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**COMPLEX ANALYSIS METHODS IN THE THEORY OF  
INFINITESIMAL BENDINGS OF SURFACES WITH A FLAT POINT**

ABSTRACT. Using I. Vekua's analytic methods, the problem of one-to-one correspondence between infinitesimal bendings of surfaces with a flat point is studied.

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**1. Introduction.** In the paper the objects of study are two surfaces  $S_0$  and  $S$  given in a rectangular Cartesian coordinate system  $Ox_1y_1z_1$  by the equations  $S_0 : z_1 = (x_1^2 + y_1^2)^{n/2}$ ,  $S : z_1 = (x_1^2 + y_1^2)^{n/2}f(x_1, y_1)$ . It is assumed that  $S_0$  and  $S$  are defined in a domain  $G_1$ ,  $(0, 0) \in G_1$ ,  $f(x_1, y_1) \in C^3(G_1)$ ,  $f(0, 0) > 0$ ,  $n$  ( $n > 2$ ) is any real number, and for all points of  $G_1$  other than  $(0, 0)$ , the gaussian curvature of  $S$  is positive.

It is clear that the point  $(0, 0)$  is a flat one on the surfaces  $S_0$  and  $S$ . At it not only the Gaussian curvature but also all coefficients of the second quadratic forms of  $S_0$  and  $S$  vanish. At this point the surfaces have with their tangent planes a contact order greater than 1. We call the surface  $S_0$  *model* with respect to  $S$  as it is a particular case of  $S$  and can be obtained from  $S$  under the condition  $f(x_1, y_1) = 1$ .

The aim of the paper is to establish the following result.

**Theorem 1.** *There exists a one-to-one correspondence between the sets of continuous infinitesimal bendings of the surfaces  $S_0$  and  $S$ .*

**2. An equivalent analytic problem [1].** We extend the I.Vekua analytic methods on investigating infinitesimal bendings of the above surfaces [2]. On  $S_0$  and  $S$ , we introduce a conjugate isometric parametrization  $z = x + iy$ ,  $i^2 = -1$ . Then infinitesimal bendings of these surfaces will be characterized by the functions  $\Phi(z) = z^2K_0^{1/4}(z)(\delta M_0 + i\delta L_0)$ ,  $w(z) = z^2K^{1/4}(z)(\delta M + i\delta L)$ , where  $K_0(z)$  and  $K(z)$  are the Gaussian

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curvatures of  $S_0$  and  $S$ , and  $\delta M_0, \delta L_0$  and  $\delta M, \delta L$  are variations of the coefficients of the second quadratic forms of  $S_0$  and  $S$ , respectively.

In the domain  $G$  which is the image of the domain  $G_1$  by the mapping  $z = z(x_1, y_1)$ , these functions satisfy the following generalized Cauchy–Riemann systems:

$$2\bar{z}\partial_{\bar{z}}\Phi - b(0)\bar{\Phi} = 0, \quad (1)$$

$$2\bar{z}\partial_{\bar{z}}w - b(z)\bar{w} = 0. \quad (2)$$

Here the singular point  $z=0$  belongs to the domain  $G$ ,  $b(0) = (n-2)/2\sqrt{n-1}$  and  $b(z)$  is a continuous function in  $G$  satisfying  $|b(z) - b(0)| < M|z|^\alpha$  at least in a sufficiently small neighbourhood of  $z = 0$ ;  $M, \alpha$  are positive constants.

Continuous solutions of the systems (1) and (2) are connected by the two-dimensional integral equation

$$w(z) = \Phi(z) + P_G\bar{w}, \quad (3)$$

where  $P_G\bar{w} = S_G\left(\frac{b(\zeta)-b(0)}{2\bar{\zeta}}\overline{w(\zeta)}\right)$  and

$$S_G f = -\frac{1}{\pi} \iint_G \left[ \frac{\Omega_1(z, \zeta)}{\zeta} f(\zeta) + \frac{\Omega_2(z, \zeta)}{\bar{\zeta}} \overline{f(\zeta)} \right] d\xi d\eta.$$

Here  $\zeta = \xi + i\eta$  and  $\Omega_1, \Omega_2$  are certain functions presented in [1]. It is necessary to note that  $P_G$  is a completely continuous operator mapping the class  $C(G)$  of continuous functions in itself. According to Fredholm's alternatives, the equation (3) will be uniquely solvable, and consequently a one-to-one correspondence between continuous infinitesimal bendings of the surfaces  $S_0$  and  $S$  will exist if the following assertion takes place.

**Theorem 2.** *The homogeneous equation*

$$w^+(z) = P_G\bar{w}^+, \quad z \in G, \quad (4)$$

*in the class  $C(G)$  has only the zero solution.*

*Scheme of proof of Theorem 2.* Suppose that the equation (4) has a non-trivial solution  $w^+(z)$ ,  $z \in G$ . Let us note some of its properties. First, we can check that any continuous solution  $w^+(z)$  of the equation (4) belongs to the class  $D_{1,p}(G)$ ,  $p > 2$ , and satisfies the equation (2). In this case, as was shown in [2] (see Theorem 1.1, p.74), we have

$$w^+(z) = O(|z|^{b(0)}) \quad \text{as } z \rightarrow 0, \quad (5)$$

From (4) it follows that  $w^+(z)$  is continuously extended to the domain  $E \setminus \bar{G}$  ( $E$  is the  $z$ -plane and  $\bar{G}$  is the closure of  $G$ ) by a continuous function  $w^-(z)$ ,  $z \in E \setminus \bar{G}$ , satisfying the equation (1). According to the theory of elliptic

systems,  $w^-(z)$  is an analytic function in  $E \setminus \overline{G}$  with respect to  $z$  and  $\bar{z}$ . In addition, it is established that

$$|w^-(z)| < M \| [b(\zeta) - b(0)] / 2\bar{\zeta} \|_{L_p} |z|^{-|b(0)|}, \quad (6)$$

where  $M = M(|b(0)|, R_0)$  is a constant depending on  $|b(0)|$  and  $R_0$  (a maximal distance from  $z = 0$  to the boundary of  $G$ ), and  $\|\cdot\|_{L_p}$  denotes the norm of a function  $f(z)$  in the space  $L_p(\overline{G})$ .

Thus on the plane  $z$  the continuous function  $W(z) = \begin{cases} w^+(z), & z \in \overline{G}, \\ w^-(z), & z \in E \setminus \overline{G} \end{cases}$  is defined. This function is subject to the conditions (5),(6) and satisfies the equation

$$2\bar{z}\partial_{\bar{z}}W - B(z)\overline{W} = 0, \quad (7)$$

in which

$$B(z) = \begin{cases} b(z) & \text{for } z \in \overline{G}, \\ b(0) & \text{for } z \in E \setminus \overline{G}. \end{cases}$$

We establish the following.

**Lemma 1.** *If  $W(z)$  is a function satisfying the above properties, then  $W(z) \equiv 0, z \in E$ .*

From this lemma it follows that  $w^+(z) = 0, z \in G$  and therefore Theorems 1 and 2 are proved.

**4. Generalization.** Now we consider a surface given in Cartesian coordinates  $Ox_1y_1z_1$  by the equation  $z_1 = \sum_{k=0}^n a_{k,n-k} x_1^k y_1^{n-k} + R(x_1, y_1)$ , where  $n$  ( $n \geq 3$ ) is an integer and  $a_{k,n-k}$  are constants. Let  $R(x_1, y_1)$  be a sufficiently regular function; moreover, let  $R(x_1, y_1) = O[(x_1^2 + y_1^2)^{(n+1)/2}]$  as  $x, y \mapsto 0$ . Passing over to polar coordinates ( $x_1 = r_1 \cos \varphi, y_1 = r_1 \sin \varphi$ ), we write the surface equation in the form

$$S: \quad z_1 = r_1^n f(\varphi) + R(x_1, y_1), \quad (8)$$

where  $f(\varphi) = \sum_{k=0}^n a_{k,n-k} (\cos \varphi)^k (\sin \varphi)^{n-k}$ . The requirement of positiveness of the curvature in a neighbourhood of the point  $(0, 0)$  imposes on  $f(\varphi)$  the restriction

$$-(n-1) \left( \frac{df}{d\varphi} \right)^2 + n f \frac{d^2 f}{d\varphi^2} + n^2 f^2 > 0. \quad (9)$$

Besides, we assume  $f(\varphi) > 0$ .

The first summand in the right side of (8) defines the structure of the surface in a neighbourhood of the flat point. The model surface

$$S_0: \quad z_1 = r_1^n f(\varphi) \quad (10)$$

corresponds to it.

Further we will consider the surfaces (8) and (10) under wider assumptions on  $n$  and  $f(\varphi)$ . We assume that  $n$  ( $n > 2$ ) is a real number,  $f(\varphi)$  is a  $2\pi$ -periodic function from the class  $C^3[0, 2\pi]$ , satisfying the inequality (9). It is clear that the surfaces under study are objects with sufficiently general and more complicated structure for a neighbourhood of the flat point  $(0, 0)$  compared with those which have been discussed previously.

The problem is to establish a one-to-one correspondence between infinitesimal bendings of surfaces (8) and (10).

As was stated in [3], in a conjugate isometric parametrization  $z = x + iy$  and in terms of complex-valued functions  $\Phi(z)$  and  $w(z)$  introduced earlier, infinitesimal bendings of those surfaces are described by the equations

$$2\bar{z}\partial_{\bar{z}}\Phi - b_0(\varphi)\bar{\Phi} = 0, \quad (11)$$

$$2\bar{z}\partial_{\bar{z}}w - [b_0(\varphi) + B(z)]\bar{w} = 0, \quad (12)$$

where the point  $z = 0$  is interior for the domain  $G$ ,  $b_0(\varphi)$  is a  $2\pi$ -periodic continuous function and  $B(z)$  is continuous in  $G$ , moreover  $B(z) = O(|z|^\alpha)$  as  $z \rightarrow 0$ ,  $\alpha > 0$ .

Apparently, the model equation (11), seeming simpler in comparison with (12), is nevertheless fairly complicated for investigating. This fact is maybe a main reason why a progress in this respect looks such moderate, see [4,5], and the problem of correspondence for infinitesimal bendings of the surfaces (8) and (10) remains unsolved.

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