}(d_1, d_2) - {}^1\overline{R}_{\mu\beta\varepsilon}{}^{\nu} \overline{\Omega}_{\gamma}{}^{\varepsilon}(d_1, d_2) = 0.
\end{aligned}$$

If the space F_n satisfies

$$(2.19) \quad \text{a) } [D_2 D_1] P_j{}^i{}_{hk} = 0 \quad \text{or} \quad \text{b) } [D_2 D_1] S_j{}^i{}_{hk} = 0$$

then the induced curvature tensors of the subspace ${}^0\overline{P}_{\alpha\beta\gamma}{}^{\delta}$, ${}^1\overline{P}_{\mu\beta\gamma}{}^{\nu}$, ${}^0\overline{S}_{\alpha\beta\gamma}{}^{\delta}$, ${}^1\overline{S}_{\mu\beta\gamma}{}^{\nu}$ satisfy the equations of type (2.20) and (2.21) and we get these equations when the letter R is substituted by P or S .

If (2.19) is true for every D_1, D_2 , i.e., the tensor R is parallel on the subspace F_m , then from (2.20) and (2.21) we obtain

$$(2.23) \quad \begin{aligned} \text{a)} \quad & {}^0\overline{R}_{\alpha\beta\gamma}{}^\varepsilon {}^0\overline{R}_{\varepsilon\iota\kappa}{}^\delta - {}^0\overline{R}_{\varepsilon\beta\gamma}{}^\delta {}^0\overline{R}_{\alpha\iota\kappa}{}^\varepsilon - {}^0\overline{R}_{\alpha\varepsilon\gamma}{}^\delta {}^0\overline{R}_{\beta\iota\kappa}{}^\varepsilon - {}^0\overline{R}_{\alpha\beta\varepsilon}{}^\delta {}^0\overline{R}_{\gamma\iota\kappa}{}^\varepsilon = 0 \\ \text{b)} \quad & {}^0\overline{R}_{\alpha\beta\gamma}{}^\varepsilon {}^0\overline{P}_{\varepsilon\iota\kappa}{}^\delta - {}^0\overline{R}_{\varepsilon\beta\gamma}{}^\delta {}^0\overline{P}_{\alpha\iota\kappa}{}^\varepsilon - {}^0\overline{R}_{\alpha\varepsilon\gamma}{}^\delta {}^0\overline{P}_{\beta\iota\kappa}{}^\varepsilon - {}^0\overline{R}_{\alpha\beta\varepsilon}{}^\delta {}^0\overline{P}_{\gamma\iota\kappa}{}^\varepsilon = 0 \\ \text{c)} \quad & {}^0\overline{R}_{\alpha\beta\gamma}{}^\varepsilon {}^0\overline{S}_{\varepsilon\iota\kappa}{}^\delta - {}^0\overline{R}_{\varepsilon\beta\gamma}{}^\delta {}^0\overline{S}_{\alpha\iota\kappa}{}^\varepsilon - {}^0\overline{R}_{\alpha\varepsilon\gamma}{}^\delta {}^0\overline{S}_{\beta\iota\kappa}{}^\varepsilon - {}^0\overline{R}_{\alpha\beta\varepsilon}{}^\delta {}^0\overline{S}_{\gamma\iota\kappa}{}^\varepsilon = 0 \\ \text{d)} \quad & {}^1\overline{R}_{\mu\beta\gamma}{}^\psi {}^1\overline{R}_{\psi\iota\kappa}{}^\nu - {}^1\overline{R}_{\psi\beta\gamma}{}^\nu {}^1\overline{R}_{\mu\iota\kappa}{}^\psi - {}^1\overline{R}_{\mu\varepsilon\gamma}{}^\nu {}^1\overline{R}_{\beta\iota\kappa}{}^\varepsilon - {}^1\overline{R}_{\mu\beta\varepsilon}{}^\nu {}^1\overline{R}_{\gamma\iota\kappa}{}^\varepsilon = 0 \\ \text{e)} \quad & {}^1\overline{R}_{\mu\beta\gamma}{}^\psi {}^1\overline{P}_{\psi\iota\kappa}{}^\nu - {}^1\overline{R}_{\psi\beta\gamma}{}^\nu {}^1\overline{P}_{\mu\iota\kappa}{}^\psi - {}^1\overline{R}_{\mu\varepsilon\gamma}{}^\nu {}^1\overline{P}_{\beta\iota\kappa}{}^\varepsilon - {}^1\overline{R}_{\mu\beta\varepsilon}{}^\nu {}^1\overline{P}_{\gamma\iota\kappa}{}^\varepsilon = 0 \\ \text{f)} \quad & {}^1\overline{R}_{\mu\beta\gamma}{}^\psi {}^1\overline{S}_{\psi\iota\kappa}{}^\nu - {}^1\overline{R}_{\psi\beta\gamma}{}^\nu {}^1\overline{S}_{\mu\iota\kappa}{}^\psi - {}^1\overline{R}_{\mu\varepsilon\gamma}{}^\nu {}^1\overline{S}_{\beta\iota\kappa}{}^\varepsilon - {}^1\overline{R}_{\mu\beta\varepsilon}{}^\nu {}^1\overline{S}_{\gamma\iota\kappa}{}^\varepsilon = 0 \end{aligned}$$

If (2.23) is true for every D_1, D_2 , then we easily obtain equations similar to (2.22) for the tensors P and S .

We shall examine what form the intrinsic connection coefficients take for case 1. In the subspace F_m with respect to the intrinsic connection coefficients DB_α^i and DN_μ^i take the form

$$\begin{aligned} DB_\alpha^i &= [(\Gamma_{\alpha\beta}^{*\delta} + \Lambda_{\alpha\beta}^\delta)du^\beta + A_{\alpha\beta}^\delta du^\beta Dl^\beta]B_\delta^i + (\theta_{\alpha\beta}^{*\mu} du^\beta + A_{\alpha\beta}^\mu Dl^\beta)N_\mu^i \\ DN_\mu^i &= -(\theta_{\mu\beta}^{*\delta} du^\beta + A_{\alpha\beta}^\delta Dl^\beta)D_\delta^i + (\lambda_{\mu\beta}^{*\nu} du^\beta + A_{\mu\beta}^\nu Dl^\beta)N_\nu^i. \end{aligned}$$

As

$$\begin{aligned} \theta_{\alpha\beta}^{*\mu} &= \overline{\theta}_{\alpha\beta}^{*\mu} - A_{\alpha\ \varkappa}^\mu A_{\nu\beta}^{\varkappa\nu} N^\nu, \\ A_{\alpha\beta\gamma} &= \overline{A}_{\alpha\beta\gamma}, \quad A_{\alpha\mu\beta} = \overline{A}_{\alpha\mu\beta}, \\ \overline{D}l^\beta &= Dl^\beta = -A_{\nu\gamma}^\beta N^\nu du^\gamma \end{aligned}$$

we have in case 1)

$$(2.24) \quad \theta_{\alpha\beta}^{*\mu} = \overline{\theta}_{\alpha\beta}^{*\mu} = 0$$

$$(2.25) \quad \begin{aligned} A_{\alpha\mu\beta} &= \overline{A}_{\alpha\mu\beta} = 0 \\ \overline{D}l^\beta &= Dl^\beta \end{aligned}$$

From the last equation and

$$Dl^k = B_\alpha^k Dl^\alpha = H_\beta^k du^\beta$$

it follows that

$$H_\beta^k = 0$$

From

$$\Lambda_{\alpha\beta}^\delta = -A_{hkj} B_\beta^j g^{\rho\delta} (H_\alpha^h B_\rho^k - B_\alpha^h H_\rho^k)$$

and $H_\beta^k = 0$ we get immediately

$$(2.26) \quad \Lambda_\alpha^\delta{}_\beta = 0$$

As $\bar{\Gamma}_{\alpha\rho\beta}^*$ and $\Gamma_{\alpha\rho\beta}^*$ are connected by

$$\Gamma_{\alpha\rho\beta}^* = \bar{\Gamma}_{\alpha\rho\beta}^* + A_{ikj} B_\beta^j (H_\alpha^i B_\rho^k - B_\alpha^i H_\rho^k) - A_{\alpha\rho\delta} A^\delta{}_{\nu\beta} \overset{\nu}{N}$$

using $H_\beta^k = 0$, $\overset{\nu}{N} = 0$ we have

$$(2.27) \quad \Gamma_{\alpha\rho\beta} = \bar{\Gamma}_{\alpha\rho\beta}^*.$$

As

$$(2.28) \quad A_\mu{}^\nu{}_\beta = \bar{A}_\mu{}^\nu{}_\beta$$

for any subspace from (2.24) — (2.28) we have:

THEOREM 2.1. *If the subspace F_m of the Finsler space F_n has the property $DB_\alpha^i \in T_H$ for the mixed lineelement $P(u, \dot{u})$ and every $(du^\alpha, d\dot{u}^\alpha)$, then the induced and intrinsic connection coefficients are the same, from which it follows that the induced and intrinsic curvature tensors are the same, and satisfy the same equations at P .*

In all previous equations every quantity and tensor was considered at the fixed lineelement $P(u, \dot{u})$. Let us denote by HF_m the subspace of case 1) for all lineelements (u, \dot{u}) where u is a fixed point and \dot{u} is any direction in the subspace. Then we have the following:

THEOREM 2.2. *The subspace F_m of the Finsler space F_n is HF_m iff one of the following equivalent equations (2.1) (2.5) or (2.10) is satisfied for all directions \dot{u} at fixed point u .*

Proof. From the definition it is obvious that the subspace F_m is HF_m iff (2.1) for fixed u and \dot{u} . Furthermore

$$(2.1) \Rightarrow (2.4) \Rightarrow (2.5)$$

To prove (2.5) \Rightarrow (2.1) from $l^k = B_\alpha^k l^\alpha$, $g_{ij}(x, \dot{x}) N_\mu^i l^j = 0$ we have

$$g_{ij} D N_\mu^i l^j + g_{ij} N_\mu^i D l^j = 0$$

From (2.5) and the equation above we obtain $g_{ij} D N_\alpha^i B_\alpha^j l^\alpha = 0$ for all l^α ; so

$$D N_\mu^i D l^j = (\lambda_\mu{}^\nu{}_\beta du^\beta + \bar{A}_\mu{}^\nu{}_\beta \bar{D} l^\beta) N_\nu^i$$

from which (2.1) follows.

To prove (2.5) \Leftrightarrow (2.10) i. e., $\bar{H}_\alpha^k = 0 \Leftrightarrow \overset{\mu}{N} = 0$ for all \dot{u} and $\mu = m+1, \dots, n$ we have the relation

$$\bar{H}_\beta^i l^\beta = \bar{\theta}_{\alpha\beta}^*{}^\mu l^\alpha l^\beta N_\mu^i = \overset{\mu}{N}(u, \dot{u}) N_\mu^i.$$

3. Case 2). $DB_\alpha^i \in T_V$.

In this case the absolute differentials of tangent and normal vectorc take the form:

$$(3.1) \quad DB_\alpha^i = (\bar{\theta}_{\alpha\beta}^* du^\beta + \bar{A}_{\alpha\beta}^\mu \bar{D}l^\beta)_\mu N^i$$

$$(3.2) \quad DN_\mu^i = (\bar{\theta}_{\mu\beta}^\delta du^\beta + A_{\alpha\beta}^\mu \bar{D}l^\beta) B_\delta^i + (\bar{\lambda}_{\mu\beta}^\nu du^\beta + \bar{A}_{\mu\beta}^\nu \bar{D}l^\beta)_\nu N^i.$$

As in this case $A_{\alpha\beta}^\mu = 0$, we have:

$$(3.3) \quad \bar{A}_{\alpha\beta\gamma} = 2^{-1} L(u, \dot{u}) \dot{\partial}_\gamma g_{\alpha\beta}(u, \dot{u}) = 0,$$

from which we conclude that the metric tensor of the subspace is not a function of the direction \dot{u} , i. e.,

$$g_{\alpha\beta} = g_{\alpha\beta}(u)$$

and the subspace F_m of the Finsler space F_n is Riemannian. From the equations

$$(3.4) \quad \begin{aligned} \Gamma_{\alpha\beta\gamma} + \Lambda_{\alpha\beta\gamma} &= \bar{\Gamma}_{\alpha\beta\gamma}^* - \bar{A}_{\alpha\beta\delta} \bar{A}_{\nu\gamma}^\delta \bar{N}^\nu \\ \bar{A}_{\alpha\beta\gamma} &= 0, \quad \bar{\Gamma}_{\alpha\beta\gamma}^* = 0, \end{aligned}$$

we obtain that in case 2) the intrinsic connection coefficient is the tensor $-\Lambda_{\alpha\beta\gamma}$, i.e. $\Gamma_{\alpha\beta\gamma} = -\Lambda_{\alpha\beta\gamma}$.

The other connection coefficients are obtained from the same formulae as in any other subspace.

Using the equations $\bar{A}_{\alpha\beta}^\delta = 0, \bar{\Gamma}_{\alpha\beta}^\delta = 0$ for case 2) we get

$$(3.5) \quad \begin{aligned} [\Delta D] B_\alpha^i &= \{2^{-1} \bar{\theta}_{\alpha[\beta}^* \bar{\theta}_{|\mu|\gamma]}^\delta [du^\beta \delta u^\gamma] + (\bar{\theta}_{\alpha\beta}^* \bar{A}_{\mu\gamma}^\delta - \bar{\theta}_{\mu\beta}^* \bar{A}_{\alpha\gamma}^\delta) [du^\beta \bar{D}l^\gamma] + \\ &2^{-1} \bar{A}_{\alpha[\beta}^\mu \bar{A}_{|\mu|\gamma]}^\delta [\bar{D}l^\beta \bar{D}l^\gamma]\} B_\delta^i + \{2^{-1} (\partial_{[\gamma} \bar{\theta}_{|\alpha|\beta]}^* + \bar{\theta}_{\alpha[\beta}^* \bar{\lambda}_{|\nu|\gamma]}^\nu) [du^\beta \delta u^\gamma] + \\ &(L \dot{\partial}_\gamma \bar{\theta}_{\alpha\beta}^* - \partial_\beta A_{\alpha\gamma}^\mu - \bar{A}_{\alpha\beta}^\nu \bar{\lambda}_{\nu\gamma}^* + \bar{\theta}_{\alpha\beta}^* \bar{A}_{\mu\gamma}^\nu) [du^\beta \bar{D}l^\gamma] + \\ &2^{-1} (L \dot{\partial}_{[\gamma} \bar{A}_{|\alpha|\beta]}^\mu + \bar{A}_{\alpha[\beta}^\nu \bar{A}_{|\nu|\gamma]}^\mu) [\bar{D}l^\beta \bar{D}l^\gamma]\} N_\mu^i. \\ [\Delta D] N_\mu^i &= \{2^{-1} (\partial_{[\gamma} \bar{\theta}_{|\mu|\beta]}^* + \bar{\theta}_{|\nu|[\gamma}^* \bar{\lambda}_{|\mu|\beta]}^\nu) [du^\beta \delta u^\gamma] + \\ &2^{-1} (L \dot{\partial}_\gamma \bar{\theta}_{\mu\beta}^* - \partial_\beta \bar{A}_{\mu\gamma}^\delta + \bar{A}_{\nu\gamma}^\delta \bar{\lambda}_{\mu\beta}^* - \bar{A}_{\mu\gamma}^\nu \bar{\theta}_{\nu\beta}^*) [du^\beta \bar{D}l^\gamma] + \\ &2^{-1} (L \dot{\partial}_{[\gamma} \bar{A}_{|\mu|\beta]}^\delta + \bar{A}_{\mu[\beta}^\nu \bar{A}_{|\nu|\gamma]}^\delta) [\bar{D}l^\beta \bar{D}l^\gamma]\} B_\beta^i \\ &\{2^{-1} (\partial_{[\gamma} \bar{\lambda}_{|\mu|\beta]}^* + \bar{\theta}_{\mu[\beta}^* \bar{\theta}_{|\delta|\gamma]}^\delta + \bar{\lambda}_{\mu[\beta}^\psi \bar{\lambda}_{|\psi|\gamma]}^*) [du^\beta \delta u^\gamma] + \\ &(L \dot{\partial}_\gamma \bar{\lambda}_{\mu\beta}^* - \partial_\beta A_{\mu\gamma}^\nu + \bar{\lambda}_{\mu\beta}^\psi A_{\psi\gamma}^\nu - \bar{A}_{\mu\gamma}^\psi \bar{\lambda}_{\psi\beta}^* - \bar{A}_{\mu\gamma}^\delta \bar{\theta}_{\delta\beta}^* + \bar{\theta}_{\mu\beta}^* \bar{A}_{\delta\gamma}^\nu) [du^\beta \bar{D}l^\gamma] + \\ &2^{-1} (L \dot{\partial}_{[\gamma} \bar{A}_{|\mu|\beta]}^\nu - A_{\mu[\beta}^\psi \bar{A}_{|\psi|\gamma]}^\nu + \bar{A}_{\mu[\beta}^\delta \bar{A}_{|\delta|\gamma]}^\nu) [du^\beta \bar{D}l^\gamma]. \end{aligned}$$

Comparing the above formulae with those in [6] we obtain that in case 2) the curvature tensors

$${}^0\bar{R}_{\alpha\beta\gamma}^\delta, \quad {}^0\bar{P}_{\alpha\beta\gamma}^\delta, \quad {}^0\bar{S}_{\alpha\beta\gamma}^\delta.$$

and some others are reduced, because of $\bar{\Gamma}_{\alpha\beta}^*{}^\delta = 0$, $\bar{A}_{\alpha\beta}{}^\delta = 0$.

4. Case 3). $DN_\mu^i \in T_H$.

In this case the absolute differentials of tangent and normal vectors take the form:

$$(4.1) \quad DB_\alpha^i = (\bar{\Gamma}_{\alpha\beta}^*{}^\delta du^\beta + \bar{A}_{\alpha\beta}{}^\delta Dl^\beta) B_\delta^i + (\bar{\theta}_{\alpha\beta}^*{}^\mu du^\beta + \bar{A}_{\alpha\beta}{}^\mu Dl^\beta) N_\mu^i,$$

$$(4.2) \quad DN^i = (\bar{\theta}_{\mu\beta}^*{}^\delta du^\beta + \bar{A}_{\mu\beta}{}^\delta Dl^\beta) B_\delta^i$$

Also

$$\bar{\lambda}_{\mu\gamma}^*{}^\nu = 0, \quad \bar{A}_{\mu\gamma}{}^\nu = 0,$$

hence

$$\bar{\lambda}_{\mu\gamma}^*{}^\nu = \overset{\nu}{N}_i (\partial_\gamma N_\mu^i - \partial_\delta N_\mu^i \bar{\Gamma}_{\gamma}^*{}^\delta + \bar{\Gamma}_{jk}^*{}^i N^j B_\mu^k + A_{jk}^i N^j \bar{H}_\mu^k) = 0,$$

$$\bar{A}_{\mu\gamma}{}^\nu = \overset{\nu}{N}_i (L\partial_\gamma N_\mu^i + A_{jk}^i N^j B_\mu^k) = 0.$$

The other connection coefficients we get from the same formulae as in any other subspace.

We also have that the absolute differentials of tangent and normal vectors take the form:

$$\begin{aligned} [\Delta D]B_\alpha^i &= \{2^{-1}({}^0\bar{R}_{\alpha\beta\gamma}{}^\varepsilon + \bar{\theta}_{\alpha[\beta}^*{}^\mu \bar{\theta}_{|\mu|\gamma]}^*{}^\varepsilon)[du^\beta \delta u^\gamma] + \\ &({}^0\bar{P}_{\alpha\beta\gamma}{}^\varepsilon \bar{\theta}_{\alpha\beta}^*{}^\mu \bar{A}_{\mu\gamma}{}^\varepsilon - \bar{A}_{\alpha\gamma}{}^\mu \bar{\theta}_{\mu\beta}^*{}^\varepsilon)[du^\beta \bar{\Delta}l^\gamma] + 2^{-1}{}^0\bar{S}_{\alpha\beta\gamma}{}^\varepsilon \bar{A}_{\alpha[\beta}{}^\nu \bar{A}_{|\nu|\gamma]}{}^\varepsilon][\bar{D}l^\beta \bar{\Delta}l^\gamma]\} B_\varepsilon^i + \\ &2^{-1}{}^0\bar{R}_{\alpha\beta\gamma}{}^\mu [du^\beta \delta u^\gamma] + {}^0\bar{P}_{\alpha\beta\gamma}{}^\mu [du^\beta \Delta u^\gamma] + 2^{-1}{}^0\bar{S}_{\alpha\beta\gamma}{}^\mu [\bar{D}l^\beta \bar{\Delta}l^\gamma] N_\mu^i, \\ [\Delta D]N_\mu^i &= \{2^{-1}({}^0\bar{R}_{\mu\beta\gamma}{}^\varepsilon [du^\beta \delta u^\gamma] + {}^0\bar{P}_{\mu\beta\gamma}{}^\varepsilon [du^\beta \Delta u^\gamma] + {}^0\bar{S}_{\mu\beta\gamma}{}^\varepsilon [du^\beta \bar{\Delta}u^\gamma])\} B_\varepsilon^i + \\ &2^{-1}(\bar{\theta}_{\mu[\beta}^*{}^\varkappa (\bar{\theta}_{|\varkappa|\gamma]}^*{}^\nu [du^\beta \Delta u^\gamma] + (\bar{A}_{\varkappa\gamma}{}^\nu \bar{\theta}_{\mu\beta}^*{}^\varkappa - \bar{A}_{\mu\gamma}{}^\varkappa \bar{\theta}_{\varkappa\beta}^*{}^\nu)[du^\beta \bar{\Delta}l^\gamma] + \\ &2^{-1}\bar{A}_{\mu[\beta}{}^\varkappa \bar{A}_{|\varkappa|\gamma]}{}^\nu + [\bar{D}l^\beta \bar{\Delta}l^\gamma])\} N_\nu^i. \end{aligned}$$

Finally we have

$${}^1\bar{R}_{\mu\beta\gamma}{}^\nu = 0, \quad {}^1\bar{P}_{\mu\beta\gamma}{}^\nu = 0, \quad {}^1\bar{S}_{\mu\beta\gamma}{}^\nu = 0.$$

5. Case 2a) or 3a) $(DB_\alpha^i \in T_V) \wedge (DN^i \in T_H)$.

In this case we have

$$(5.1) \quad \bar{\Gamma}_{\alpha\gamma}^*{}^\beta = 0, \quad \bar{A}_{\alpha\gamma}{}^\beta = 0, \quad \bar{A}_{\mu\beta}{}^\nu = 0, \quad \bar{\lambda}_{\mu\beta}^*{}^\nu = 0$$

and

$$(5.2) \quad DB_\alpha^i = (\bar{\theta}_{\alpha\beta}^*{}^\mu du^\beta + \bar{A}_{\alpha\beta}{}^\mu \bar{D}l^\beta) N_\mu^i$$

$$(5.2) \quad DN_\mu^i = (\bar{\theta}_{\mu\beta}^*{}^\alpha du^\beta + \bar{A}_{\alpha\beta}{}^\mu \bar{D}l^\beta) B_\alpha^i$$

For the absolute differential of tangent and normal vectors we obtain:

$$(5.4) \quad \begin{aligned} [\Delta D]B_\alpha^i &= \{2^{-1}\bar{\theta}_{\alpha[\beta}^*\delta\bar{\theta}_{|\mu|\gamma]}^\delta[du^\beta\delta u^\gamma] + (\bar{\theta}_{\alpha\beta}^*\bar{A}_{\mu\gamma}^\delta - \bar{\theta}_{\mu\beta}^*\bar{A}_{\alpha\gamma}^\mu)[du^\beta\bar{\Delta}l^\gamma] + \\ &\quad 2^{-1}\bar{A}_{\alpha[\beta}^\mu\bar{A}_{|\mu|\gamma]}^\delta[\bar{D}l^\beta\bar{\Delta}l^\gamma]\}B_\delta^i + \\ &\quad \{2^{-1}(\partial_{[\gamma}\bar{\theta}_{|\alpha|\beta]}^*\mu[du^\beta\bar{\Delta}l^\gamma] + (L\dot{\partial}_\gamma\bar{\theta}_{\alpha\beta}^*\mu - \partial_\beta A_{\alpha\beta}^\mu)[du^\beta\bar{\Delta}l^\gamma] + \\ &\quad 2^{-1}(L\dot{\partial}_{[\gamma}\bar{A}_{|\alpha|\beta]}^\mu)[\bar{D}l^\beta\bar{\Delta}l^\gamma]\}N^i, \end{aligned}$$

$$(5.5) \quad \begin{aligned} [\Delta D]N_\mu^i &= \{2^{-1}(\partial_{[\gamma}\bar{\theta}_{|\mu|\beta]}^*\alpha[du^\beta\delta u^\gamma] + (L\dot{\partial}_\gamma\bar{\theta}_{\mu\beta}^*\alpha - \partial_\beta\bar{A}_{\mu\delta}^\alpha)[du^\beta\bar{\Delta}l^\gamma] + \\ &\quad 2^{-1}L\dot{\partial}_{[\gamma}\bar{A}_{|\mu|\beta]}^\alpha[\bar{D}l^\beta\bar{\Delta}l^\gamma]\}B_\alpha^i + \\ &\quad \{2^{-1}\bar{\theta}_{\mu[\beta}^*\delta\bar{\theta}_{|\delta|\gamma]}^\nu[du^\beta\delta u^\gamma] + (\bar{\theta}_{\mu\beta}^*\bar{A}_{\delta\gamma}^\nu - \bar{\theta}_{\delta\beta}^*\bar{A}_{\mu\gamma}^\delta)[du^\beta\bar{\Delta}l^\gamma] + \\ &\quad 2^{-1}A_{\mu[\beta}^\delta\bar{A}_{|\delta|\gamma]}^\nu[\bar{D}l^\beta\bar{\Delta}l^\gamma]\}N_\nu^i. \end{aligned}$$

We also have:

$$\begin{aligned} {}^0\bar{R}_{\mu\beta\gamma}^\delta &= 0, & {}^0\bar{P}_{\mu\beta\gamma}^\delta &= 0, & {}^0\bar{S}_{\mu\beta\gamma}^\delta &= 0 \\ {}^1\bar{R}_{\mu\beta\gamma}^\nu &= 0, & {}^1\bar{P}_{\mu\beta\gamma}^\nu &= 0, & {}^1\bar{S}_{\mu\beta\gamma}^\nu &= 0. \end{aligned}$$

The intrinsic connection coefficients are:

$$\Gamma_{\alpha\beta\gamma}^* = -\Lambda_{\alpha\beta\gamma}, \quad \lambda_{\mu\nu\beta} = 0, \quad A_{\mu\nu\beta} = 0, \quad A_{\alpha\beta\gamma} = 0$$

and the corresponding equations for the intrinsic curvature tensors are the same as (4.6), (4.7) except for

$${}^0\bar{R}_{\mu\beta\gamma}^\delta = -\partial_{[\gamma}\Lambda_{|\alpha|\beta]}^\delta + \Lambda_{\alpha[\beta}^\varkappa + \Lambda_{|\varkappa|\gamma]}^\delta.$$

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